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EXOTIC DEFORMATION QUANTIZATION

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1. Introduction

Let \mathcal{A} be one of the following commutative associative algebras: the algebra of all smooth functions on the plane: $\mathcal{A} = C^{\infty}(\mathbf{R}^2)$, or the algebra of polynomials $\mathcal{A} = \mathbf{C}[p,q]$ over \mathbf{R} or \mathbf{C} . There exists a non-trivial formal *associative* deformation of \mathcal{A} called the *Moyal* \star -product (or the standard \star -product). It is defined as an associative operation $\mathcal{A}^{\otimes 2} \to \mathcal{A}[[\hbar]]$ where \hbar is a formal variable. The explicit formula is:

(1)
$$F \star_{\hbar} G = FG + \sum_{k \ge 1} \frac{(i\hbar)^k}{2^k k!} \{F, G\}_k,$$

where $\{F, G\}_1 = \frac{\partial F}{\partial p} \frac{\partial G}{\partial q} - \frac{\partial F}{\partial q} \frac{\partial G}{\partial p}$ is the standard Poisson bracket, and the higher order terms are:

(2)
$$\{F,G\}_k = \sum_{i=0}^k (-1)^i \begin{pmatrix} k \\ i \end{pmatrix} \frac{\partial^k F}{\partial p^{k-i} \partial q^i} \frac{\partial^k G}{\partial p^i \partial q^{k-i}}.$$

The Moyal product is the *unique* (modulo equivalence) non-trivial formal deformation of the associative algebra \mathcal{A} (see [13]).

Definition 1. A formal associative deformation of \mathcal{A} given by formula (1) is called a \star -product if the following hold:

1) the first order term coincides with the Poisson bracket: $\{F, G\}_1 = \{F, G\};$

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2) the higher order terms $\{F, G\}_k$ are given by differential operators vanishing on constants: $\{1, G\}_k = \{F, 1\}_k = 0$;

3) $\{F,G\}_k = (-1)^k \{G,F\}_k$.

Definition 2. Two \star -products \star_{\hbar} and \star'_{\hbar} on \mathcal{A} are called *equivalent* if there exists a linear mapping $A_{\hbar} : \mathcal{A} \to \mathcal{A}[[\hbar]]$ such that

$$A_{\hbar}(F) = F + \sum_{k=1}^{\infty} A_k(F)\hbar^k$$

intertwining the operations \star_{\hbar} and \star'_{\hbar} : $A_{\hbar}(F) \star'_{\hbar} A_{\hbar}(G) = A_{\hbar}(F \star_{\hbar} G)$.

Consider now \mathcal{A} as a Lie algebra; the commutator is given by the Poisson bracket. The Lie algebra \mathcal{A} has a unique (modulo equivalence) non-trivial formal deformation called the *Moyal bracket* or the Moyal \star -commutator: $\{F, G\}_t = \frac{1}{i\hbar}(F \star_{\hbar} G - G \star_{\hbar} F)$, where $t = -\hbar^2/2$.

The well-known De Wilde-Lecomte theorem [4] states the existence of a non-trivial \star -product for an arbitrary symplectic manifold. The theory of \star -products is a subject of *deformation quantization*. The geometrical proof of the existence theorem was given by B. Fedosov [9] (see [8] and [20] for clear explanation and survey of recent progress).

The main idea of this paper is to consider the algebra $\mathcal{F}(M)$ of functions (with singularities) on the cotangent bundle T^*M which are *Laurent polynomials* on the fibers. In contrast to the above algebra \mathcal{A} it turns out that for such algebras the standard \star -product is no more unique at least if M is one-dimensional: dim M = 1.

We consider $M = S^1$, **R** in the real case, and $M = \mathcal{H}$ (the upper half-plane) in the holomorphic case. The main result of this paper is an explicit construction of a new \star -product on the algebra $\mathcal{F}(M)$ nonequivalent to the standard Moyal product. This \star -product is equivariant with respect to the Möbius transformations. The construction is based on the bilinear SL_2 -equivariant operations on tensor-densities on M, known as Gordan transvectants and Rankin-Cohen brackets.

We study the relations between the new \star -product and extensions of the Lie algebra $\operatorname{Vect}(S^1)$.

The results of this paper are closely related to those of the recent work of P. Cohen, Yu. Manin and D. Zagier [3] where a one-parameter family of associative products on the space of classical modular forms is constructed using the same SL_2 -equivariant bilinear operations.

