# EXOTIC DEFORMATION QUANTIZATION 

VALENTIN OVSIENKO

## 1. Introduction

Let $\mathcal{A}$ be one of the following commutative associative algebras: the algebra of all smooth functions on the plane: $\mathcal{A}=C^{\infty}\left(\mathbf{R}^{2}\right)$, or the algebra of polynomials $\mathcal{A}=\mathbf{C}[p, q]$ over $\mathbf{R}$ or $\mathbf{C}$. There exists a nontrivial formal associative deformation of $\mathcal{A}$ called the Moyal $\star$-product (or the standard $\star$-product). It is defined as an associative operation $\mathcal{A}^{\otimes 2} \rightarrow \mathcal{A}[[\hbar]]$ where $\hbar$ is a formal variable. The explicit formula is:

$$
\begin{equation*}
F \star_{\hbar} G=F G+\sum_{k \geq 1} \frac{(i \hbar)^{k}}{2^{k} k!}\{F, G\}_{k}, \tag{1}
\end{equation*}
$$

where $\{F, G\}_{1}=\frac{\partial F}{\partial p} \frac{\partial G}{\partial q}-\frac{\partial F}{\partial q} \frac{\partial G}{\partial p}$ is the standard Poisson bracket, and the higher order terms are:

$$
\begin{equation*}
\{F, G\}_{k}=\sum_{i=0}^{k}(-1)^{i}\binom{k}{i} \frac{\partial^{k} F}{\partial p^{k-i} \partial q^{i}} \frac{\partial^{k} G}{\partial p^{i} \partial q^{k-i}} . \tag{2}
\end{equation*}
$$

The Moyal product is the unique (modulo equivalence) non-trivial formal deformation of the associative algebra $\mathcal{A}$ (see [13]).

Definition 1. A formal associative deformation of $\mathcal{A}$ given by formula (1) is called a $\star$-product if the following hold:

1) the first order term coincides with the Poisson bracket: $\{F, G\}_{1}=$ $\{F, G\} ;$

[^0]2) the higher order terms $\{F, G\}_{k}$ are given by differential operators vanishing on constants: $\{1, G\}_{k}=\{F, 1\}_{k}=0$;
3) $\{F, G\}_{k}=(-1)^{k}\{G, F\}_{k}$.

Definition 2. Two $\star$-products $\star_{\hbar}$ and $\star_{\hbar}^{\prime}$ on $\mathcal{A}$ are called equivalent if there exists a linear mapping $A_{\hbar}: \mathcal{A} \rightarrow \mathcal{A}[[\hbar]]$ such that

$$
A_{\hbar}(F)=F+\sum_{k=1}^{\infty} A_{k}(F) \hbar^{k}
$$

intertwining the operations $\star_{\hbar}$ and $\star_{\hbar}^{\prime}: A_{\hbar}(F) \star_{\hbar}^{\prime} A_{\hbar}(G)=A_{\hbar}\left(F \star_{\hbar} G\right)$.
Consider now $\mathcal{A}$ as a Lie algebra; the commutator is given by the Poisson bracket. The Lie algebra $\mathcal{A}$ has a unique (modulo equivalence) non-trivial formal deformation called the Moyal bracket or the Moyal *-commutator: $\{F, G\}_{t}=\frac{1}{i \hbar}\left(F \star_{\hbar} G-G \star_{\hbar} F\right)$, where $t=-\hbar^{2} / 2$.

The well-known De Wilde-Lecomte theorem [4] states the existence of a non-trivial $\star$-product for an arbitrary symplectic manifold. The theory of $\star$-products is a subject of deformation quantization. The geometrical proof of the existence theorem was given by B. Fedosov [9] (see [8] and [20] for clear explanation and survey of recent progress).

The main idea of this paper is to consider the algebra $\mathcal{F}(M)$ of functions (with singularities) on the cotangent bundle $T^{*} M$ which are Laurent polynomials on the fibers. In contrast to the above algebra $\mathcal{A}$ it turns out that for such algebras the standard $\star$-product is no more unique at least if $M$ is one-dimensional: $\operatorname{dim} M=1$.

We consider $M=S^{1}, \mathbf{R}$ in the real case, and $M=\mathcal{H}$ (the upper half-plane) in the holomorphic case. The main result of this paper is an explicit construction of a new $\star$-product on the algebra $\mathcal{F}(M)$ nonequivalent to the standard Moyal product. This $\star$-product is equivariant with respect to the Möbius transformations. The construction is based on the bilinear $S L_{2}$-equivariant operations on tensor-densities on $M$, known as Gordan transvectants and Rankin-Cohen brackets.

We study the relations between the new $\star$-product and extensions of the Lie algebra $\operatorname{Vect}\left(S^{1}\right)$.

The results of this paper are closely related to those of the recent work of P. Cohen, Yu. Manin and D. Zagier [3] where a one-parameter
family of associative products on the space of classical modular forms is constructed using the same $S L_{2}$-equivariant bilinear operations.

