

Near-critical Ising model

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ENS Lyon and CNRS

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Plan

- 1 Near-critical behavior, case of percolation
 - ▶ Notion of **correlation length** $L(p)$

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 - ▶ Joint work with **H. Duminil-Copin** and **Gábor Pete**.

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- 3 Near-critical Ising model as the **external magnetic field** varies
 - ▶ Joint work with **F. Camia** and **C. Newman**.

Near criticality

Consider your favorite statistical physics model, for example:

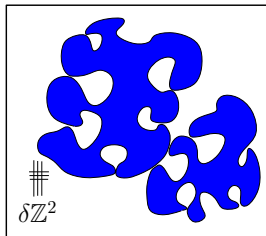
- ▶ percolation
- ▶ FK percolation
- ▶ Ising model etc ...

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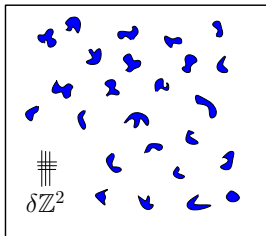


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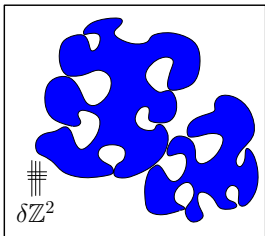
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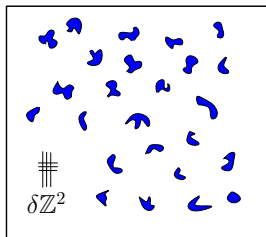


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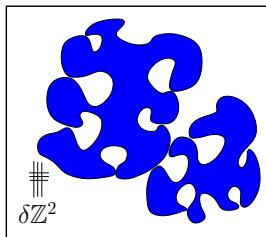
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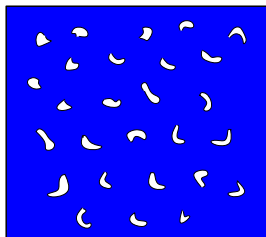
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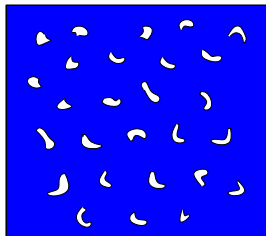
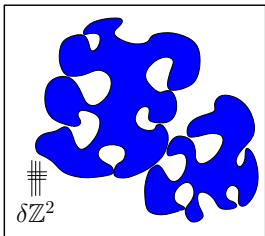
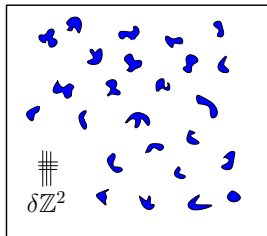


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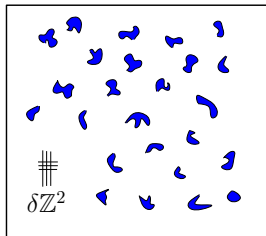


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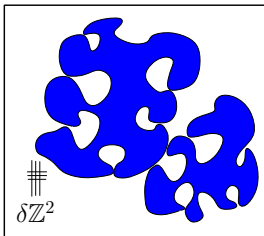
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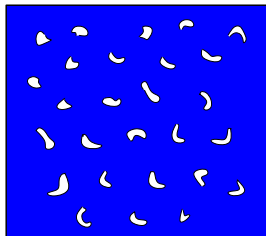
Sub-critical



$T = T_c$ and $h = 0$



$T = T_c$ and $h > 0$

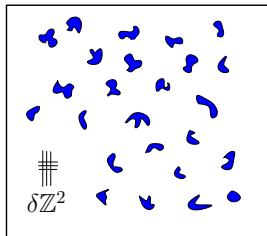


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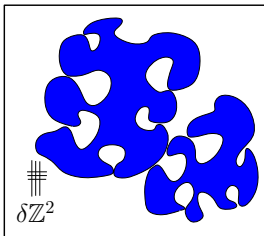
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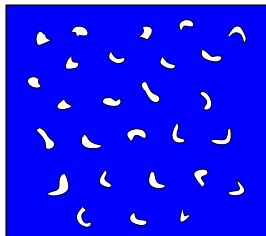
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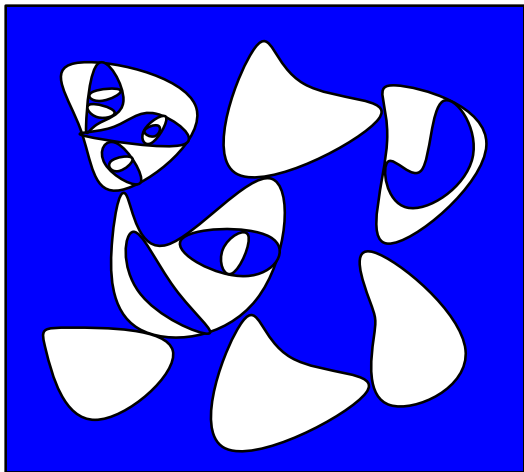


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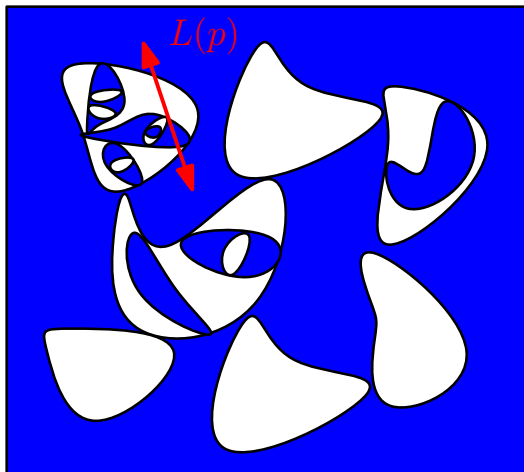
What happens if $T \approx T_c$ or $h \approx 0$??

