Deterministic, stochastic and strategical dynamics under density constraints

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Deterministic Dynamics - Micro and macro models for crowd motions with constraints

- Disks with no overlapping
- $\bullet~{\rm Density}~\rho\leq 1$
- The continuity equation
- Optimistic Dynamics Solution by optimal transport
 - A splitting method
 - Few words about optimal transport and Wasserstein distance
 - Recovering the PDE
- Strategical Dynamics MFG with density penalization or constraints
 - Classical MFG with density penalization
 - How to replace penalizations by constraints
 - Questions
- Stochastic Dynamics Let's introduce diffusion
 - The PDEs for MFG and crow motion
 - Three possible different schemes

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Deterministic - models

Micro and macro models with constraints

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A general model

A particle population needs to move, and each particle - if alone - would follow its own velocity u (depend on time, position...) Yet, particles are rigid disks that cannot overlap, hence, the actual velocity v will not be u if u is too concentrating. Let us suppose $v = P_{adm(q)}(u)$, where q is the particle configuration, adm(q) the set of velocities that do not induce overlapping, $P_{adm(q)}$ the projection on this set.

If every particle is a disk with radius R, located at q_i , we have

$$q \in K := \{q = (q_i)_i \in \Omega^N : |q_i - q_j| \ge 2R\}$$

 $adm(q) = \{v = (v_i)_i : (v_i - v_j) \cdot (q_i - q_j) \ge 0 \ \forall (i,j) : |q_i - q_j| = 2R\}$

and we solve $q'(t) = P_{adm(q(t))}u(t)$ (with q(0) given).

B. MAURY, J. VENEL, Handling of contacts in crowd motion simulations, *Traffic* and Granular Flow, 2007.

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Continuous formulation

• The particles population will be described by a density $\rho \in \mathcal{P}(\Omega)$,

• the constraint by $K = \{ \rho \in \mathcal{P}(\Omega) : \rho \leq 1 \}$,

- $u: \Omega \to \mathbb{R}^d$ will be a vector field, possibly depending on time or ρ ,
- $adm(\rho) = \{ v : \Omega \to \mathbb{R}^d : \nabla \cdot v \ge 0 \text{ on } \{ \rho = 1 \} \},$
- *P* is the projection in $L^2(dx)$ or (which is the same) in $L^2(\rho)$,
- we'll solve $\partial_t \rho_t + \nabla \cdot \left(\rho_t \left(P_{adm(\rho_t)} u_t \right) \right) = 0.$

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- $u: \Omega \to \mathbb{R}^d$ will be a vector field, possibly depending on time or ρ ,
- $adm(\rho) = \{ \mathbf{v} : \Omega \to \mathbb{R}^d : \nabla \cdot \mathbf{v} \ge 0 \text{ on } \{ \rho = 1 \} \},$
- P is the projection in L²(dx) or (which is the same) in L²(ρ),
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Difficulty : $v = P_{adm(\rho_t)}u_t$ is not regular, neither depends regularly on ρ .

Pressures and duality

The set $adm(\rho)$ may be better described by duality :

$$\mathit{adm}(
ho)=\{v\in L^2(
ho)\ :\ \int v\cdot
abla p\leq 0\quad orall p\ :\ p\geq 0,\ p(1-
ho)=0\}.$$

In this way we can characterize $v = P_{adm(\rho)}(u)$ through

$$u = v + \nabla p, \quad v \in adm(
ho), \quad \int v \cdot \nabla p = 0,$$

 $p \in press(
ho) := \{p \in H^1(\Omega), \ p \ge 0, \ p(1 -
ho) = 0\}$

This function *p* plays the role of the pressure affecting the movement.

$$\partial_t \rho_t + \nabla \cdot \left(\rho_t (u_t - \nabla p_t) \right) = 0$$

$$\rho_t \in K, \ p_t \in press(\rho_t)$$

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Deterministic - solutions

Using optimal transport and Wasserstein distance

A splitting scheme for the PDE

Fix a time step $\tau > 0$. We look for a sequence $(\rho_n^{\tau})_n$ where ρ_n^{τ} stands for ρ at time $n\tau$. We first define

$$\tilde{\rho}_{n+1}^{\tau} = (id + \tau u_{n\tau})_{\#} \rho_n^{\tau}; \quad \rho_{n+1}^{\tau} = P_{\mathcal{K}}(\tilde{\rho}_{n+1}^{\tau})$$

where the projection $P_{\mathcal{K}}$ is in the sense of the Wasserstein distance, induced by optimal transport.