## 2. Definition of the exotic $\star$-product

2.1 Algebras of Laurent polynomials. Let $\mathcal{F}$ be one of the following associative algebras of functions:

$$
\mathcal{F}=\mathbf{C}[p, 1 / p] \otimes C^{\infty}(\mathbf{R}) \quad \text { or } \quad \mathcal{F}=\mathbf{C}[p, 1 / p] \otimes \operatorname{Hol}(\mathcal{H})
$$

This means, it consists of functions of the type:

$$
\begin{equation*}
F(p, q)=\sum_{i=-N}^{N} p^{i} f_{i}(q) \tag{3}
\end{equation*}
$$

where $f_{i}(q) \in C^{\infty}(\mathbf{R})$ in the real case, or $f_{i}(q)$ are holomorphic functions on the upper half-plane $\mathcal{H}$ or $f_{i} \in C[q]$ (respectively).

We will also consider the algebra of polynomials: $\mathbf{C}[p, 1 / p, q]$ (Laurent polynomials in $p$ ).
2.2 Transvectants. Consider the following bilinear operators on functions of one variable:

$$
\begin{equation*}
J_{k}^{m, n}(f, g)=\sum_{i+j=k}(-1)^{i}\binom{k}{i} \frac{(2 m-i)!(2 n-j)!}{(2 m-k)!(2 n-k)!} f^{(i)} g^{(j)} \tag{4}
\end{equation*}
$$

where $f=f(z), g=g(z), f^{(i)}(z)=\frac{d^{i} f(z)}{d z^{i}}$.
These operators satisfy a remarkable property: they are equivariant under Möbius (linear-fractional) transformations. Namely, suppose that the transformation $z \mapsto \frac{a z+b}{c z+d}$ (with $a d-b c=1$ ) acts on the arguments as follows:

$$
f(z) \mapsto f\left(\frac{a z+b}{c z+d}\right)(c z+d)^{2 m}, \quad g(z) \mapsto g\left(\frac{a z+b}{c z+d}\right)(c z+d)^{2 n}
$$

then $J_{k}^{m, n}(f, g)$ transforms as:

$$
J_{k}^{m, n}(f, g)(z) \mapsto J_{k}^{m, n}(f, g)\left(\frac{a z+b}{c z+d}\right)(c z+d)^{2(m+n-k)}
$$

In other words, the operations (4) are bilinear $S L_{2}$-equivariant mappings on tensor-densities:

$$
J_{k}^{m n}: \mathcal{F}_{m} \otimes \mathcal{F}_{n} \rightarrow \mathcal{F}_{m+n-k}
$$

where $\mathcal{F}_{l}$ is the space of tensor-densities of degree $-l: \phi=\phi(z)(d z)^{-l}$.
The operations (4) were discovered more than one hundred years ago by Gordan [11] who called them the transvectants. They have been rediscovered many times: in the theory of modular functions by Rankin [18] and by Cohen [2] (so-called Rankin-Cohen brackets), in differential projective geometry by Janson and Peetre [12]. The "multi-dimensional transvectants" were defined in [14] in the context of the the Virasoro algebra and symplectic and contact geometry.
2.3 Main definition. Define the following bilinear mapping $\mathcal{F}^{\otimes 2} \rightarrow$ $\mathcal{F}[[\hbar]]$, for $F=p^{m} f(q), G=p^{n} g(q)$, where $m, n \in \mathbf{Z}$, by putting:

$$
\begin{equation*}
F \widetilde{\star}_{h} G=\sum_{k=0}^{\infty} \frac{(i \hbar)^{k}}{2^{2 k}} p^{(m+n-k)} J_{k}^{m, n}(f, g), \tag{5}
\end{equation*}
$$

Note, that the first order term coincides with the Poisson bracket.
This operation will be the main subject of this paper. We call it the exotic *-product.
2.4 Remark. Another one-parameter family of operations on modular forms: $f \star^{\kappa} g=\sum_{n=0}^{\infty} t_{n}^{\kappa}(k, l) J_{n}^{k l}(f, g)$, where $f$ and $g$ are modular forms of weight $k$ and $l$ respectively, and $t_{n}^{\kappa}(k, l)$ are very interesting and complicated coefficients, is defined in [3].

## 3. Main theorems

We formulate here the main results of this paper. All the proofs will be given in Sections 4-7.
3.1 Non-equivalence. The Moyal $\star$-product (1) defines a nontrivial formal deformation of $\mathcal{F}$. We will show that the formula (5) defines a $\star$-product non-equivalent to the standard Moyal product.

Theorem 1. The operation (5) is associative; it defines a formal deformation of the algebra $\mathcal{F}$ which is not equivalent to the Moyal product.

The associativity of the product (5) is a trivial corollary of Proposition 1 below. To prove the non-equivalence, we will use the relations with extensions of the Lie algebra of vector fields on $S^{1}: \operatorname{Vect}\left(S^{1}\right) \subset \mathcal{F}$ (cf. Sec.5).

It is interesting to note that the constructed $\star$-product is equivalent to the standard Moyal product if we consider it on the algebra $C^{\infty}\left(T^{*} M \backslash M\right)$ of all smooth functions (not only Laurent polynomials on fibers) ; cf. Corollary 1 below.
$3.2 s l_{2}$-equivariance. The Lie algebra $s l_{2}(\mathbf{R})$ has two natural embeddings into the Poisson Lie algebra on $\mathbf{R}^{2}$ : the symplectic Lie algebra $s p_{2}(\mathbf{R}) \cong s l_{2}(\mathbf{R})$ generated by quadratic polynomials $\left(p^{2}, p q, q^{2}\right)$ and another one with generators: $\left(p, p q, p q^{2}\right)$ which is called the Möbius algebra.