2. Definition of the exotic *-product

2.1 Algebras of Laurent polynomials. Let \mathcal{F} be one of the following associative algebras of functions:

$$\mathcal{F} = \mathbf{C}[p, 1/p] \otimes C^{\infty}(\mathbf{R})$$
 or $\mathcal{F} = \mathbf{C}[p, 1/p] \otimes \operatorname{Hol}(\mathcal{H})$

This means, it consists of functions of the type:

(3)
$$F(p,q) = \sum_{i=-N}^{N} p^{i} f_{i}(q),$$

where $f_i(q) \in C^{\infty}(\mathbf{R})$ in the real case, or $f_i(q)$ are holomorphic functions on the upper half-plane \mathcal{H} or $f_i \in C[q]$ (respectively).

We will also consider the algebra of polynomials: $\mathbf{C}[p, 1/p, q]$ (Laurent polynomials in p).

2.2 Transvectants. Consider the following bilinear operators on functions of one variable:

(4)
$$J_k^{m,n}(f,g) = \sum_{i+j=k} (-1)^i \binom{k}{i} \frac{(2m-i)!(2n-j)!}{(2m-k)!(2n-k)!} f^{(i)}g^{(j)},$$

where $f = f(z), g = g(z), f^{(i)}(z) = \frac{d^i f(z)}{dz^i}$. These operators satisfy a remarkable property: they are equivariant

These operators satisfy a remarkable property: they are equivariant under Möbius (linear-fractional) transformations. Namely, suppose that the transformation $z \mapsto \frac{az+b}{cz+d}$ (with ad - bc = 1) acts on the arguments as follows:

$$f(z) \mapsto f\left(\frac{az+b}{cz+d}\right)(cz+d)^{2m}, \quad g(z) \mapsto g\left(\frac{az+b}{cz+d}\right)(cz+d)^{2n},$$

then $J_k^{m,n}(f,g)$ transforms as:

$$J^{m,n}_k(f,g)(z)\mapsto J^{m,n}_k(f,g)(\frac{az+b}{cz+d})(cz+d)^{2(m+n-k)}.$$

In other words, the operations (4) are bilinear SL_2 -equivariant mappings on tensor-densities:

$$J_k^{mn}:\mathcal{F}_m\otimes\mathcal{F}_n\to\mathcal{F}_{m+n-k}$$

where \mathcal{F}_l is the space of tensor-densities of degree -l: $\phi = \phi(z)(dz)^{-l}$.

The operations (4) were discovered more than one hundred years ago by Gordan [11] who called them the *transvectants*. They have been rediscovered many times: in the theory of modular functions by Rankin [18] and by Cohen [2] (so-called Rankin-Cohen brackets), in differential projective geometry by Janson and Peetre [12]. The "multi-dimensional transvectants" were defined in [14] in the context of the the Virasoro algebra and symplectic and contact geometry.

2.3 Main definition. Define the following bilinear mapping $\mathcal{F}^{\otimes 2} \to \mathcal{F}[[\hbar]]$, for $F = p^m f(q), G = p^n g(q)$, where $m, n \in \mathbb{Z}$, by putting:

(5)
$$F\widetilde{\star}_{\hbar}G = \sum_{k=0}^{\infty} \frac{(i\hbar)^k}{2^{2k}} p^{(m+n-k)} J_k^{m,n}(f,g),$$

Note, that the first order term coincides with the Poisson bracket.

This operation will be the main subject of this paper. We call it the *exotic* \star -*product*.

2.4 Remark. Another one-parameter family of operations on modular forms: $f \star^{\kappa} g = \sum_{n=0}^{\infty} t_n^{\kappa}(k,l) J_n^{kl}(f,g)$, where f and g are modular forms of weight k and l respectively, and $t_n^{\kappa}(k,l)$ are very interesting and complicated coefficients, is defined in [3].

3. Main theorems

We formulate here the main results of this paper. All the proofs will be given in Sections 4-7.

3.1 Non-equivalence. The Moyal \star -product (1) defines a non-trivial formal deformation of \mathcal{F} . We will show that the formula (5) defines a \star -product non-equivalent to the standard Moyal product.

Theorem 1. The operation (5) is associative; it defines a formal deformation of the algebra \mathcal{F} which is not equivalent to the Moyal product.

The associativity of the product (5) is a trivial corollary of Proposition 1 below. To prove the non-equivalence, we will use the relations with extensions of the Lie algebra of vector fields on S^1 : Vect $(S^1) \subset \mathcal{F}$ (cf. Sec.5). It is interesting to note that the constructed \star -product is equivalent to the standard Moyal product if we consider it on the algebra $C^{\infty}(T^*M \setminus M)$ of all smooth functions (not only Laurent polynomials on fibers); cf. Corollary 1 below.

