

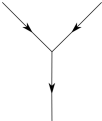
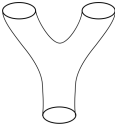


Classifying Conformal Field Theories

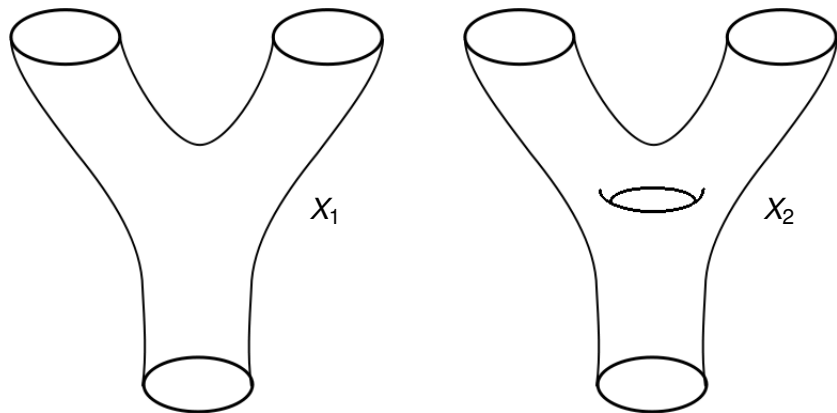
Leonard Hardiman

String Theory

String theory replaces the *point-like particles* of particle physics with one-dimensional objects called *strings*.

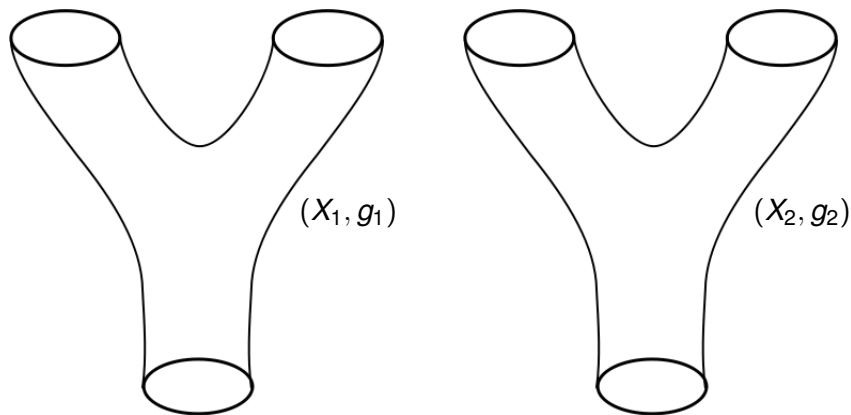
Quantum Field Theory	String Theory
	
	

Topological Quantum Field Theories



Setting $X_1 \sim X_2$ if X_1 is topologically equivalent to X_2 would lead us towards a *topological quantum field theory*.

Conformal Quantum Field Theories



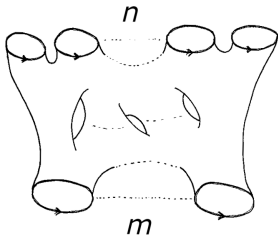
Setting $(X_1, g_1) \sim (X_2, g_2)$ if (X_1, g_1) is conformally equivalent to (X_2, g_2) leads us towards a *conformal field theory*.

Segal's Category

$$\text{Obj}(\mathcal{C}) = \left\{ \underline{n} := \bigsqcup_n S^1 \mid n \in \mathbb{N} \right\}$$

$$\underline{n} = \underbrace{\bigcirc \bigcirc \bigcirc \dots \bigcirc}_n$$

$$\text{Hom}(\underline{n}, \underline{m}) = \{ [X] \mid X \text{ is a Riemannian Manifold with } n \text{ incoming boundaries and } m \text{ outgoing boundaries.} \}$$



Definition of a Conformal Field Theory

Let $H \in \mathbf{VECT}$ be a Hilbert space. Then a conformal field theory on H is a monoidal functor

$$\mathcal{U} : \mathcal{C} \rightarrow \mathbf{VECT}$$

that sends $\underline{1}$ to H (and satisfies certain other axioms).