It is well-known that the Moyal product (1) is the unique non-trivial formal deformation of the associative algebra of functions on $\mathbf{R}^{2}$ equivariant under the action of the symplectic algebra. This means, (1) satisfies the Leibnitz property:

$$
\begin{equation*}
\left\{F, G \star_{\hbar} H\right\}=\{F, G\} \star_{\hbar} H+G \star_{\hbar}\{F, H\} \tag{6}
\end{equation*}
$$

where $F$ is a quadratic polynomial (note that $\{F, G\}_{t}=\{F, G\}$ if $F$ is a quadratic polynomial).

Theorem 2. The product (5) is the unique formal deformation of the associative algebra $\mathcal{F}$ equivariant under the action of the Möbius algebra.

The product (5) is the unique non-trivial formal deformation of $\mathcal{F}$ satisfying (6) for $F$ from the Möbius $s l_{2}$ algebra.
3.3 Symplectomorphism $\Phi$. The relation between the Moyal product and the product (5) is as follows. Consider the symplectic mapping

$$
\begin{equation*}
\Phi(p, q)=\left(\frac{p^{2}}{2}, \frac{q}{p}\right) \tag{7}
\end{equation*}
$$

defined on $\mathbf{R}^{2} \backslash \mathbf{R}$ in the real case and on $\mathcal{H}$ in the complex case.
Proposition 1. The product (5) is the $\Phi$-conjugation of the Moyal product:

$$
\begin{equation*}
F \widetilde{\star}_{\hbar} G=F \star_{\hbar}^{\Phi} G:=\left(F \circ \Phi \star_{\hbar} G \circ \Phi\right) \circ \Phi^{-1} . \tag{8}
\end{equation*}
$$

Remark. The mapping (7) (in the complex case) can be interpreted as follows. It transforms the space of holomorphic tensor-densities of
degree $-k$ on $\mathbf{C} \mathbf{P}^{1}$ to the space $\mathbf{C}^{k}[p, q]$ of polynomials of degree $k$. Indeed, there exists a natural isomorphism $z^{n}(d z)^{-m} \mapsto p^{m} q^{n}$ (where $m \geq 2 n)$ and $\left(p^{m} q^{n}\right) \circ \Phi=p^{2 m-n} q^{n}$.
3.4 Operator formalism. The Moyal product is related to the following Weil quantization procedure. Define the following differential operators:

$$
\begin{align*}
& \widehat{p}=i \hbar \frac{\partial}{\partial q},  \tag{9}\\
& \widehat{q}=q
\end{align*}
$$

satisfying the canonical relation: $[\widehat{p}, \widehat{q}]=i \hbar \mathbf{I}$. Associate to each polynomial $F=F(p, q)$ the differential operator $\widehat{F}=\operatorname{Sym} F(\widehat{p}, \widehat{q})$ symmetric in $\widehat{p}$ and $\widehat{q}$. The Moyal product on the algebra of polynomials coincides with the product of differential operators: $\widehat{F \star_{\hbar} G}=\widehat{F} \widehat{G}$.

We will show that the $\star$-product (5) leads to the operators:

$$
\begin{align*}
& \widehat{p}^{\Phi}=\left(\frac{i \hbar}{2}\right)^{2} \Delta \\
& \widehat{q}^{\Phi}=\frac{1}{4 i \hbar}\left(\Delta^{-1} \circ A+A \circ \Delta^{-1}\right) \tag{10}
\end{align*}
$$

(where $\Delta=\frac{\partial^{2}}{\partial q^{2}}$ and $A=2 q \frac{\partial}{\partial q}+1$ is the dilation operator) also satisfying the canonical relation.

Remark that $\widehat{p}^{\Phi}$ and $\widehat{q}^{\Phi}$ given by (10) on the Hilbert space $L_{2}(\mathbf{R})$ are not equivalent to the operators (9) since $\widehat{q}^{\Phi}$ is symmetric but not self-adjoint (see [6] on this subject).
3.5 "Symplectomorphic" deformations. Let us consider the general situation.

Proposition 2. Given a symplectic manifold $V$ endowed with a $\star$-product $\star_{\hbar}$ and a symplectomorphism $\Psi$ of $V$, if there exists a hamiltonian isotopy of $\Psi$ to the identity, then the $\Psi$-conjugate product $\star_{\hbar}^{\Psi}$ defined according to the formula (8) is equivalent to $\star_{\hbar}$.

Corollary 1. The $\star$-product (5) considered on the algebra of all smooth functions $C^{\infty}\left(T^{*} \mathbf{R} \backslash \mathbf{R}\right)$ is equivalent to the Moyal product.

## 4. Möbius-invariance

In this section we prove Theorem 2 . We show that the operations of transvectant (4) are $\Phi$-conjugate of the terms of the Moyal product.
4.1 Lie algebra $\operatorname{Vect}(\mathbf{R})$ and modules of tensor-densities. Let $\operatorname{Vect}(\mathbf{R})$ be the Lie algebra of smooth (or polynomial) vector fields on R:

$$
X=X(x) \frac{d}{d x}
$$

with the commutator

$$
\left[X(x) \frac{d}{d x}, Y(x) \frac{d}{d x}\right]=\left(X(x) Y^{\prime}(x)-X^{\prime}(x) Y(x)\right) \frac{d}{d x} .
$$

The natural embedding of the Lie algebra $s l_{2} \subset \operatorname{Vect}(\mathbf{R})$ is generated by the vector fields $d / d x, x d / d x, x^{2} d / d x$.