Notion of correlation length (informal)



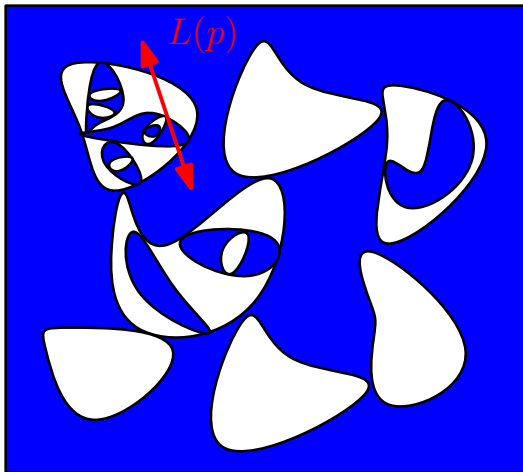
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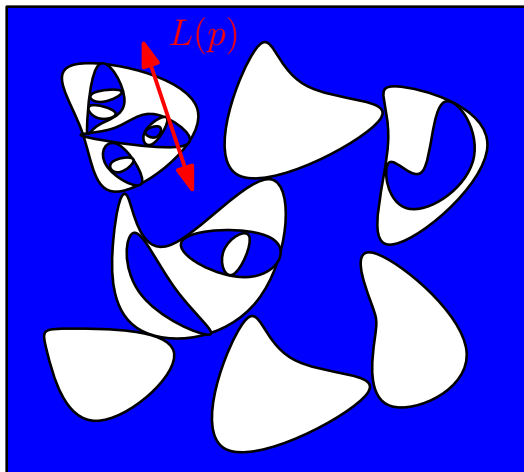
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Example (critical percolation):

Theorem (Smirnov-Werner 2001):

$$L(p) = \left| \frac{1}{p-p_c} \right|^{4/3+o(1)}$$

The models we shall consider

Percolation:

$$P_p(\omega) = p^o (1 - p)^c$$

$o = o(\omega) = \text{Nb of open ed}$

$c = c(\omega) = \text{Nb of closed ed}$

FK Percolation (or random cluster model)

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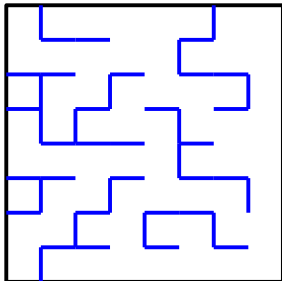
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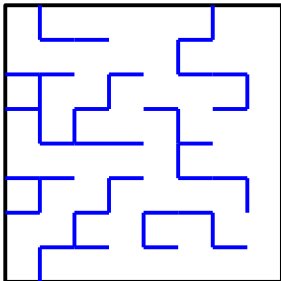
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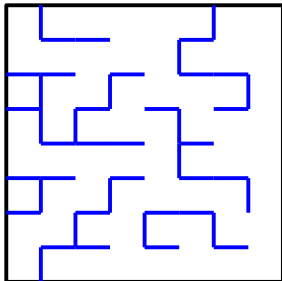
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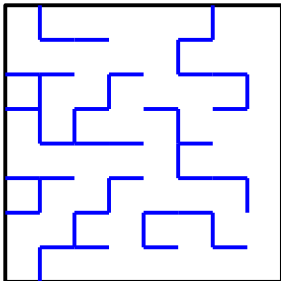
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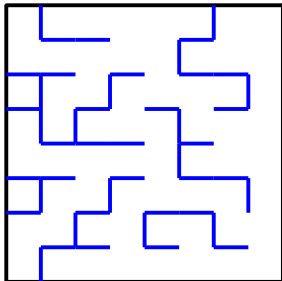
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Theorem (Kesten 1980)

$$p_c(\mathbb{Z}^2) = \frac{1}{2}$$



Theorem (Beffara, Duminil-Copin 2010)

$$p_c(q) = \frac{\sqrt{q}}{1 + \sqrt{q}}$$

Notion of correlation length (precise definition)

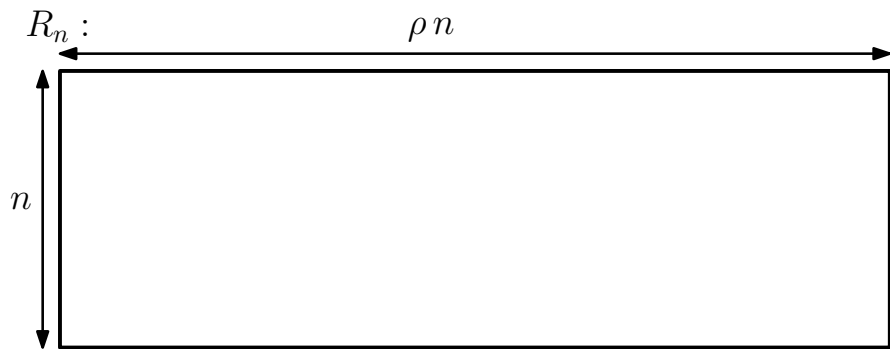
Definition

Fix $\rho > 0$.