3.2 sl_2 -equivariance. The Lie algebra $sl_2(\mathbf{R})$ has two natural embeddings into the Poisson Lie algebra on \mathbf{R}^2 : the symplectic Lie algebra $sp_2(\mathbf{R}) \cong sl_2(\mathbf{R})$ generated by quadratic polynomials (p^2, pq, q^2) and another one with generators: (p, pq, pq^2) which is called the *Möbius algebra*.

It is well-known that the Moyal product (1) is the unique non-trivial formal deformation of the associative algebra of functions on \mathbf{R}^2 equivariant under the action of the symplectic algebra. This means, (1) satisfies the Leibnitz property:

(6)
$$\{F, G \star_{\hbar} H\} = \{F, G\} \star_{\hbar} H + G \star_{\hbar} \{F, H\}$$

where F is a quadratic polynomial (note that $\{F, G\}_t = \{F, G\}$ if F is a quadratic polynomial).

Theorem 2. The product (5) is the unique formal deformation of the associative algebra \mathcal{F} equivariant under the action of the Möbius algebra.

The product (5) is the unique non-trivial formal deformation of \mathcal{F} satisfying (6) for F from the Möbius sl_2 algebra.

3.3 Symplectomorphism Φ . The relation between the Moyal product and the product (5) is as follows. Consider the symplectic mapping

(7)
$$\Phi(p,q) = \left(\frac{p^2}{2}, \frac{q}{p}\right)$$

defined on $\mathbf{R}^2 \setminus \mathbf{R}$ in the real case and on \mathcal{H} in the complex case.

Proposition 1. The product (5) is the Φ -conjugation of the Moyal product:

(8)
$$F \check{\star}_{\hbar} G = F \star_{\hbar}^{\Phi} G := (F \circ \Phi \star_{\hbar} G \circ \Phi) \circ \Phi^{-1}.$$

Remark. The mapping (7) (in the complex case) can be interpreted as follows. It transforms the space of *holomorphic tensor-densities* of

degree -k on \mathbb{CP}^1 to the space $\mathbb{C}^k[p,q]$ of polynomials of degree k. Indeed, there exists a natural isomorphism $z^n(dz)^{-m} \mapsto p^m q^n$ (where $m \geq 2n$) and $(p^m q^n) \circ \Phi = p^{2m-n}q^n$.

3.4 Operator formalism. The Moyal product is related to the following Weil quantization procedure. Define the following differential operators:

(9)
$$\widehat{p} = i\hbar\frac{\partial}{\partial q},$$

$$\widehat{q} = q$$

satisfying the canonical relation: $[\hat{p}, \hat{q}] = i\hbar \mathbf{I}$. Associate to each polynomial F = F(p, q) the differential operator $\hat{F} = \operatorname{Sym} F(\hat{p}, \hat{q})$ symmetric in \hat{p} and \hat{q} . The Moyal product on the algebra of polynomials coincides with the product of differential operators: $\widehat{F \star_h G} = \widehat{FG}$.

We will show that the \star -product (5) leads to the operators:

(10)
$$\widehat{p}^{\Phi} = (\frac{i\hbar}{2})^2 \Delta, \widehat{q}^{\Phi} = \frac{1}{4i\hbar} (\Delta^{-1} \circ A + A \circ \Delta^{-1})$$

(where $\Delta = \frac{\partial^2}{\partial q^2}$ and $A = 2q\frac{\partial}{\partial q} + 1$ is the dilation operator) also satisfying the canonical relation.

Remark that \hat{p}^{Φ} and \hat{q}^{Φ} given by (10) on the Hilbert space $L_2(\mathbf{R})$ are not equivalent to the operators (9) since \hat{q}^{Φ} is symmetric but not self-adjoint (see [6] on this subject).

3.5 "Symplectomorphic" deformations. Let us consider the general situation.

Proposition 2. Given a symplectic manifold V endowed with a \star -product \star_{\hbar} and a symplectomorphism Ψ of V, if there exists a hamiltonian isotopy of Ψ to the identity, then the Ψ -conjugate product \star_{\hbar}^{Ψ} defined according to the formula (8) is equivalent to \star_{\hbar} .

Corollary 1. The \star -product (5) considered on the algebra of all smooth functions $C^{\infty}(T^*\mathbf{R} \setminus \mathbf{R})$ is equivalent to the Moyal product.

4. Möbius-invariance

In this section we prove Theorem 2. We show that the operations of transvectant (4) are Φ -conjugate of the terms of the Moyal product.