Define a 1-parameter family of $\operatorname{Vect}(\mathbf{R})$-actions on $C^{\infty}(\mathbf{R})$ given by

$$
\begin{equation*}
L_{X}^{(\lambda)} f=X(x) f^{\prime}(x)-\lambda X^{\prime}(x) f(x), \tag{11}
\end{equation*}
$$

where $\lambda \in \mathbf{R}$. Geometrically, $L_{X}^{(\lambda)}$ is the operator of Lie derivative on tensor-densities of degree $-\lambda$ :

$$
f=f(x)(d x)^{-\lambda} .
$$

Denote $\mathcal{F}_{\lambda}$ the $\operatorname{Vect}(\mathbf{R})$-module structure on $C^{\infty}(\mathbf{R})$ given by (11).
4.2 Transvectant as a bilinear $s l_{2}$-equivariant operator. The operations (4) can be defined as bilinear mappings on $C^{\infty}(\mathbf{R})$ which are sl2-equivariant:

Statement 4.1. For each $k=0,1,2, \ldots$ there exists a unique (up to a constant) bilinear $s l_{2}$-equivariant mapping

$$
\mathcal{F}_{\mu} \otimes \mathcal{F}_{\nu} \rightarrow \mathcal{F}_{\mu+\nu-k}
$$

It is given by $f \otimes g \mapsto J_{k}^{\mu, \nu}(f, g)$.
Proof. Straightforward (cf. [11], [12]).
4.3 Algebra $\mathcal{F}$ as a module over $\operatorname{Vect}(\mathbf{R})$. The Lie algebra $\operatorname{Vect}(\mathbf{R})$ can be considered as a Lie subalgebra of $\mathcal{F}$. The embedding $\operatorname{Vect}(\mathbf{R}) \subset \mathcal{F}$ is given by:

$$
X(x) \frac{d}{d x} \mapsto p X(q) .
$$

The algebra $\mathcal{F}$ is therefore, a $\operatorname{Vect}(\mathbf{R})$-module.

Lemma 4.2. The algebra $\mathcal{F}$ is decomposed to a direct sum of $\operatorname{Vect}(\mathbf{R})$-modules:

$$
\mathcal{F}=\oplus_{m \in \mathbf{Z}} \mathcal{F}_{m}
$$

Proof. Consider the subspace of $\mathcal{F}$ consisting of functions homogeneous of degree $m$ in $p: F=p^{m} f(q)$. This subspace is a $\operatorname{Vect}(\mathbf{R})$-module isomorphic to $\mathcal{F}_{m}$. Indeed, $\left\{p X(q), p^{m} f(q)\right\}=p^{m}\left(X f^{\prime}-m X^{\prime} f\right)=$ $p^{m} L_{X}^{(m)} f$.
4.4 Projective property of the diffeomorphism $\Phi$. The transvectants (4) coincide with the $\Phi$-conjugate operators (2) from the Moyal product:

Proposition 4.3. Let $F=p^{m} f(q), G=p^{n} g(q)$. Then

$$
\begin{equation*}
\Phi^{*-1}\left\{\Phi^{*} F, \Phi^{*} G\right\}_{k}=\frac{k!}{2^{k}} p^{m+n-k} J_{k}^{m, n}(f, g) . \tag{12}
\end{equation*}
$$

Proof. The symplectomorphism $\Phi$ of $\mathbf{R}^{2}$ intertwines the symplectic algebra $s p_{2} \equiv s l_{2}$ and the Möbius algebra: $\Phi^{*}\left(p, p q, p q^{2}\right)=$ $\left(\frac{1}{2} p^{2}, \frac{1}{2} p q, \frac{1}{2} q^{2}\right)$. Therefore, the operation $\Phi^{*-1}\left\{\Phi^{*} F, \Phi^{*} G\right\}_{k}$ is Möbiusequivariant.

On the other hand, one has: $\Phi^{*} F=\frac{1}{2^{m}} p^{2 m} f\left(\frac{q}{p}\right)$ and $\Phi^{*} G=\frac{1}{2^{n}} p^{2 m} g\left(\frac{q}{p}\right)$. Since $\Phi^{*} F$ and $\Phi^{*} G$ are homogeneous of degree $2 m$ and $2 n$ (respectively), the function $\left\{\Phi^{*} F, \Phi^{*} G\right\}_{k}$ is also homogeneous of degree $2(m+$ $n-k)$. Thus, the operation $\{F, G\}_{k}^{\Phi}=\Phi^{*-1}\left\{\Phi^{*} F, \Phi^{*} G\right\}_{k}$ defines a bilinear mapping on the space of tensor-densities $\mathcal{F}_{m} \otimes \mathcal{F}_{n} \rightarrow \mathcal{F}_{m+n-k}$ which is $s l_{2}$-equivariant.

Statement 4.1 implies that it is proportional to $J_{k}^{m, n}$. One easily verifies the coefficient of proportionality for $F=p^{m}, G=p^{n} q^{k}$, to obtain the formula (12).

Proposition 4.3 is proven.
Remark. Proposition 4.3 was proven in [15]. We do not know whether this elementary fact has been mentioned by classics.
4.5 Proof of Theorem 2. Proposition 4.3 implies that the formula (5) is a $\Phi$-conjugation of the Moyal product and is given by the formula (8).