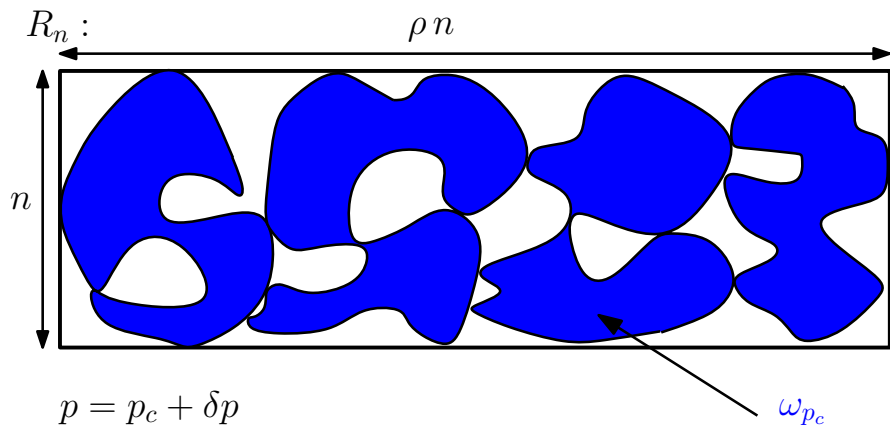
For any $n \geq 0$, let R_n be the rectangle $[0, \rho n] \times [0, n]$. If $p > p_c$, then define for all $\epsilon > 0$ and all “boundary conditions” ξ around R_n ,

$$L_{\rho, \epsilon}^{\xi}(p) := \inf_{n > 0} \left\{ \mathbb{P}_p^{\xi}(\text{there is a left-right crossing in } R_n) > 1 - \epsilon \right\}$$

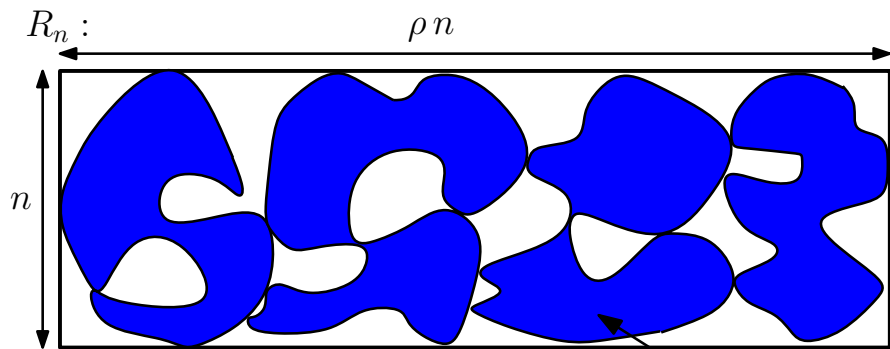
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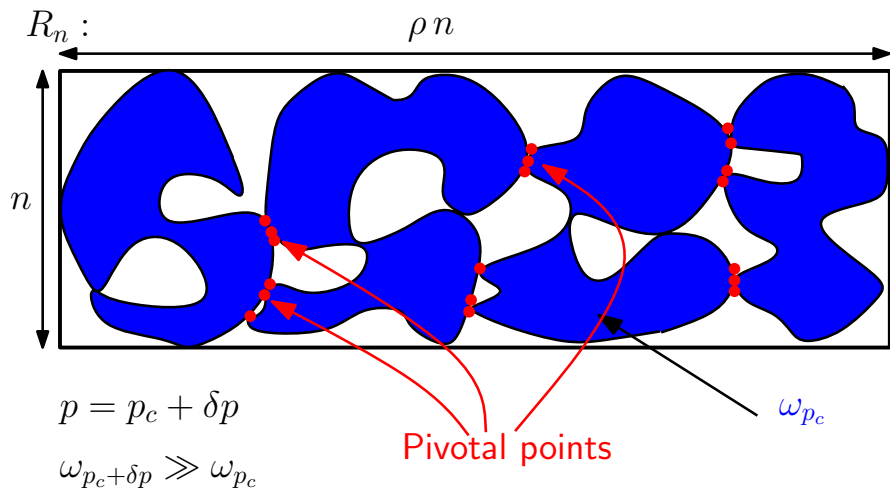


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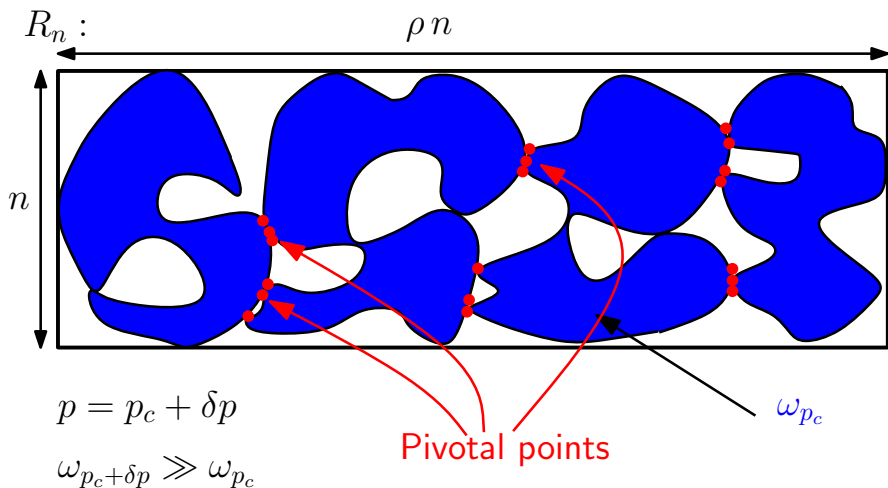
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ω_{p_c}