4.1 Lie algebra $Vect(\mathbf{R})$ and modules of tensor-densities. Let $Vect(\mathbf{R})$ be the Lie algebra of smooth (or polynomial) vector fields on \mathbf{R} :

$$X = X(x)\frac{d}{dx}$$

with the commutator

$$[X(x)\frac{d}{dx}, Y(x)\frac{d}{dx}] = (X(x)Y'(x) - X'(x)Y(x))\frac{d}{dx}.$$

The natural embedding of the Lie algebra $sl_2 \subset \text{Vect}(\mathbf{R})$ is generated by the vector fields $d/dx, xd/dx, x^2d/dx$.

Define a 1-parameter family of Vect(**R**)-actions on $C^{\infty}(\mathbf{R})$ given by

(11)
$$L_X^{(\lambda)}f = X(x)f'(x) - \lambda X'(x)f(x),$$

where $\lambda \in \mathbf{R}$. Geometrically, $L_X^{(\lambda)}$ is the operator of Lie derivative on *tensor-densities* of degree $-\lambda$:

$$f = f(x)(dx)^{-\lambda}$$

Denote \mathcal{F}_{λ} the Vect(**R**)-module structure on $C^{\infty}(\mathbf{R})$ given by (11).

4.2 Transvectant as a bilinear sl_2 -equivariant operator. The operations (4) can be defined as bilinear mappings on $C^{\infty}(\mathbf{R})$ which are sl_2 -equivariant:

Statement 4.1. For each k = 0, 1, 2, ... there exists a unique (up to a constant) bilinear sl_2 -equivariant mapping

$$\mathcal{F}_{\mu} \otimes \mathcal{F}_{\nu} \to \mathcal{F}_{\mu+\nu-k}.$$

It is given by $f \otimes g \mapsto J_k^{\mu,\nu}(f,g)$.

Proof. Straightforward (cf. [11], [12]).

4.3 Algebra \mathcal{F} as a module over $\operatorname{Vect}(\mathbf{R})$. The Lie algebra $\operatorname{Vect}(\mathbf{R})$ can be considered as a Lie subalgebra of \mathcal{F} . The embedding $\operatorname{Vect}(\mathbf{R}) \subset \mathcal{F}$ is given by:

$$X(x)\frac{d}{dx} \mapsto pX(q).$$

The algebra \mathcal{F} is therefore, a Vect(**R**)-module.

Lemma 4.2. The algebra \mathcal{F} is decomposed to a direct sum of Vect(**R**)-modules:

$$\mathcal{F} = \oplus_{m \in \mathbf{Z}} \mathcal{F}_m.$$

Proof. Consider the subspace of \mathcal{F} consisting of functions homogeneous of degree m in p: $F = p^m f(q)$. This subspace is a Vect(**R**)-module isomorphic to \mathcal{F}_m . Indeed, $\{pX(q), p^m f(q)\} = p^m (Xf' - mX'f) = p^m L_X^{(m)} f$.

4.4 Projective property of the diffeomorphism Φ . The transvectants (4) coincide with the Φ -conjugate operators (2) from the Moyal product:

Proposition 4.3. Let
$$F = p^m f(q), G = p^n g(q)$$
. Then
(12) $\Phi^{*-1} \{ \Phi^* F, \Phi^* G \}_k = \frac{k!}{2^k} p^{m+n-k} J_k^{m,n}(f,g).$

Proof. The symplectomorphism Φ of \mathbf{R}^2 intertwines the symplectic algebra $sp_2 \equiv sl_2$ and the Möbius algebra: $\Phi^*(p, pq, pq^2) = (\frac{1}{2}p^2, \frac{1}{2}pq, \frac{1}{2}q^2)$. Therefore, the operation $\Phi^{*-1}{\Phi^*F, \Phi^*G}_k$ is Möbius-equivariant.

On the other hand, one has: $\Phi^*F = \frac{1}{2^m}p^{2m}f(\frac{q}{p})$ and $\Phi^*G = \frac{1}{2^n}p^{2m}g(\frac{q}{p})$. Since Φ^*F and Φ^*G are homogeneous of degree 2m and 2n (respectively), the function $\{\Phi^*F, \Phi^*G\}_k$ is also homogeneous of degree 2(m + n - k). Thus, the operation $\{F, G\}_k^{\Phi} = \Phi^{*-1}\{\Phi^*F, \Phi^*G\}_k$ defines a bilinear mapping on the space of tensor-densities $\mathcal{F}_m \otimes \mathcal{F}_n \to \mathcal{F}_{m+n-k}$ which is sl_2 -equivariant.