Remark

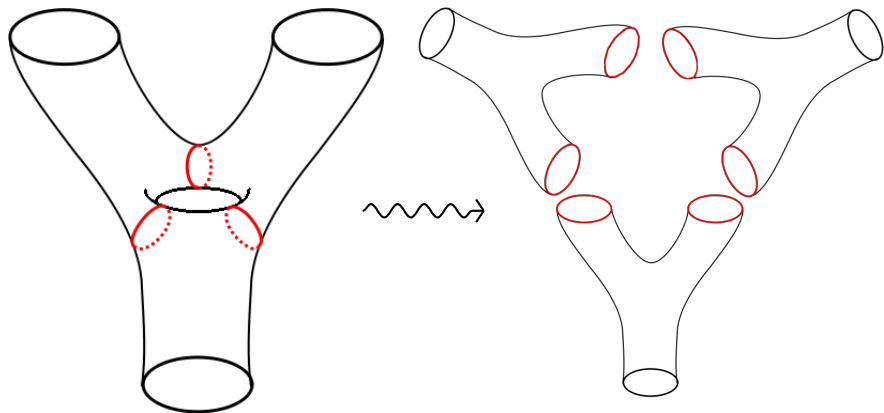
The fact that \mathcal{U} is monoidal implies $\mathcal{U}(\underline{n}) = H^{\otimes n}$. Therefore $X \in \text{Hom}(\underline{n}, \underline{m})$ is sent onto a linear map from $H^{\otimes n}$ to $H^{\otimes m}$.

An important question in the field is the classification of possible conformal field theories. In other words how much data do we need to construct such a functor?

The answer can be summed up via the following mnemonic:

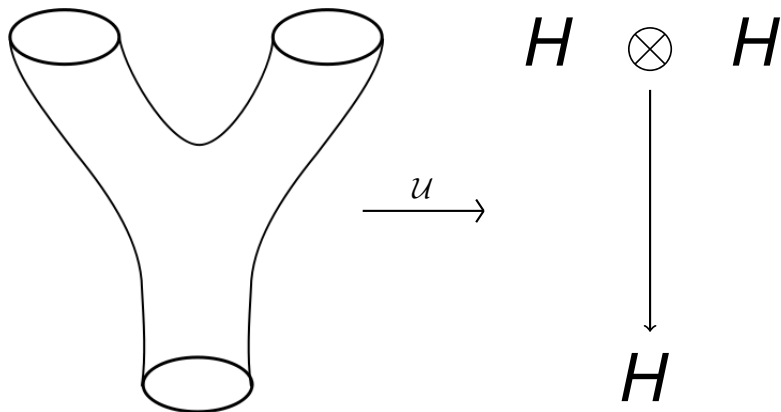
Donut Pants

Pants (i.e. Trousers)



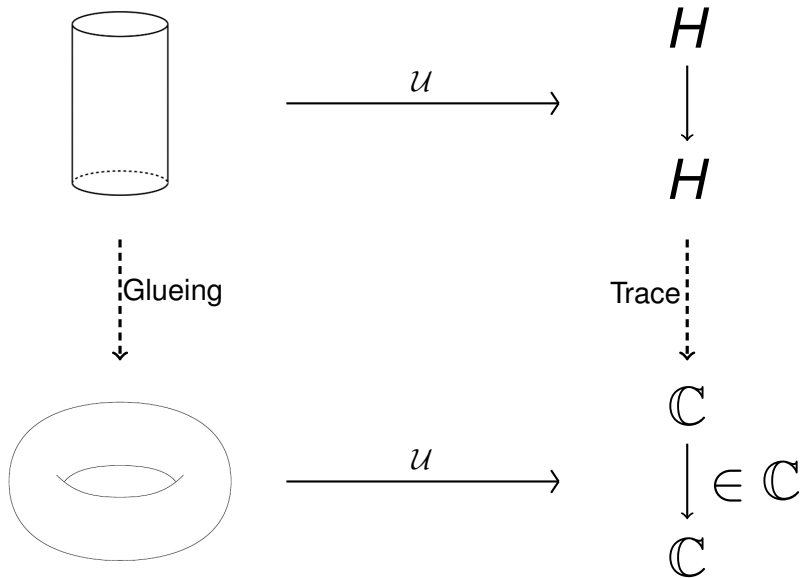
All elements in $\text{Hom}(\underline{n}, \underline{m})$ can be cut up into *pairs of pants*.

Pants (i.e. Trousers)

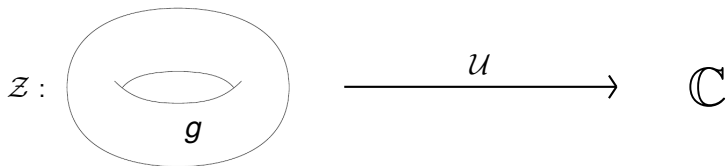


Under our functor \mathcal{U} a pair of pants (equipped with a conformal structure) is sent onto a linear map $H \otimes H \rightarrow H$. The Hilbert space H is therefore equipped with an infinite sequence of inner products, this gives rise to a V.O.A.