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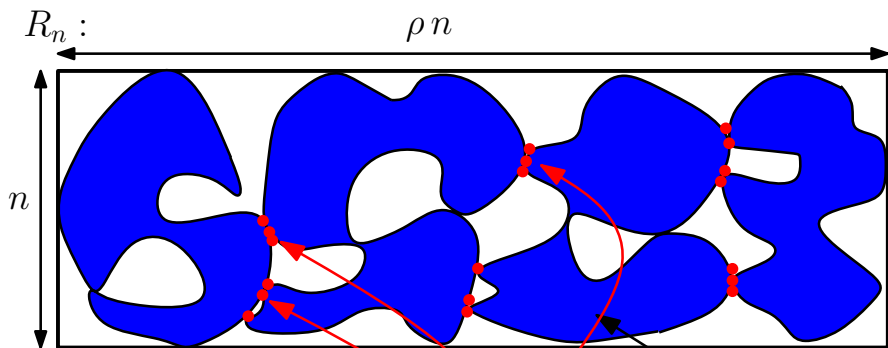
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$$\begin{aligned} \#(\text{Pivotal points}) \\ \approx n^2 \alpha_4(n) \approx n^{3/4} \end{aligned}$$

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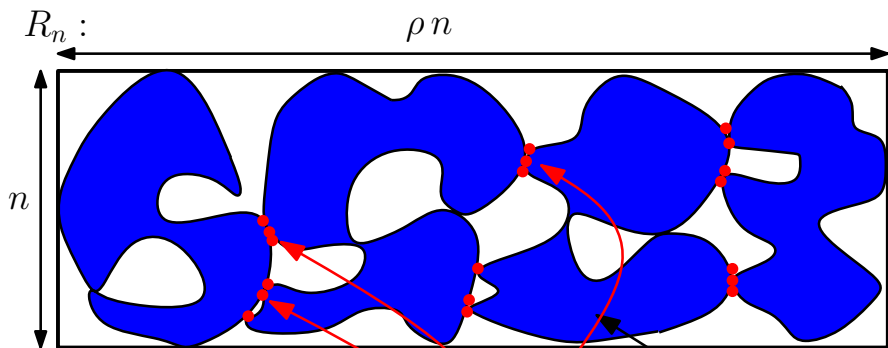
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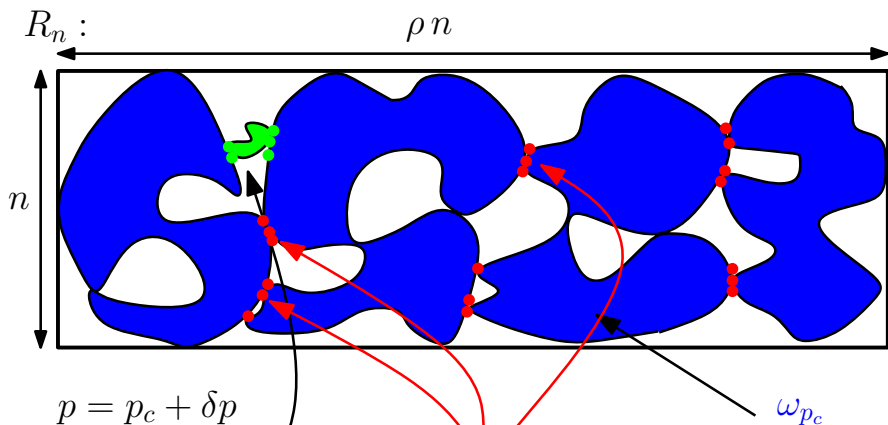
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Difficulty !

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Sharp threshold

To analyze the behavior of the correlation length, it is useful to rely on Russo's formula: if $\phi_n(p) := \mathbb{P}_p(\text{there is a left-right crossing in } R_n)$, then

$$\begin{aligned} \frac{d}{dp} \phi_n(p) &= \mathbb{E}_p(\text{Number of pivotal points in } \omega_p) \\ &= \sum_{x \in R_n} \mathbb{P}_p(x \text{ is a pivotal point}) \end{aligned}$$

This point of view also leads to the identity

$$|p - p_c| L(p)^2 \alpha_4(L(p)) \asymp 1$$

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But this does not match with related results known since Onsager which suggest that $L(p)$ should instead scale like $\left| \frac{1}{p - p_c} \right| \ll \left| \frac{1}{p - p_c} \right|^{24/13} !!$

So what is wrong here !?

Monotone couplings of FK percolation, $q = 2$

Grimmett constructed in 1995 a somewhat explicit **monotone coupling** of FK percolation configurations $(\omega_p)_{p \in [0,1]}$. This monotone coupling differs in several essential ways from the standard monotone coupling ($q = 1$):

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Most remains unknown regarding the structure of these random clouds.