Statement 4.1 implies that it is proportional to $J_k^{m,n}$. One easily verifies the coefficient of proportionality for $F = p^m, G = p^n q^k$, to obtain the formula (12).

Proposition 4.3 is proven.

Remark. Proposition 4.3 was proven in [15]. We do not know whether this elementary fact has been mentioned by classics.

4.5 Proof of Theorem 2. Proposition 4.3 implies that the formula (5) is a Φ -conjugation of the Moyal product and is given by the formula (8).

Proposition 1 is proven.

It follows that (5) is a \star -product on \mathcal{F} equivariant under the action of the Möbius sl_2 algebra. Moreover, it is the unique \star -product with this property since the Moyal product is the unique \star -product equivariant under the action of the symplectic algebra.

Theorem 2 is proven.

5. Relation with extensions of the Lie algebra $Vect(S^1)$

We prove here that the \star -product (5) is not equivalent to the Moyal product.

Let $\operatorname{Vect}(S^1)$ be the Lie algebra of vector fields on the circle. Consider the embedding $\operatorname{Vect}(S^1) \subset \mathcal{F}$ given by functions on \mathbb{R}^2 of the type: X = pX(q) where X(q) is periodical: X(q+1) = X(q).

5.1 An idea of the proof of Theorem 1. Consider the formal deformations of the *Lie algebra* \mathcal{F} associated to the \star -products (1) and (5). The restriction of the Moyal bracket to $\operatorname{Vect}(S^1)$ is identically zero. We show that the restriction of the \star -commutator

$$\widetilde{\{F,G\}}_t = \frac{1}{i\hbar} (F \widetilde{\star}_\hbar G - G \widetilde{\star}_\hbar F), \quad t = -\frac{\hbar^2}{2}$$

associated to the *-product (5) defines a series of non-trivial extensions of the Lie algebra $\operatorname{Vect}(S^1)$ by the modules $\mathcal{F}_k(S^1)$ of tensor-densities on S^1 of degree -k.

5.2 Extensions and the cohomology group $H^2(\operatorname{Vect}(S^1); \mathcal{F}_{\lambda})$. Recall that an extension of a Lie algebra by its module is defined by a 2-cocycle on it with values in this module. To define an extension of $\operatorname{Vect}(S^1)$ by the module \mathcal{F}_{λ} one needs therefore a bilinear mapping $c: \operatorname{Vect}(S^1)^{\otimes 2} \to \mathcal{F}_{\lambda}$ which satisfies the identity $\delta c = 0$:

$$c(X, [Y, Z]) + L_X^{(\lambda)} c(Y, Z) + (cycle_{X,Y,Z}) = 0.$$

(See [10]).

The cohomology group $H^2(\operatorname{Vect}(S^1); \mathcal{F}_{\lambda})$ were calculated in [19] (see [10]). This group is trivial for each value of λ except $\lambda = 0, -1, -2, -5, -7$. The explicit formulæ for the corresponding non-trivial cocycles are given in [17]. If $\lambda = -5, -7$, then dim $H^2(\operatorname{Vect}(S^1); \mathcal{F}_{\lambda}) = 1$, the cohomology group is generated by the unique (up to equivalence) non-trivial cocycle. We will obtain these cocycles from the \star -commutator.

5.3 Non-trivial cocycles on $Vect(S^1)$.

Consider the restriction of the *-commutator $\{,\}_t$ (corresponding to the *-product (5)) to Vect $(S^1) \subset \mathcal{F}$: let

$$X = pX(q), \ Y = pY(q);$$

then from (5) we have

$${X,Y}_t = {X,Y} + \sum_{k=1}^{\infty} \frac{t^k}{2^{2k+1}} \frac{1}{p^{2k-1}} J_{2k+1}^{1,1}(X,Y).$$

It follows from the Jacobi identity that the first non-zero term of the series $\{X, Y\}_t$ is a 2-cocycle on $\operatorname{Vect}(S^1)$ with values in one of the $\operatorname{Vect}(S^1)$ -modules $\mathcal{F}_k(S^1)$.

Denote for simplicity $J_{2k+1}^{1,1}$ by J_{2k+1} .