The key point is actually using the W_2 projection (instead of L^2 or other projections). It corresponds to the L^2 projection of velocity fields.

B. MAURY, A. ROUDNEFF-CHUPIN AND F. SANTAMBROGIO, A macroscopic crowd motion model of gradient flow type, *Math. Mod. Meth. Appl. Sci.*

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Optimal transport and Wasserstein distances

If two probabilities $\mu,\nu\in\mathcal{P}(\Omega)$ are given on a compact domain, the Monge-Kantorovitch problem reads

$$\begin{split} W_2^2(\mu,\nu) &= \inf \left\{ \int |x - T(x)|^2 d\mu \ : \ T : \Omega \to \Omega, \ T_{\#}\mu = \nu \right\} \\ &= \inf \left\{ \int |x - y|^2 d\gamma \ : \ \gamma \in \mathcal{P}(\Omega^2), \ (\pi_x)_{\#}\gamma = \mu, \ (\pi_y)_{\#}\gamma = \nu \right\} \\ &= 2 \sup \left\{ \int \phi \ d\mu + \int \psi \ d\nu \ : \ \phi(x) + \psi(y) \le \frac{1}{2} |x - y|^2 \right\}. \end{split}$$

Under suitable assumptions, there exist an optimal transport T and an optimal function ϕ , called Kantorovich potential, which is Lipschitz continuous. They are linked by $T(x) = x - \nabla \phi(x)$.

Moreover, $W_2(\mu, \nu)$, the square root of the minimal value, is a distance on $\mathcal{P}(\Omega)$ which metrizes the weak-* convergence of probabilities (on compact domains).

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Projections and pressures

Fix a measure $\nu \in \mathcal{P}(\Omega)$ and solve

$$\min \quad \frac{1}{2}W_2^2(\rho,\nu) \ : \ \rho \in \mathcal{K} \quad = \quad \min_{\rho \leq 1} \ \sup_{\phi,\psi} \ \int \phi \ d\rho + \int \psi \ d\nu.$$

By duality and inf-sup exchange arguments, the optimal ho must also solve

$$\min\int \phi \, d
ho ~:~
ho \leq 1,$$

where ϕ is the Kantorovich potential in the transport from ρ to $\nu.$ This implies

 $\exists t \ : \
ho = egin{cases} 1 & ext{on } \phi < t, \ 0 & ext{on } \phi > t, \Rightarrow p := (t-\phi)_+ \geq 0, \ p(1ho) = 0. \ \in [0,1] & ext{on } \phi = t \end{cases}$

Hence, $p \in press(\rho)$ and, passing to gradients, we have

$$\rho$$
 - a.e. $\nabla p = -\nabla \phi = T(x) - x$.

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Getting back to the PDE

 $T(x) = x + \nabla p(x)$ is the optimal transport from ρ_{n+1}^{τ} to $\tilde{\rho}_{n+1}^{\tau}$. Notice

$$|\nabla \rho||_{L^{2}(\rho_{n+1}^{\tau})} = W_{2}(\rho_{n+1}^{\tau}, \tilde{\rho}_{n+1}^{\tau}) \leq W_{2}(\rho_{n}^{\tau}, \tilde{\rho}_{n+1}^{\tau}) \leq \tau ||u_{n\tau}||_{L^{2}(\rho_{n}^{\tau})}.$$

This suggest to scale the pressure (we call it now τp) and get the following situation



Notice that $(id + \tau u_{n\tau})^{-1}(id + \tau \nabla p) = id - \tau (u_{(n+1)\tau} - \nabla p) + o(\tau)$ provided u is regular enough. This allows to get, in the limit $\tau \to 0$, the vector field $v_t = P_{adm(\rho_t)}[u_t]$ and get a solution of the PDE.