What we can prove

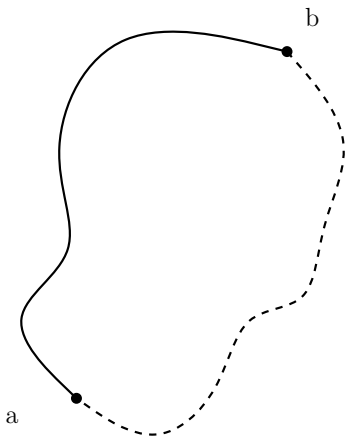
Theorem (Duminil-Copin, G., Pete, 2011)

Fix $q = 2$. For every $\epsilon, \rho > 0$, there is a constant $c = c(\epsilon, \rho) > 0$ s.t.

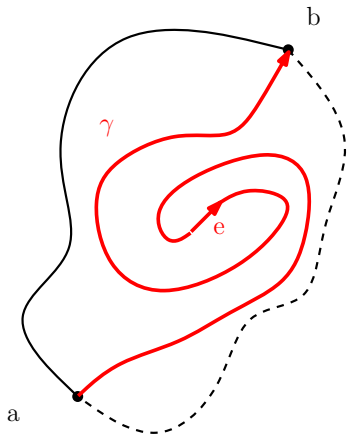
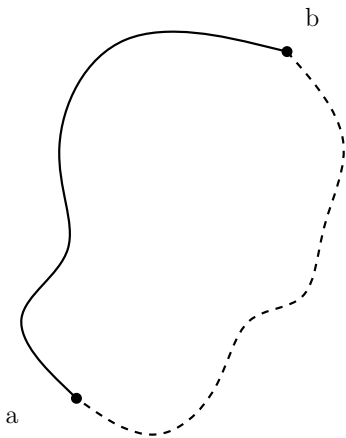
$$c \frac{1}{|p - p_c|} \leq L_{\rho, \epsilon}^{\xi}(p) \leq c^{-1} \frac{1}{|p - p_c|} \sqrt{\log \frac{1}{|p - p_c|}}$$

for all $p \neq p_c$, whatever the choice of the boundary condition ξ is.

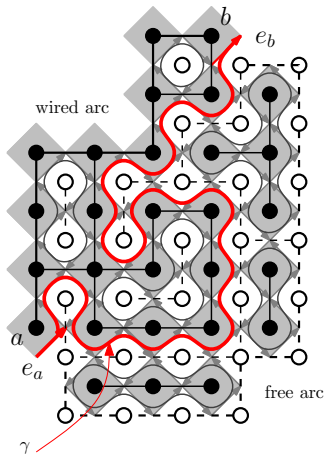
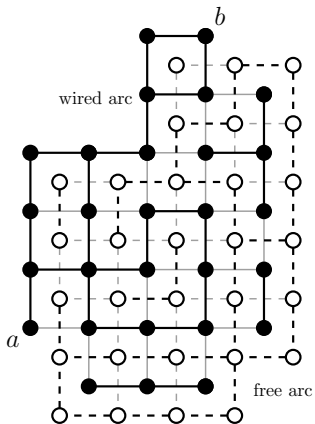
Techniques behind the proof: Smirnov's observable



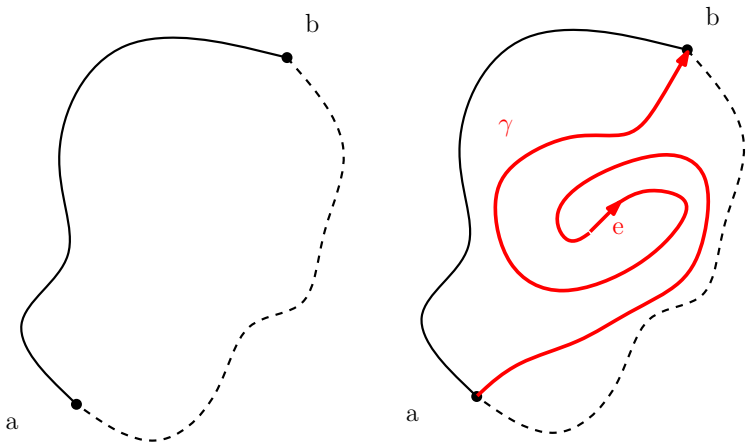
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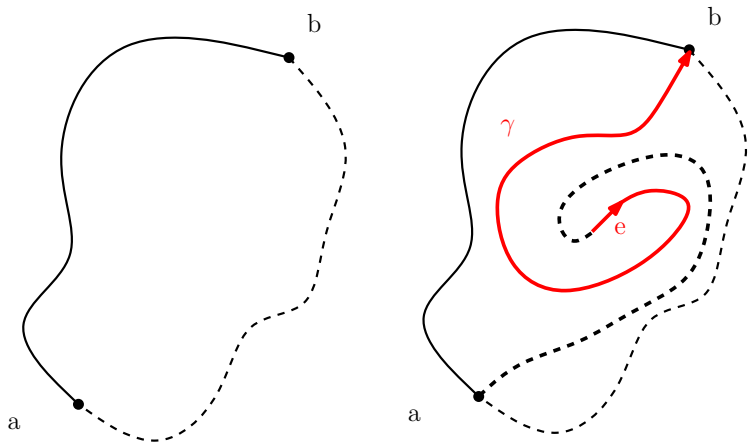