From the general formula (4) one obtains:

Lemma 5.1. First two terms of $\{X, Y\}_t$ are identically zero: $J_3(X, Y) = 0$, $J_5(X, Y) = 0$, the next two terms are proportional to:

(13)
$$J_{7}(X,Y) = X'''Y^{(IV)} - X^{(IV)}Y''', J_{9}(X,Y) = 2(X'''Y^{(VI)} - X^{(VI)}Y''') -9(X^{(IV)}Y^{(V)} - X^{(V)}Y^{(IV)}).$$

The transvectant J_7 defines therefore a 2-cocycle. It is a remarkable fact that the same fact is true for J_9 :

Lemma 5.2. (See [17]). The mappings

$$J_7: \operatorname{Vect}(S^1)^{\otimes 2} \to \mathcal{F}_{-5} \quad and \quad J_9: \operatorname{Vect}(S^1)^{\otimes 2} \to \mathcal{F}_{-7}$$

are 2-cocycles on $\operatorname{Vect}(S^1)$ representing the unique non-trivial classes of the cohomology groups $H^2(\operatorname{Vect}(S^1); \mathcal{F}_{-5})$ and $H^2(\operatorname{Vect}(S^1); \mathcal{F}_{-7})$ respectively.

Proof. Let us prove that J_9 is a 2-cocycle on Vect (S^1) . The Jacobi identity for the bracket $\{, \}_t$ implies:

$$\{X, J_9(Y, Z)\} + J_9(X, \{Y, Z\}) + J_3(X, J_7(Y, Z)) + (cycle_{X, Y, Z}) = 0$$

for any X = pX(q), Y = pY(q), Z = pZ(q). One checks that the expression $J_3(X, J_7(Y, Z))$ is proportional to $X'''(Y''Z^{(IV)} - Y^{(IV)}Z''')$, so that

$$J_3(X, J_7(Y, Z)) + (cycle_{X,Y,Z}) = 0.$$

We obtain the following relation:

$$\{X, J_9(Y, Z)\} + J_9(X, \{Y, Z\}) + (cycle_{X,Y,Z}) = 0,$$

which means that J_9 is a 2-cocycle. Indeed, recall that for any tensor density a, $\{pX, p^m a\} = p^m L_{Xd/dx}^{(m)}(a)$. Thus, the last relation coincides with the relation $\delta J_9 = 0$.

Let us now show that the cocycle J_7 on $\operatorname{Vect}(S^1)$ is not trivial. Consider a linear differential operator $A : \operatorname{Vect}(S^1) \to \mathcal{F}_5$, given by: $A(X(q)d/dq) = (\sum_{i=0}^{K} a_i X^{(i)}(q))(dq)^5$. Then $\delta A(X,Y) = L_X^{(5)}A(Y) - L_Y^{(5)}A(X) - A([X,Y])$. The higher order part of this expression has a non-zero term $(5 - K)a_K X'Y^{(K)}$ and therefore $J_7 \neq \delta A$.

In the same way one proves that the cocycle J_9 on $\operatorname{Vect}(S^1)$ is non-trivial.

Lemma 5.2 is proven.

It follows that the \star -product (5) on the algebra \mathcal{F} is not equivalent to the Moyal product.

Theorem 1 is proven.

6. Operator representation

We are looking for an linear mapping (depending on \hbar) $F \mapsto \hat{F}^{\Phi}$ of the associative algebra of Laurent polynomials $\mathcal{F} = \mathbf{C}[p, 1/p, q]$ into the algebra of formal pseudodifferential operators on \mathbf{R} such that

$$\widehat{F\check{\star}_{\hbar}G}^{\Phi} = \widehat{F}^{\Phi}\widehat{G}^{\Phi}.$$

Recall that the algebra of Laurent polynomials $\mathbf{C}[p, 1/p, q]$ with the Moyal product is isomorphic to the associative algebra of pseudodifferential operators on \mathbf{R} with polynomial coefficients (see [1]). This isomorphism is defined on the generators $p \mapsto \hat{p}, q \mapsto \hat{q}$ by the operators (9) and $p^{-1} \mapsto \hat{p}^{-1}$:

$$\widehat{p}^{-1} = \frac{1}{i\hbar} (\partial/\partial q)^{-1}.$$

6.1 Definition. Put:

(14)
$$\widehat{F}^{\Phi} = \widehat{\Phi^* F}.$$

Then $\widehat{F}^{\Phi}\widehat{G}^{\Phi} = \widehat{\Phi^*F}\widehat{\Phi^*G} = \Phi^*F \star_{\hbar} \Phi^*G = \Phi^*(F\widetilde{\star}_{\hbar}G) = \widehat{F\widetilde{\star}_{\hbar}G}^{\Phi}.$

One obtains the formulæ (10). Indeed,

$$\widehat{p}^{\Phi} = \widehat{p^2/2} = \frac{(i\hbar)^2}{2} \frac{\partial^2}{\partial q^2}.$$