Proposition 1 is proven.

It follows that (5) is a $\star$-product on $\mathcal{F}$ equivariant under the action of the Möbius $s l_{2}$ algebra. Moreover, it is the unique $\star$-product with this property since the Moyal product is the unique $\star$-product equivariant under the action of the symplectic algebra.

Theorem 2 is proven.

## 5. Relation with extensions of the Lie algebra Vect $\left(S^{1}\right)$

We prove here that the $\star$-product (5) is not equivalent to the Moyal product.

Let $\operatorname{Vect}\left(S^{1}\right)$ be the Lie algebra of vector fields on the circle. Consider the embedding $\operatorname{Vect}\left(S^{1}\right) \subset \mathcal{F}$ given by functions on $\mathbf{R}^{2}$ of the type: $X=p X(q)$ where $X(q)$ is periodical: $X(q+1)=X(q)$.
5.1 An idea of the proof of Theorem 1. Consider the formal deformations of the Lie algebra $\mathcal{F}$ associated to the $\star$-products (1) and (5). The restriction of the Moyal bracket to $\operatorname{Vect}\left(S^{1}\right)$ is identically zero. We show that the restriction of the $*$-commutator

$$
\widetilde{\{F, G\}_{t}}=\frac{1}{i \hbar}\left(F \widetilde{\star}_{\hbar} G-G \widetilde{\star}_{\hbar} F\right), \quad t=-\frac{\hbar^{2}}{2}
$$

associated to the $\star$-product (5) defines a series of non-trivial extensions of the Lie algebra $\operatorname{Vect}\left(S^{1}\right)$ by the modules $\mathcal{F}_{k}\left(S^{1}\right)$ of tensor-densities on $S^{1}$ of degree $-k$.
5.2 Extensions and the cohomology group $H^{2}\left(\operatorname{Vect}\left(S^{1}\right) ; \mathcal{F}_{\lambda}\right)$. Recall that an extension of a Lie algebra by its module is defined by a 2-cocycle on it with values in this module. To define an extension of $\operatorname{Vect}\left(S^{1}\right)$ by the module $\mathcal{F}_{\lambda}$ one needs therefore a bilinear mapping $c: \operatorname{Vect}\left(S^{1}\right)^{\otimes 2} \rightarrow \mathcal{F}_{\lambda}$ which satisfies the identity $\delta c=0$ :

$$
c(X,[Y, Z])+L_{X}^{(\lambda)} c(Y, Z)+\left(\text { cycle }_{X, Y, Z}\right)=0
$$

(See [10]).
The cohomology group $H^{2}\left(\operatorname{Vect}\left(S^{1}\right) ; \mathcal{F}_{\lambda}\right)$ were calculated in [19] (see [10]). This group is trivial for each value of $\lambda$ except $\lambda=0,-1,-2,-5$, -7 . The explicit formulæ for the corresponding non-trivial cocycles are given in [17]. If $\lambda=-5,-7$, then $\operatorname{dim} H^{2}\left(\operatorname{Vect}\left(S^{1}\right) ; \mathcal{F}_{\lambda}\right)=1$, the cohomology group is generated by the unique (up to equivalence) nontrivial cocycle. We will obtain these cocycles from the $\star$-commutator.
5.3 Non-trivial cocycles on $\operatorname{Vect}\left(S^{1}\right)$.

Consider the restriction of the $\star$-commutator $\widetilde{\{,\}_{t}}$ (corresponding to the $\star$-product (5)) to $\operatorname{Vect}\left(S^{\mathbf{1}}\right) \subset \mathcal{F}$ : let

$$
X=p X(q), \quad Y=p Y(q)
$$

then from (5) we have

$$
\{X, Y\}_{t}=\{X, Y\}+\sum_{k=1}^{\infty} \frac{t^{k}}{2^{2 k+1}} \frac{1}{p^{2 k-1}} J_{2 k+1}^{1,1}(X, Y)
$$

It follows from the Jacobi identity that the first non-zero term of the series $\{\widetilde{X, Y}\}_{t}$ is a 2 -cocycle on $\operatorname{Vect}\left(S^{1}\right)$ with values in one of the $\operatorname{Vect}\left(S^{1}\right)$-modules $\mathcal{F}_{k}\left(S^{\mathbf{1}}\right)$.

Denote for simplicity $J_{2 k+1}^{1,1}$ by $J_{2 k+1}$.
From the general formula (4) one obtains:
Lemma 5.1. First two terms of $\left\{\widetilde{X, Y\}_{t}}\right.$ are identically zero: $J_{3}(X, Y)=$ $0, J_{5}(X, Y)=0$, the next two terms are proportional to:

$$
\begin{align*}
J_{7}(X, Y)= & X^{\prime \prime \prime} Y^{(I V)}-X^{(I V)} Y^{\prime \prime \prime} \\
J_{9}(X, Y)= & 2\left(X^{\prime \prime \prime} Y^{(V I)}-X^{(V I)} Y^{\prime \prime \prime}\right)  \tag{13}\\
& -9\left(X^{(I V)} Y^{(V)}-X^{(V)} Y^{(I V)}\right) .
\end{align*}
$$

The transvectant $J_{7}$ defines therefore a 2-cocycle. It is a remarkable fact that the same fact is true for $J_{9}$ :

Lemma 5.2. (See [17]). The mappings

$$
J_{7}: \operatorname{Vect}\left(S^{1}\right)^{\otimes 2} \rightarrow \mathcal{F}_{-5} \quad \text { and } \quad J_{9}: \operatorname{Vect}\left(S^{1}\right)^{\otimes 2} \rightarrow \mathcal{F}_{-7}
$$

are 2-cocycles on $\operatorname{Vect}\left(S^{1}\right)$ representing the unique non-trivial classes of the cohomology groups $H^{2}\left(\operatorname{Vect}\left(S^{1}\right) ; \mathcal{F}_{-5}\right)$ and $H^{2}\left(\operatorname{Vect}\left(S^{1}\right) ; \mathcal{F}_{-7}\right)$ respectively.