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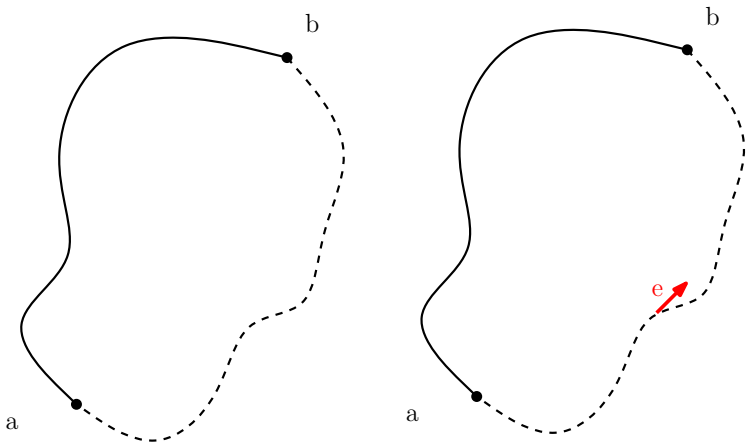
$$F_p(e) := \mathbb{E}_{p,2} \left(e^{\frac{i}{2} W_\gamma(e,b)} 1_{e \in \gamma} \right)$$

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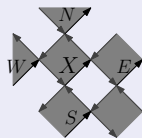
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“Near-harmonicity” of Smirnov’s observable

Theorem (Smirnov, exact harmonicity at criticality)

For $q = 2$ and $p = p_c(2) = \sqrt{2}/(1 + \sqrt{2})$, once restricted to a proper sub-lattice (NE pointing edges), the observable F_{p_c} is **harmonic**:

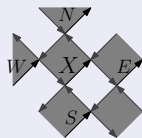
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Theorem (Beffara, Duminil-Copin)

When $p < p_c$, the observable F_p is now **massive harmonic**: namely

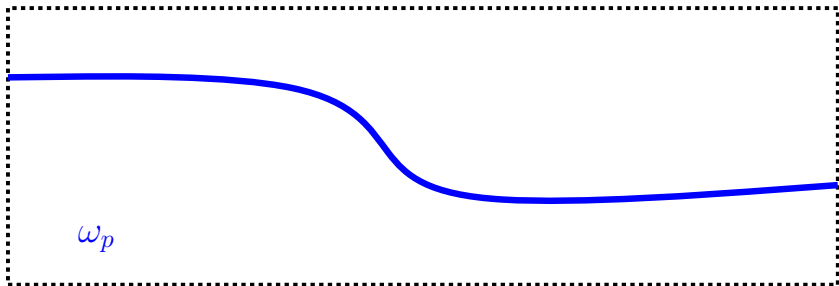
$$\Delta F_p(e_X) = m(p) F_p(e_X),$$

where the **mass** $m(p) \asymp |p - p_c|^2$.

Upper-bound on the correlation length

Fix $\rho, \epsilon > 0$. For any $\rho > \rho_c$, we want to find a scale n so that the rectangle R_n is crossed horizontally with high probability ($> 1 - \epsilon$).

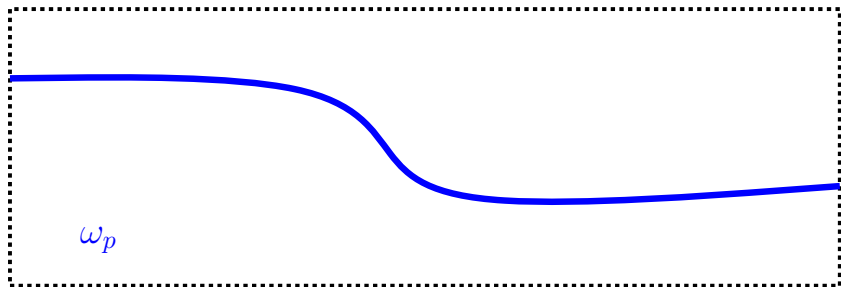
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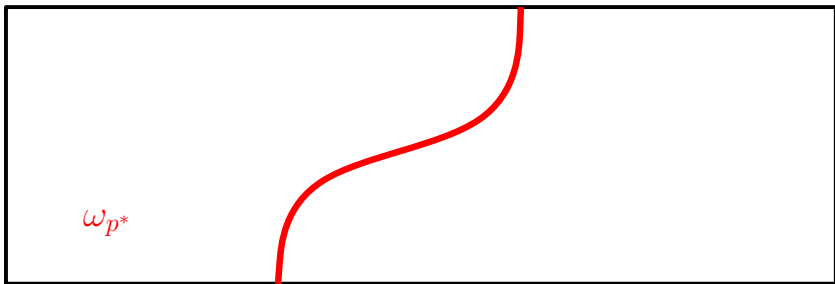