Since $q = \frac{1}{2} \left(\left(\frac{1}{p} \right) \breve{\star}_{\hbar} p q + p q \breve{\star}_{\hbar} \left(\frac{1}{p} \right) \right)$, one gets:

$$\widehat{q}^{\Phi} = \frac{1}{4i\hbar} (\Delta \circ A + A \circ \Delta).$$

6.2 sl_2 -equivariance. For the Möbius sl_2 algebra one has:

Lemma 6.1. The mapping $F \mapsto \widehat{F}^{\Phi}$ satisfies the Möbius-equivariance condition:

$$\widehat{\{X,F\}}^{\Phi} = [\widehat{X}^{\Phi}, \widehat{F}^{\Phi}]$$

for $X \in sl_2$.

Proof. It follows immediately from Theorem 2. Indeed, the \star -product (5) is sl_2 -equivariant (that is, satisfying the relation: $\{X, F\}_t = \{X, F\}$ for $X \in sl_2$).

Remark. Beautiful explicit formulæ for sl_2 -equivariant mappings from the space of tensor-densities to the space of pseudodifferential operators are given in [3].

7. Hamiltonian isotopy

The simple calculations below are quite standard for the cohomological technique. We need them to prove Corollary 1 of Sec. 3.

Given a syplectomorphism Ψ of a symplectic manifold V and a formal deformation $\{, \}_t$ of the Poisson bracket on V, we prove that if Ψ is isotopic to the identity, then the formal deformation $\{, \}_t^{\Psi}$ defined by:

$$\{F,G\}_t^{\Psi} = \Psi^{*-1}\{\Psi^*F,\Psi^*G\}_t$$

is equivalent to $\{ , \}_t$. The similar proof is valid in the case of \star -products.

Recall that two symplectomorphisms Ψ and Ψ' of a symplectic manifold V are *isotopic* if there exists a family of functions $H_{(s)}$ on V such that the symplectomorphism $\Psi_1 \circ \Psi_2^{-1}$ is the flow of the Hamiltionian vector field with the Hamiltonian function $H_{(s)}$, $0 \le s \le 1$.

Let $\Psi_{(s)}$ be the flow of a family of functions $H = H_{(s)}$. We will prove that the equivalence class of the formal deformation $\{ , \}_t^{\Psi_{(s)}}$ does not depend on s.

7.1 Equivalence of homotopic cocycles. Let us first show that the cohomology class of the cocycle $C_3^{\Psi(s)}$:

$$C_3^{\Psi_{(s)}}(F,G) = \Psi_{(s)}^{*-1} C_3(\Psi_{(s)}^*F, \Psi_{(s)}^*G)$$

does not depend on s. To do this, it is sufficient to prove that the derivative $\dot{C}_3 = \frac{d}{ds} C_3^{\Psi(s)}|_{s=0}$ is a coboundary. One has

$$\dot{C}_3(F,G) = C_3(\{H,F\},G) + C_3(F,\{H,G\}) - \{H,C_3(F,G)\}.$$

The relation $\delta C_3 = 0$ implies:

$$\dot{C}_3(F,G) = \{F, C_3(G,H)\} - \{G, C_3(F,H)\} - C_3(\{F,G\},H).$$

This means, $\frac{d}{ds}C_3^{\Psi(s)}|_{s=0} = \delta B_H$, where $B_H(F) = C_3(F, H)$.

7.2 General case. Let us apply the same arguments to prove that the deformations $\{ , \}_t^{\Psi(s)}$ are equivalent to each other for all values of s. For this purpose we must show that there exists a family of mappings $A_{(s)}(F) = F + \sum_{k=1}^{\infty} A_{(s)_k}(F)t^k$ such that $A_{(s)}^{-1}(\{A_{(s)}(F), A_{(s)}(G)\}_t) = \{F, G\}_t$.

It is sufficient to verify the existence of a mapping $a(F) = \sum_{k=1}^{\infty} a_k(F) t^k$ (the derivative: $a(F) = d/ds(A_{(s)}(F))|_{s=s_0}$) such that

$$\frac{d}{ds} \{F, G\}_t^{\Psi(s)}|_{s=s_0} = \{a(F), G\}_t + \{F, a(G)\} - a(\{F, G\}_t)\}$$