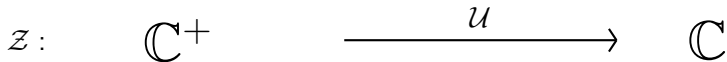
Donuts (i.e. Glueing)



Modularity of the Partition Function

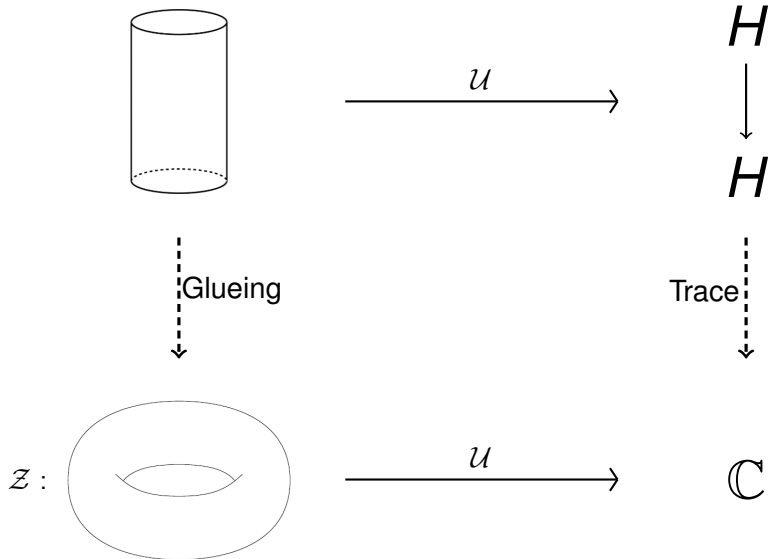


Conformal structures on the torus are parametrized by $\tau \in \mathbb{C}^+ / \mathrm{PSL}_2(\mathbb{Z})$.

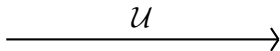


and z is a *modular* function (i.e. invariant under the action of $\mathrm{PSL}_2(\mathbb{Z})$).

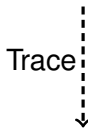
WZW Models



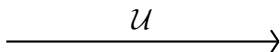
WZW Models

 H  H 

Glueing



Trace

 $z :$ \mathbb{C}^+  \mathbb{C}

WZW Models

$$\begin{array}{ccc} \mathfrak{g} \oplus \bar{\mathfrak{g}} & \xrightarrow{\mathcal{U}} & \begin{array}{c} H \\ \downarrow \\ H \end{array} \\ \downarrow \text{Glueing} & & \text{Trace} \downarrow \\ \mathcal{Z} : \mathbb{C}^+ & \xrightarrow{\mathcal{U}} & \mathbb{C} \end{array}$$

Where \mathfrak{g} is an affine algebra. We therefore get a representation of $\mathfrak{g} \oplus \bar{\mathfrak{g}}$. Any such representation is a direct sum of irreducible representations. When we take the trace of this representation we get the character, which will be the sum of the characters of the irreducibles.

$$\mathcal{Z}(\tau) = \sum_{\lambda, \mu} M_{\lambda\mu} H_{\lambda} \otimes \overline{H_{\mu}}, \quad M_{\lambda\mu} \in \mathbb{Z}^{\geq 0}$$

Modularity in WZW models

The partition function

$$\mathcal{Z}(\tau) = \sum_{\lambda, \mu} M_{\lambda\mu} H_{\lambda} \otimes \overline{H_{\mu}}, \quad M_{\lambda\mu} \in \mathbb{Z}^{\geq 0}$$

is completely determined by the integer matrix $M = (M_{\lambda\mu})$.

For every affine algebra \mathfrak{g} there exists matrices S and T such that

$$\mathcal{Z} \text{ modular} \iff M \text{ commutes with } S \text{ and } T.$$

Therefore the problem of classifying WZW models has been reduced to the purely algebraic problem of finding integer matrices that commute with S and T ! Such matrices are called physical invariants.

The CIZ Classification

Physical invariants of WZW models for which $\mathfrak{g} = \widehat{\mathfrak{sl}}_2$ were classified by Cappelli, Itzykson and Zuber in 1987.

Their classification had an **A-D-E** meta-pattern!

The goal is to understand why this occurs. Current research is focused around the fact that $\mathbf{TL}(-[2]_q)$ is a subcategory in $\text{Rep}(\widehat{\mathfrak{sl}}_2)$ and representations of $\mathbf{TL}(-[2]_q)$ that generate rational conformal theories are classified by **A-D-E**.