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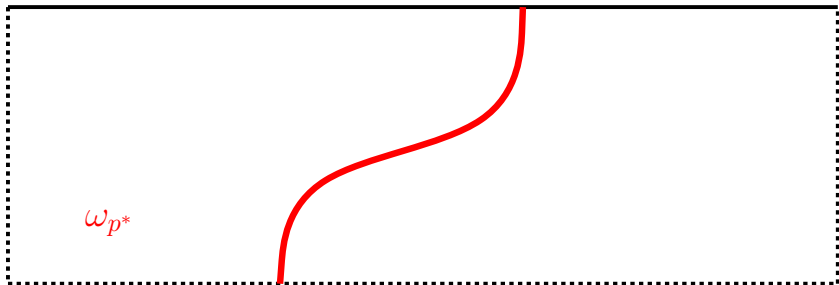
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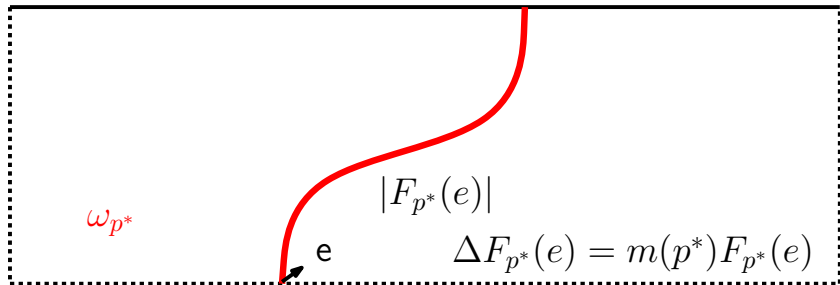
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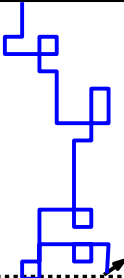
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Random Walk interpretation: at each step in the **bulk**, the mass of the particle is divided by $1+m$

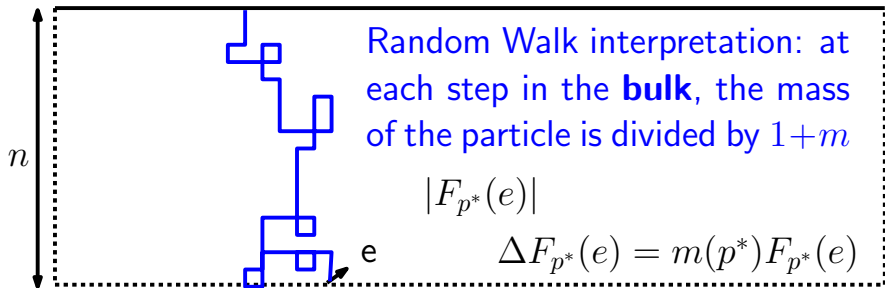
$$|F_{p^*}(e)|$$
$$\Delta F_{p^*}(e) = m(p^*)F_{p^*}(e)$$

e

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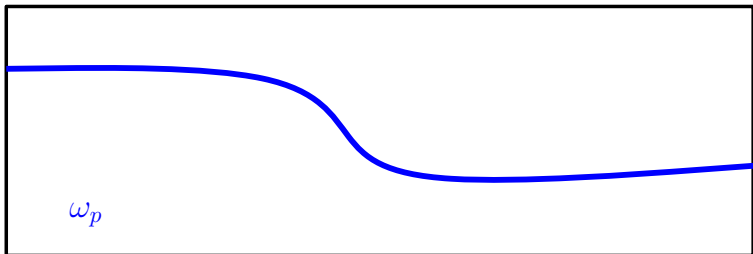
Recall $m(p^*) \asymp |p^* - p_c|^2 \asymp |p - p_c|^2$

If scale $n \gg |p - p_c|^{-1}$, then the RW does more than $|p - p_c|^{-2}$ steps and thus, in average its mass goes to zero.

Lower-bound on the correlation length

Fix $\rho, \epsilon > 0$. For any $p > p_c$, we want to find scales n so that the rectangle R_n is NOT crossed horizontally with high probability (i.e. with prob $< 1 - \epsilon$).