Proof. Let us prove that $J_{9}$ is a 2 -cocycle on $\operatorname{Vect}\left(S^{1}\right)$. The Jacobi identity for the bracket $\{,\}_{t}$ implies:

$$
\left\{X, J_{9}(Y, Z)\right\}+J_{9}(X,\{Y, Z\})+J_{3}\left(X, J_{7}(Y, Z)\right)+\left(c y c l e_{X, Y, Z}\right)=0
$$

for any $X=p X(q), Y=p Y(q), Z=p Z(q)$. One checks that the expression $J_{3}\left(X, J_{7}(Y, Z)\right)$ is proportional to $X^{\prime \prime \prime}\left(Y^{\prime \prime \prime} Z^{(I V)}-Y^{(I V)} Z^{\prime \prime \prime}\right)$, so that

$$
J_{3}\left(X, J_{7}(Y, Z)\right)+\left(\text { cycle }_{X, Y, Z}\right)=0
$$

We obtain the following relation:

$$
\left\{X, J_{9}(Y, Z)\right\}+J_{9}(X,\{Y, Z\})+\left(\text { cycle }_{X, Y, Z}\right)=0
$$

which means that $J_{9}$ is a 2 -cocycle. Indeed, recall that for any tensor density $a,\left\{p X, p^{m} a\right\}=p^{m} L_{X d / d x}^{(m)}(a)$. Thus, the last relation coincides with the relation $\delta J_{9}=0$.

Let us now show that the cocycle $J_{7}$ on $\operatorname{Vect}\left(S^{\mathbf{1}}\right)$ is not trivial. Consider a linear differential operator $A: \operatorname{Vect}\left(S^{1}\right) \rightarrow \mathcal{F}_{5}$, given by: $A(X(q) d / d q)=\left(\sum_{i=0}^{K} a_{i} X^{(i)}(q)\right)(d q)^{5}$. Then $\delta A(X, Y)=L_{X}^{(5)} A(Y)-$ $L_{Y}^{(5)} A(X)-A([X, Y])$. The higher order part of this expression has a non-zero term $(5-K) a_{K} X^{\prime} Y^{(K)}$ and therefore $J_{7} \neq \delta A$.

In the same way one proves that the cocycle $J_{9}$ on $\operatorname{Vect}\left(S^{1}\right)$ is nontrivial.

Lemma 5.2 is proven.
It follows that the $\star$-product (5) on the algebra $\mathcal{F}$ is not equivalent to the Moyal product.

Theorem 1 is proven.

## 6. Operator representation

We are looking for an linear mapping (depending on $\hbar) F \mapsto \widehat{F}^{\Phi}$ of the associative algebra of Laurent polynomials $\mathcal{F}=\mathbf{C}[p, 1 / p, q]$ into the algebra of formal pseudodifferential operators on $\mathbf{R}$ such that

$$
{\widehat{F \star_{\hbar} G}}{ }^{1}=\widehat{F}^{\Phi} \widehat{G}^{\Phi} .
$$

Recall that the algebra of Laurent polynomials $\mathbf{C}[p, 1 / p, q]$ with the Moyal product is isomorphic to the associative algebra of pseudodifferential operators on $\mathbf{R}$ with polynomial coefficients (see [1]). This isomorphism is defined on the generators $p \mapsto \widehat{p}, q \mapsto \widehat{q}$ by the operators (9) and $p^{-1} \mapsto \widehat{p}^{-1}$ :

$$
\widehat{p}^{-1}=\frac{1}{i \hbar}(\partial / \partial q)^{-1} .
$$

### 6.1 Definition. Put:

$$
\begin{equation*}
\widehat{F}^{\Phi}=\widehat{\Phi^{*} F} . \tag{14}
\end{equation*}
$$

Then $\widehat{F}^{\Phi} \widehat{G}^{\Phi}=\widehat{\Phi^{\star} F} \widehat{\Phi^{*} G}=\Phi^{*} F \star_{\hbar} \Phi^{*} G=\Phi^{*}\left(F \widetilde{\star}_{\hbar} G\right)=\widehat{\mathcal{F}_{\hbar} G}{ }^{\Phi}$.