Strategical dynamics

Mean Field Games with density penalizations or constraints

-

MFG with density penalization- 1

In a population of agents everybody follows controlled trajectories

$$y'(t) = f(t, y(t), \alpha(t)), \quad t \in [0, T].$$

For every t, the goal of each agent is to maximize

$$-\int_t^T \left(\frac{|\alpha(s)|^2}{2} + g(\rho_s(y(s)))\right) ds + \Phi(y(T)),$$

where g is a given increasing function of the density ρ_s at time s. The agent hence tries to avoid overcrowded regions.

Let φ be the value function for this problem : it satisfies

$$\partial_t \varphi(t, x) + H(t, x, \nabla \varphi(t, x)) = 0, \quad \varphi(T, x) = \Phi(x) \quad : \quad y(t) = x$$

for a Hamiltonian function H, depending on f and $g(\rho_t)$. The optimal $\alpha(t)$, and hence the evolution of ρ_t , depends on $\nabla \varphi(t, x)$.

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$$\partial_t \varphi(t,x) + H(t,x,\nabla \varphi(t,x)) = 0, \quad \varphi(T,x) = \Phi(x) \quad : \quad y(t) = x$$

for a Hamiltonian function H, depending on f and $g(\rho_t)$. The optimal $\alpha(t)$, and hence the evolution of ρ_t , depends on $\nabla \varphi(t, x)$.

MFG with density penalization- 2

The evolution follows a coupled system : φ solves HJB with ρ , which on turn evolves according to $\partial_t \rho + \nabla \cdot (\rho v) = 0$, where v(t, x) depends on $\nabla \varphi(t, x)$.

Typical example : if $f(t, x, \alpha) = \alpha$ then we have

$$\begin{cases} \partial_t \varphi + \frac{|\nabla \varphi|^2}{2} - g(\rho) = 0, \\ \partial_t \rho + \nabla \cdot (\rho \nabla \varphi) = 0, \\ \varphi(\mathcal{T}, x) = \Phi(x), \quad \rho(0, x) = \rho_0(x). \end{cases}$$

Stochastic case : we can also insert random effects $dY = f(t, Y, \alpha)dt + dB$, which lets $\Delta \varphi$ and $\Delta \rho$ appear :

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J.-M. LASRY, P.-L. LIONS, Mean-Field Games, Japan. J. Math. 2007 P. CARDALIAGUET, lecture notes, www.ceremade.daughingfr/gcardaliag and

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MFG with density constraints - 1

How to define a mean field game if we want to replace the penalization $+g(\rho)$ with the constraint $\rho \leq 1$?

Naif idea : when $(\rho_s)_s$ is given, every agent minimizes his own cost paying attention to the constraint $\rho_s(y(s)) \leq 1$. But if ρ already satisfies $\rho \leq 1$, one extra agent will not violate the constraint (non-atomic game). Hence the constraint becomes empty.

Good idea : use the pressure. The unnown is now the pair $(\overline{\rho}, \overline{\alpha})$. We want

$$\partial_t \overline{\rho}_t + \nabla \cdot \left(\overline{\rho}_t \left(P_{adm(\rho_t)}[\overline{\alpha}_t] \right) \right) = 0.$$

The projection of $\overline{\alpha}_t$ onto $adm(\overline{\rho}_t)$ rises a pressure p_t . Every agent tries to solve

$$\max \quad -\int_t^T \frac{|\alpha(s)|^2}{2} ds + \Phi(y(T)), \quad y'(s) = \alpha(s) - \nabla p_s(y(s)), \quad y(t) = x.$$

A configuration $(\overline{\rho}, \overline{\alpha})$ will be an equilibrium if the original effort field $\overline{\alpha}$ is equal to the optimal one in this problem and if the original densities $\overline{\rho}_t$ are equal to those realized at time t by these optimal trajectories, $\overline{\rho}_t \approx 1$

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JMFG with density constraints - 2

By computing the Hamiltonian of this problem and the optimal α we get the equations satisfied by the equilibrium, i.e.

$$\begin{cases} \partial_t \varphi + \frac{|\nabla \varphi|^2}{2} - \nabla \varphi \cdot \nabla p = 0, \\ \partial_t \rho + \nabla \cdot (\rho(\nabla \varphi - \nabla p)) = 0, \\ p \ge 0, \ p(1 - \rho) = 0, \\ \varphi(T, x) = \Phi(x), \quad \rho(0, x) = \rho_0(x). \end{cases}$$

Unfortunately, no result (existence, uniqueness. . .) is available for the moment.