Since

$$\frac{d}{ds}\{F,G\}_t^{\Psi(s)}|_{s=s_0} = \{\{F,H\},G\}_t + \{F,\{G,H\}\}_t - \{\{F,G\}_t,H\},\$$

from the Jacobi identity:

$$\{\{F,H\}_t,G\}_t + \{F,\{G,H\}_t\}_t - \{\{F,G\}_t,H\}_t = 0$$

one obtains that the mapping a(F) can be written in the form:

$$a(F) = \sum_{k=1}^{\infty} \frac{1}{(2k+1)!} C_{2k+1}(F, H_{(s_0)}) t^k.$$

7.3 Proof of Corollary 2. Consider the \star -product (8) given by $F \star_{\hbar}^{\Phi} G$, where $F \star_{\hbar} G$ is the Moyal product (1), and $\Phi : (p,q) \mapsto (p^2/2, q/p)$. It is defined on $\mathbf{R}^2 \setminus \mathbf{R}$.

The *-product (8) on the algebra $C^{\infty}(\mathbf{R}^2 \setminus \mathbf{R})$ is equivalent to the Moyal product. Indeed, the symplectomorphism Φ is isotopic to the identity in the group of all smooth symplectomorphisms of $\mathbf{R}^2 \setminus \mathbf{R}$. The isotopy is: $\Phi_s : (p,q) \mapsto (\frac{p^{1+s}}{1+s}, \frac{q}{p^s})$, where $s \in [0,1]$.

Recall that the \star -product $F \star^{\Phi}_{\hbar} G$ on the algebra \mathcal{F} is not equivalent to the Moyal product since it coincides with the product (5).

The family Φ_s does not preserve the algebra \mathcal{F} . Theorem 1 implies that F is not isotopic to the identity in the group of symplectomorphisms of $\mathbf{R}^2 \setminus \mathbf{R}$ preserving the algebra \mathcal{F} .

8. Discussion

8.1 Difficulties in multi-dimensional case.

There exist multi-dimensional analogues of transvectants [14] and [16].

Consider the projective space \mathbf{RP}^{2n+1} endowed with the standard contact structure (or an open domain of the complex projective space \mathbf{CP}^{2n+1}). There exists an unique bilinear differential operator of order k on tensor-densities equivariant with respect to the action of the group Sp_{2n} (see [14], [16]):

(15)
$$J_k^{\lambda,\mu}: \mathcal{F}_\lambda \otimes \mathcal{F}_\mu \to \mathcal{F}_{\lambda+\mu-\frac{k}{n+1}},$$

where $\mathcal{F}_{\lambda} = \mathcal{F}_{\lambda}(\mathbf{P}^{2n+1})$ is the space of tensor-densities on \mathbf{P}^{2n+1} of degree $-\lambda$:

$$f = f(x_1, \dots, x_{2n+1})(dx_1 \wedge \dots dx_{2n+1})^{-\lambda}.$$

The space of tensor-densities $\mathcal{F}_{\lambda}(\mathbf{RP}^{2n+1})$ is isomorphic as a module over the group of contact diffeomorphisms to the space of homogeneous functions on \mathbf{R}^{2n+2} , and the isomorphism is given by:

$$f \mapsto F(y_1, \dots, y_{2n+2}) = y_{2n+2}^{-\lambda(n+1)} f(\frac{y_1}{y_{2n+2}}, \dots, \frac{y_{2n+1}}{y_{2n+2}}).$$

Then the operations (15) are defined as the restrictions of the terms of the standard \star -product on \mathbf{R}^{2n+2} .

The same formula (5) defines a \star -product on the space of tensordensities on \mathbb{CP}^{2n+1}) (cf. [16]). However, there is no analogues of the symplectomorphism (7). I do not know if there exists a \star -product on the Poisson algebra $\mathbb{C}[y_{2n+2}, y_{2n+2}^{-1}] \otimes C^{\infty}(\mathbb{RP}^{2n+1})$ non-equivalent to the standard.

8.2 Classification problem. The classification (modulo equivalence) of \star -products on the Poisson algebra \mathcal{F} is an interesting open problem. It is related to the calculation of cohomology groups $H^2(\mathcal{F}; \mathcal{F})$ and $H^3(\mathcal{F}; \mathcal{F})$. The following result was announced in [7]: dim $H^2(\mathcal{F}; \mathcal{F}) =$ 2.

Let us formulate a conjecture in the compact case. Consider the Poisson algebra $\mathcal{F}(S^1)$ of functions on $T^*S^1 \setminus S^1$ which are Laurent polynomials on the fiber: $F(p,q) = \sum_{-N \leq i \leq N} p^i f_i(q)$ where $f_i(q+1) = f_i(q)$.

Conjecture. Every \star -product on $\mathcal{F}(S^1)$ is equivalent to (1) or (5).

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