

Renormalization Theory, III, Lyon, September 2008 Introduction

Noncommutative Field Theory

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Physical phenomena occur over a wide range of scales from the Planck scale $\ell_P = \sqrt{\hbar G/c^3} \simeq 1.6 \ 10^{-35} \ m$ to the radius of the observable universe, in practice about 45 billion light-years, hence around 4.4 $10^{26} \ m$ or better, 2.7 $10^{61} \ell_P$.

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The universe therefore is made of roughly 61 powers of 10 or 140 powers of e.

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But the fifteen to sixteen scales between ℓ_P and 2 10^{-19} meters (about 1 Tev), are still terra incognita for physics.

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After that treat no new spectacular advance of this type is planned on terra incognita, hence we have some time to deepen our theoretical and mathematical understanding.

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Space-time itself could be of this type; for instance at a certain scale new uncertainty relations could appear between length and width which would generalize Heisenberg's relations.

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Noncommutative Field Theory

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where $\Theta^{\mu\nu}$ is an antisymmetric constant tensor which in the simplest case can be written as:

$$\Theta^{\mu
u}=oldsymbol{ heta}egin{pmatrix} 0&1&(0)\-1&0&(0)\(0)&-1&0\end{pmatrix} \end{pmatrix}$$

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Noncommutative Field Theory

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The unique product (associative, but noncommutative) generated by these relations on (Schwarz class) functions is called the Moyal-Weyl product and writes:

$$(f \star g)(x) = \int \frac{d^4y}{(2\pi)^4} d^4z f(x+y)g(x+z)e^{2iy\wedge z}$$

where $x \wedge y = x^{\mu} \theta_{\mu\nu} y^{
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 - Alternative to current ideas, eg on supersymmetry

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> The relevant species is now the species of ribbon graphs.

Noncommutative Field Theory

Ribbon Graphs

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- **EF** is the number of *external* faces (i.e. containing arriving external lines).

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g = 1 - (V - V)

$$\begin{array}{c} V=3\\ L=3\\ F=2\\ E=6\\ EF=2 \end{array} \end{array} \right\} \implies g=0,$$

















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For instance,



The Euclidean theory with an additional harmonic potential

$$S = \int d^4x \, \left(\frac{1}{2} \partial_\mu \phi \star \partial^\mu \phi + \frac{\Omega^2}{2} (\tilde{x}_\mu \phi) \star (\tilde{x}^\mu \phi) + \frac{\mu_0^2}{2} \phi \star \phi + \frac{\lambda}{4!} \phi \star \phi \star \phi \right) (x)$$

where $\tilde{x}_{\mu} = 2\Theta_{\mu\nu}^{-1} x^{\nu}$, is covariant under a symmetry $p_{\mu} \leftrightarrow \tilde{x}_{\mu}$ called Langmann-Szabo symmetry and is renormalizable at every order in λ !

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- parametric representation and dimensional renormalization

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Noncommutative Field Theory

Random matrices

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Among other aspects, the GW model is a theory of non independent, non identically distributed random matrices for which one can understand the large N limit.

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Noncommutative Field Theory



$\phi_4^{\star 4}$ in the Matrix Base

In the matrix base, the action is

$$S = (2\pi\theta)^2 \sum_{m,n,k,l} \left\{ \frac{1}{2} \phi_{mn} \Delta_{mn;kl} \phi_{kl} + \frac{\lambda}{4!} \phi_{mn} \phi_{nk} \phi_{kl} \phi_{lm} \right\}$$

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Recall the interaction is very simple namely $Tr\phi^4$. But the propagator is not.

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Noncommutative Field Theory

The Propagator at $\Omega \neq 1$

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Computation of he propagator at $\Omega \neq 1$ was done by Grosse and Wulkenhaar. It depends on 3 indices, not four. More precisely one finds that G(m, n; k, l is zero unless n - m = k - l = h. It writes

$$G_{m,m+h;l+h,l} = \frac{\theta}{8\Omega} \int_0^1 d\alpha \, \frac{(1-\alpha)^{\frac{\mu_0^2 \theta}{8\Omega}}}{(1+C\alpha)^2} \left(\frac{\sqrt{1-\alpha}}{1+C\alpha}\right)^{m+l+h} \\ \times \sum_{u=\max(0,-h)}^{\min(m,l)} \mathcal{A}(m,l,h,u) \, \left(\frac{C\alpha(1+\Omega)}{\sqrt{1-\alpha}(1-\Omega)}\right)^{m+l-2u} + \frac{1}{2} \left(\frac{C\alpha(1+\Omega)}{\sqrt{1-\alpha}$$

with
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Noncommutative Field Theory

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It is a Gaussian distribution for random matrices which is independent, non-identically distributed.

The propagator in *x*-space

The propagator in x-space

The propagator for this theory is best understood through its parametric representation. In dimension d:

$$G(x,y) = \int_0^\infty dt \ e^{-\frac{\mu_0^2 \theta}{4\Omega}t} \\ \frac{\Omega}{(2\pi \sinh \Omega t)^{d/2}} \exp\left(-\frac{\Omega \cosh \Omega t}{2\sinh \Omega t} (\|x\|^2 + \|y\|^2) + \frac{\Omega x \cdot y}{\sinh \Omega t}\right).$$

The propagator in *x*-space

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and involves the Mehler kernel rather than the heat kernel.

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The Moyal vertex

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It can be also computed explicitly in direct space. V is proportional to

$$\int \prod_{i=1}^{4} d^{4} x^{i} \phi(x^{i}) \, \delta(x_{1} - x_{2} + x_{3} - x_{4}) \, \exp\left(2 \imath \theta^{-1} \left(x_{1} \wedge x_{2} + x_{3} \wedge x_{4}\right)\right)$$

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It relies again on the combination of three elements, but they are

A new scale decomposition

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- A new locality principle (Moyality)

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As in the commutative case the first two elements are quite universal. The third depends on details of the model, such as dimension and interaction.
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Scale decomposition

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- In direct space

$$G^{i}(x,y) = \int_{M^{-2i}}^{M^{-2(i-1)}} dt \cdots \leqslant KM^{2i} e^{-c_{1}M^{2i} ||x-y||^{2} - c_{2}M^{-2i} (||x||^{2} + ||y||^{2})}$$

The corresponding new renormalization group corresponds to a completely new mixture of the previous ultraviolet and infrared notions. Furthermore there exists only a half direction which is infinite for this RG.

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Moyality

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This principle applies only to planar graphs with a single external face.

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Power Counting

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$$\omega = \frac{d}{2}(F - EF) - L = \left(2 - \frac{E}{2}\right) - 4g - 2(EF - 1)$$
 if $d = 4$.

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Unexpected: No Landau ghost, so complete constructive analysis (with loop vertices exansion) should be possible!