A PhD thesis (A. Mészáros, Orsay) is ongoing on these questions.

F. SANTAMBROGIO, A Modest Proposal for MFG with Density Constraints, *Net. Het. Media*, 2012.

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Stochastic Dynamics

Some perspectives and methods

Which PDE for diffusion and density constraints?

If we want to model a population which diffuse but is also subject to $\rho \leq 1$, should we apply the projection operator on the drift only or should it also interact with the Laplacian term?

Good definitions (and existence/uniqueness results) for this notion without strategical issues are a necessary starting point in order to attack later the 2nd order MFG system with density constraints (which could turn out to be simpler than the 1st order one, because of higher regularity).

Fortunately, the Laplacian is consistent with the constraint. Moreover $\Delta \rho = \nabla \cdot \left(\rho \frac{\nabla \rho}{\rho} \right)$ and $\frac{\nabla \rho}{\rho} = 0$ on $\{ \rho = 1 \}$. Hence the equation is

$$0 = \partial_t \rho_t - \Delta \rho_t + \nabla \cdot \left(P_{adm(\rho_t)}[u_t] \rho_t \right) = \partial_t \rho_t + \nabla \cdot \left(P_{adm(\rho_t)}\left[u_t - \frac{\nabla \rho_t}{\rho_t} \right] \rho_t \right)$$

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Different catching-up methods

Take $\tau > 0$ and a density ρ_n^{τ} . How to choose ρ_{n+1}^{τ} ?

The cleanest way Take a r.v. $X \sim \rho_n^{\tau}$ and build $(id + \tau u_{n\tau}) \circ X + B_{\tau}$ with *B* a Brownian motion independent of *X*. Define $\tilde{\rho}_{n+1}^{\tau}$ as the law of this new r.v., which is given by

$$\tilde{\rho}_{n+1}^{\tau} = \left((id + \tau u_{n\tau})_{\#} \rho_n^{\tau} \right) * \eta_{\sqrt{\tau}},$$

where η_r is a standard Gaussian of size r. Then set $\rho_{n+1}^{\tau} = P_{\mathcal{K}}(\tilde{\rho}_{n+1}^{\tau})$. **A similar one** Take the solution ρ of the Fokker-Planck equation

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Estimates as $\tau \to 0$ on the two previous methods are difficult, since they require estimates on the distance $W_2(\rho_n^{\tau}, \tilde{\rho}_{n+1}^{\tau})$, which roughly corresponds to estimate on the Heat equation between time 0 and time τ . They are available, but under additional regularity assumptions on the initial datum (namely, it should be at least BV).

Hence here is a less meaningful but easier to handle way of defining the next step.

The most efficient way First build $\tilde{\rho}_{n+1}^{\tau} := (id + \tau u_{n\tau})_{\#} \rho_n^{\tau}$. Then set

$$\rho_{n+1}^{\tau} = \operatorname{argmin} \int \rho \ln \rho + \frac{1}{2\tau} W_2^2(\rho, \tilde{\rho}_{n+1}^{\tau}) \, : \, \rho \in K.$$

This problem is strictly convex and admits a unique solution. This recalls (and actually is taken from) the theory of gradient flows in the Wasserstein space.

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Where the PDE comes from

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We then define $p = (t - \ln \rho - \frac{\phi}{\tau})_+$ and we get $p \in press(\rho)$. Moreover, $\rho - \text{a.e. } \nabla p = -\frac{\nabla \rho}{\rho} - \frac{\nabla \phi}{\tau}$, which shows that we have

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Determinsitic - solutions Strategical - MFG Stochastic - perspectives

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The End

Thanks for your attention

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Filippo Santambrogio Deterministic, stochastic and strategical dynamics under density constraints