One obtains the formulæ (10). Indeed,

$$
\widehat{p}^{\Phi}=\widehat{p^{2} / 2}=\frac{(i \hbar)^{2}}{2} \frac{\partial^{2}}{\partial q^{2}} .
$$

Since $q=\frac{1}{2}\left(\left(\frac{1}{p}\right) \widetilde{\star}_{\hbar} p q+p q \widetilde{\star}_{\hbar}\left(\frac{1}{p}\right)\right)$, one gets:

$$
\widehat{q}^{\Phi}=\frac{1}{4 i \hbar}(\Delta \circ A+A \circ \Delta) .
$$

$6.2 s l_{2}$-equivariance. For the Möbius $s l_{2}$ algebra one has:

$$
\begin{aligned}
\widehat{p}^{\Phi} & =\frac{(i \hbar)^{2}}{2} \frac{\partial^{2}}{\partial q^{2}}, \\
\widehat{p q}^{\Phi} & =\frac{i \hbar}{4}+\frac{i \hbar}{2} q \frac{\partial}{\partial q}, \\
{\widehat{p q^{2}}}^{\Phi} & =\frac{q^{2}}{2} .
\end{aligned}
$$

Lemma 6.1. The mapping $F \mapsto \widehat{F}^{\Phi}$ satisfies the Möbius-equivariance condition:

$$
\widehat{\{X, F}\}^{\Phi}=\left[\widehat{X}^{\Phi}, \widehat{F}^{\Phi}\right]
$$

for $X \in s l_{2}$.
Proof. It follows immediately from Theorem 2. Indeed, the $\star$ product (5) is $s l_{2}$-equivariant (that is, satisfying the relation: $\{X, F\}_{t}=$ $\{X, F\}$ for $X \in s l_{2}$ ).

Remark. Beautiful explicit formulæ for $s l_{2}$-equivariant mappings from the space of tensor-densities to the space of pseudodifferential operators are given in [3].

## 7. Hamiltonian isotopy

The simple calculations below are quite standard for the cohomological technique. We need them to prove Corollary 1 of Sec. 3.

Given a syplectomorphism $\Psi$ of a symplectic manifold $V$ and a formal deformation $\{,\}_{t}$ of the Poisson bracket on $V$, we prove that if $\Psi$ is isotopic to the identity, then the formal deformation $\{,\}_{t}^{\Psi}$ defined by:

$$
\{F, G\}_{t}^{\Psi}=\Psi^{*-1}\left\{\Psi^{*} F, \Psi^{*} G\right\}_{t}
$$

is equivalent to $\{,\}_{t}$. The similar proof is valid in the case of $\star$ products.

Recall that two symplectomorphisms $\Psi$ and $\Psi^{\prime}$ of a symplectic manifold $V$ are isotopic if there exists a family of functions $H_{(s)}$ on $V$ such that the symplectomorphism $\Psi_{1} \circ \Psi_{2}^{-1}$ is the flow of the Hamiltionian vector field with the Hamiltonian function $H_{(s)}, 0 \leq s \leq 1$.

Let $\Psi_{(s)}$ be the the flow of a family of functions $H=H_{(s)}$. We will prove that the equivalence class of the formal deformation $\{,\}_{t}^{\Psi_{(s)}}$ does not depend on $s$.
7.1 Equivalence of homotopic cocycles. Let us first show that the cohomology class of the cocycle $C_{3}^{\Psi_{(s)}}$ :

$$
C_{3}^{\Psi_{(s)}}(F, G)=\Psi_{(s)}^{*-1} C_{3}\left(\Psi_{(s)}^{*} F, \Psi_{(s)}^{*} G\right)
$$

does not depend on $s$. To do this, it is sufficient to prove that the derivative $\dot{C}_{3}=\left.\frac{d}{d s} C_{3}^{\Psi(s)}\right|_{s=0}$ is a coboundary. One has

$$
\dot{C}_{3}(F, G)=C_{3}(\{H, F\}, G)+C_{3}(F,\{H, G\})-\left\{H, C_{3}(F, G)\right\} .
$$

The relation $\delta C_{3}=0$ implies:

$$
\dot{C}_{3}(F, G)=\left\{F, C_{3}(G, H)\right\}-\left\{G, C_{3}(F, H)\right\}-C_{3}(\{F, G\}, H) .
$$

This means, $\left.\frac{d}{d s} C_{3}^{\Psi(s)}\right|_{s=0}=\delta B_{H}$, where $B_{H}(F)=C_{3}(F, H)$.
7.2 General case. Let us apply the same arguments to prove that the deformations $\{,\}_{t}^{\Psi_{(s)}}$ are equivalent to each other for all values of $s$. For this purpose we must show that there exists a family of mappings $A_{(s)}(F)=F+\sum_{k=1}^{\infty} A_{(s)_{k}}(F) t^{k}$ such that $A_{(s)}^{-1}\left(\left\{A_{(s)}(F), A_{(s)}(G)\right\}_{t}\right)=$ $\{F, G\}_{t}$.