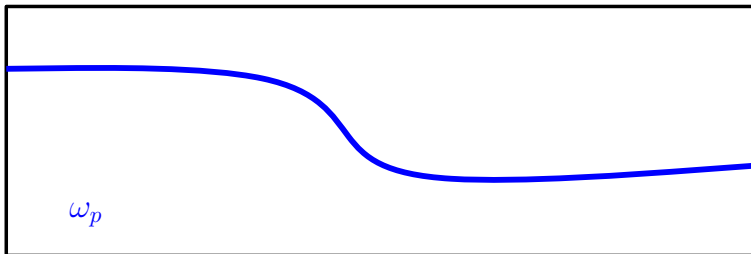
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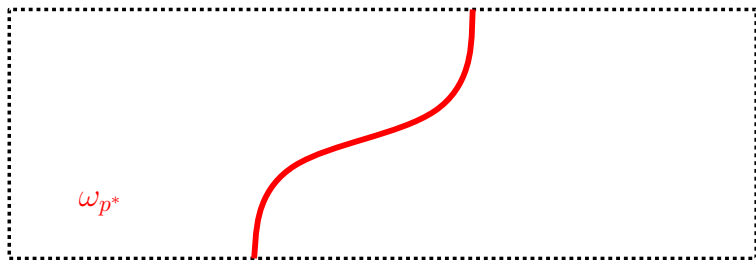
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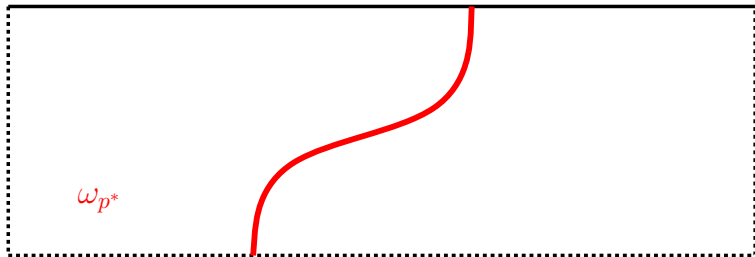
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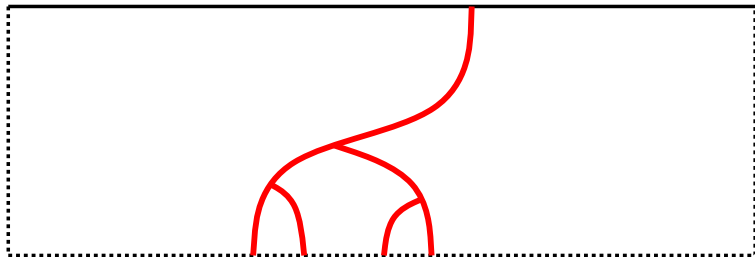
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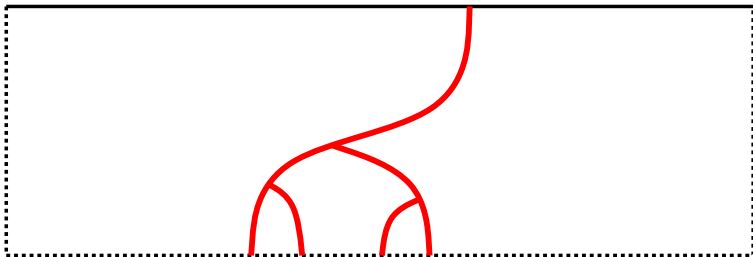


Introduce $N := \text{Nb of lower points connected to the top}$

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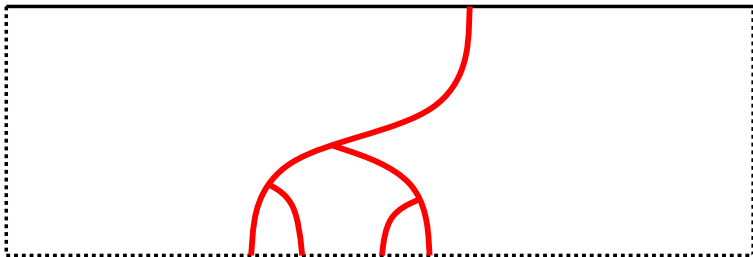
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Using the RSW proof at p_c from Duminil-Copin, Hongler, Nolin, $E_{p^*}(N^2) \leq E_{p_c}(N^2) < c^{-1}n$

Ising model

To each configuration $\sigma \in \{-1, 1\}^{N^2}$, one associates the Hamiltonian

$$H_h(\sigma) := - \sum_{i \sim j} \sigma_i \sigma_j - h \sum \sigma_i$$

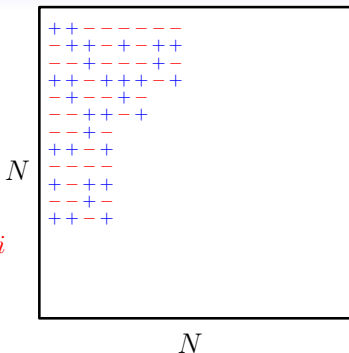
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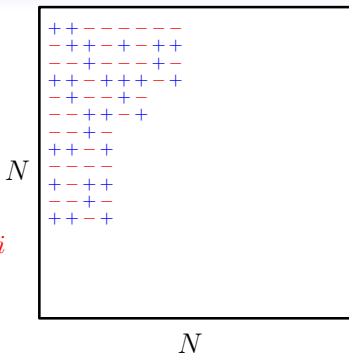
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The Ising model is intimately related with FK percolation ($q = 2$) via the following identity: If $h = 0$,

$$E_{\beta}(\sigma_x \sigma_y) = P_{p, q=2}[x \leftrightarrow y] \text{ with } 1 - p = e^{-2\beta}$$



Classical near-critical results

Theorem (Kesten - Smirnov/Werner)

For site percolation on the triangular lattice,

$$\theta(p) := \mathbb{P}[0 \leftrightarrow \infty] = |p - p_c|^{5/36+o(1)} \quad \text{as } p \searrow p_c$$

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Theorem (Onsager, 1944)

For Ising model on \mathbb{Z}^2 :

$$\langle \sigma_0 \rangle_\beta^+ \asymp |\beta - \beta_c|^{1/8} \quad \text{as } \beta \searrow \beta_c$$

Average magnetization under small external field

Theorem (Camia, G., Newman)

Consider Ising model on \mathbb{Z}^2 at β_c with a positive external magnetic field $h > 0$, then

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- ▶ Lower bound: prove that the correlation length $L(h) \asymp h^{-8/15}$ and conclude using

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- ▶ Upper bound: rely on a kind of strong “convexity” property satisfied by the Ising model, namely the **GHS inequality**.

Theorem (**GHS** inequality, Griffiths, Hurst, Sherman, 1970)

$Z_{\beta,h} := \sum_{\sigma} e^{-\beta H(\sigma) + h \sum \sigma_x}$ is such that

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