It is sufficient to verify the existence of a mapping $a(F)=\sum_{k=1}^{\infty} a_{k}(F) t^{k}$ (the derivative: $\left.a(F)=d /\left.d s\left(A_{(s)}(F)\right)\right|_{s=s_{0}}\right)$ such that

$$
\left.\frac{d}{d s}\{F, G\}_{t}^{\Psi_{(s)}}\right|_{s=s_{0}}=\{a(F), G\}_{t}+\{F, a(G)\}-a\left(\{F, G\}_{t}\right.
$$

Since

$$
\left.\frac{d}{d s}\{F, G\}_{t}^{\Psi_{(s)}}\right|_{s=s_{0}}=\{\{F, H\}, G\}_{t}+\{F,\{G, H\}\}_{t}-\left\{\{F, G\}_{t}, H\right\},
$$

from the Jacobi identity:

$$
\left\{\{F, H\}_{t}, G\right\}_{t}+\left\{F,\{G, H\}_{t}\right\}_{t}-\left\{\{F, G\}_{t}, H\right\}_{t}=0
$$

one obtains that the mapping $a(F)$ can be written in the form:

$$
a(F)=\sum_{k=1}^{\infty} \frac{1}{(2 k+1)!} C_{2 k+1}\left(F, H_{\left(s_{0}\right)}\right) t^{k} .
$$

7.3 Proof of Corollary 2. Consider the $\star$-product (8) given by $F \star_{\hbar}^{\Phi} G$, where $F \star_{\hbar} G$ is the Moyal product (1), and $\Phi:(p, q) \mapsto$ ( $p^{2} / 2, q / p$ ). It is defined on $\mathbf{R}^{2} \backslash \mathbf{R}$.

The t-product (8) on the algebra $C^{\infty}\left(\mathbf{R}^{2} \backslash \mathbf{R}\right)$ is equivalent to the Moyal product. Indeed, the symplectomorphism $\Phi$ is isotopic to the identity in the group of all smooth symplectomorphisms of $\mathbf{R}^{2} \backslash \mathbf{R}$. The isotopy is: $\Phi_{s}:(p, q) \mapsto\left(\frac{p^{1+s}}{1+s}, \frac{q}{p^{s}}\right)$, where $s \in[0,1]$.

Recall that the $\star$-product $F \star{ }_{\hbar}^{\Phi} G$ on the algebra $\mathcal{F}$ is not equivalent to the Moyal product since it coincides with the product (5).

The family $\Phi_{s}$ does not preserve the algebra $\mathcal{F}$. Theorem 1 implies that $F$ is not isotopic to the identity in the group of symplectomorphisms of $\mathbf{R}^{2} \backslash \mathbf{R}$ preserving the algebra $\mathcal{F}$.

## 8. Discussion

### 8.1 Difficulties in multi-dimensional case.

There exist multi-dimensional analogues of transvectants [14] and [16].

Consider the projective space $\mathbf{R} \mathbf{P}^{2 n+1}$ endowed with the standard contact structure (or an open domain of the complex projective space $\mathbf{C} \mathbf{P}^{2 n+1}$ ). There exists an unique bilinear differential operator of order $k$ on tensor-densities equivariant with respect to the action of the group $S p_{2 n}$ (see [14], [16]):

$$
\begin{equation*}
J_{k}^{\lambda, \mu}: \mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \rightarrow \mathcal{F}_{\lambda+\mu-\frac{k}{n+1}}, \tag{15}
\end{equation*}
$$

where $\mathcal{F}_{\lambda}=\mathcal{F}_{\lambda}\left(\mathbf{P}^{2 n+1}\right)$ is the space of tensor-densities on $\mathbf{P}^{2 n+1}$ of degree $-\lambda$ :

$$
f=f\left(x_{1}, \ldots, x_{2 n+1}\right)\left(d x_{1} \wedge \ldots d x_{2 n+1}\right)^{-\lambda}
$$

The space of tensor-densities $\mathcal{F}_{\lambda}\left(\mathbf{R P}^{2 n+1}\right)$ is isomorphic as a module over the group of contact diffeomorphisms to the space of homogeneous functions on $\mathbf{R}^{2 n+2}$, and the isomorphism is given by:

$$
f \mapsto F\left(y_{1}, \ldots, y_{2 n+2}\right)=y_{2 n+2}^{-\lambda(n+1)} f\left(\frac{y_{1}}{y_{2 n+2}}, \ldots, \frac{y_{2 n+1}}{y_{2 n+2}}\right) .
$$

Then the operations (15) are defined as the restrictions of the terms of the standard $\star$-product on $\mathbf{R}^{2 n+2}$.

The same formula (5) defines a $\star$-product on the space of tensordensities on $\mathbf{C P}^{2 n+1}$ ) (cf. [16]). However, there is no analogues of the symplectomorphism (7). I do not know if there exists a $\star$-product on the Poisson algebra $\mathbf{C}\left[y_{2 n+2}, y_{2 n+2}^{-1}\right] \otimes C^{\infty}\left(\mathbf{R} \mathbf{P}^{2 n+1}\right)$ non-equivalent to the standard.
8.2 Classification problem. The classification (modulo equivalence) of $\star$-products on the Poisson algebra $\mathcal{F}$ is an interesting open problem. It is related to the calculation of cohomology groups $H^{2}(\mathcal{F} ; \mathcal{F})$ and $H^{3}(\mathcal{F} ; \mathcal{F})$. The following result was announced in [7]: $\operatorname{dim} H^{2}(\mathcal{F} ; \mathcal{F})=$ 2.

Let us formulate a conjecture in the compact case. Consider the Poisson algebra $\mathcal{F}\left(S^{1}\right)$ of functions on $T^{*} S^{1} \backslash S^{1}$ which are Laurent polynomials on the fiber: $F(p, q)=\sum_{-N \leq i \leq N} p^{i} f_{i}(q)$ where $f_{i}(q+1)=$ $f_{i}(q)$.

Conjecture. Every $\star$-product on $\mathcal{F}\left(S^{1}\right)$ is equivalent to (1) or (5).

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Centre National De La Recherche Scientifique, Marseille


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