

HYDRODYNAMIC LIMIT FOR AN ACTIVATED EXCLUSION PROCESS

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ABSTRACT. We study the hydrodynamic limit for a periodic 1-dimensional exclusion process with a dynamical constraint, which prevents a particle at site x from jumping to site $x \pm 1$ unless site $x \mp 1$ is occupied. This process with degenerate jump rates admits transient states, which it eventually leaves to reach an ergodic component, assuming that the initial macroscopic density is larger than $\frac{1}{2}$, or one of its absorbing states if this is not the case. It belongs to the class of conserved lattice gases (CLG) which have been introduced in the physics literature as systems with active-absorbing phase transition in the presence of a conserved field. We show that, for initial profiles smooth enough and uniformly larger than the critical density $\frac{1}{2}$, the macroscopic density profile for our dynamics evolves under the diffusive time scaling according to a fast diffusion equation (FDE). The first step in the proof is to show that the system typically reaches an ergodic component in subdiffusive time.

1. INTRODUCTION

The procedure of deriving the partial differential equation (PDE) ruling the macroscopic density profile of a microscopic stochastic dynamics under some suitable space-time rescaling is referred to in the probability community as *hydrodynamic limit* (cf. [11]). There is a vast literature on the hydrodynamic limit for exclusion processes, namely systems of particles interacting on a lattice, where only one particle per site is allowed. The subject has been receiving much attention for decades since such models are simple enough for rigorous mathematical study but also rich enough to describe many interesting phenomena. In this article, we derive the hydrodynamic limit for a 1-dimensional exclusion model with dynamical constraints, which we refer to as *activated exclusion process* (AEP)¹, in which a particle present at site x jumps at rate one to site $x \pm 1$ assuming that site $x \mp 1$ is occupied by a particle.

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¹This system is also called *restricted* [2] or *facilitated exclusion process* in the literature [1, 7]. We prefer to relate this process with activated random walks (which display a similar active-absorbing transition) rather than facilitated or kinetically constrained spin models, in which the constraint is typically chosen to make the dynamics reversible.

This model was originally introduced as a system with active-absorbing phase transition in the presence of a conserved field [15]. Such models are referred to as conserved lattice gases (CLG) in the context of the study of non-equilibrium phase transition. They usually exhibit a phase transition from an absorbing phase to an active state. In the AEP for instance, if the density is bigger than $\frac{1}{2}$ the system is in the active state with a unique invariant measure, while all the invariant measures are superpositions of atoms on absorbing states if the density is less than $\frac{1}{2}$. The critical particle density is therefore $\frac{1}{2}$. The critical behavior of CLG has been studied numerically and analytically in the physics literature [12, 14, 2], in order to identify the universality classes of models displaying an active-absorbing phase transition. In particular, the AEP is not in the same universality class as directed percolation. Let us mention that, recently, the AEP appeared in its totally asymmetric version in [1], where the authors compare the behavior of its first particles with what happens in the totally asymmetric exclusion process (TASEP).

Our main result is that, in the active phase, the macroscopic behavior of this microscopic dynamics, under periodic boundary conditions, is ruled by a non-linear diffusion equation.

The derivation of a non-linear diffusion equation from general reversible exclusion processes was first established by Funaki et al. [5] for gradient models, and by Varadhan and Yau [18] for non-gradient models. In both works, the irreducibility of the process on the hyperplanes with fixed number of particles is one of the essential assumptions. Results are few without this irreducibility. For instance, degenerate dynamics have been studied as microscopic models of the 1-dimensional porous medium equation (PME)

$$\partial_t \rho = \partial_u (\rho^{m-1} \partial_u \rho) \quad (m \in \mathbb{N}, m > 1) \quad (1.1)$$

in [8], [16] (for a generalization of the PME) and more recently in [3]. In these models, the problem comes mostly from the fact that regions with trivial density do not dissolve instantaneously, which calls for some adaptation and new techniques to apply the classical entropy method or relative entropy method. The obstacles caused by the lack of irreducibility that we address in this paper are of a different nature.

First, in our dynamics the state space is divided into transient states, absorbing states and ergodic states. Depending on the initial number of particles in the system, some of configurations will lead to the ergodic component (we call such states “transient good”) and the others will be absorbed to an inactive state (resp. “transient bad”). Although the behavior of the dynamics started from a transient bad configuration is an interesting matter in its own right, in this article, we focus on the transient good case where the ergodic component is ultimately reached. To apply the entropy method, the existence of transient good states is also troublesome since, although invariant measures are supported only on the ergodic components, the process may stay in the transient good states for some macroscopic time with positive probability. To guarantee this is not the case, we show Theorem 2.4. More precisely: assume that the initial density is initially larger than $\frac{1}{2}$ and look at the microscopic system of size N at a macroscopic time $(\log N)^\alpha / N^2$ for some $\alpha > 0$; then with high probability it has already reached the ergodic component. This is the first main novelty of this paper.

Secondly, the Gibbs measures (grand canonical measures) of our process are not product, while they are Bernoulli product measures in [8]. In particular, when the density is close to $\frac{1}{2}$, these measures exhibit long range correlations, and adapting the entropy method for these non-product grand canonical measures requires significant extra technical work. We note that both papers [5, 18] also cover cases with Gibbs measures which

are not necessarily product, although for non-degenerate dynamics. Besides exclusion processes, microscopic models with non-product Gibbs measures for which the hydrodynamic limit is rigorously established are very few.

As the nature of microscopic dynamics is different from [8], the macroscopic equation is also different from the PME, and we obtain the following macroscopic equation

$$\partial_t \rho = \Delta(-\rho^{-1}) = \partial_u(\rho^{-2} \partial_u \rho).$$

The equation is in the class of fast diffusion equation (FDE), which corresponds to the case with $m < 1$ in (1.1). Though the equation has a singularity at $\rho = 0$, it is not relevant for us since we always assume that the initial density is bigger than $\frac{1}{2}$ at any spatial point. Because of the phase transition described above, for density profiles below $\frac{1}{2}$ the system cannot be governed by this macroscopic equation, since it does not reflect the absorption phenomenon. On the other hand, an extension to initial profiles with total density bigger than $\frac{1}{2}$ (but locally smaller) with a different macroscopic equation might be possible, but it remains out of reach. It would require understanding the interplay between frozen regions with density lower than $\frac{1}{2}$ and active regions with higher density, which may give rise to complex phenomena [4]. We note that the same FDE was derived recently from a non-degenerate zero-range process under a proper high density limit [9].

We now give some remarks on other aspects of our model. As discussed in the next section, the model is reversible and gradient. Moreover, it can be *formally* understood as the specific case of Examples 1 or 2 in [5, Section 5], specifically the case with parameters $b(1) = -\infty$, $b(k) = 0$ for $k \geq 2$ in Example 1, or $\alpha = \infty$ and $\beta = 0$ in Example 2 under proper normalization. In particular, the exclusion processes in the class described in Example 1 of [5] can be mapped to zero-range processes by a simple but non-linear transformation of configurations, which we discuss precisely in Section 3. This relation has been used in the literature (cf. [6, 2, 10]) and plays an essential role in our estimation of the transience time. Note that since the corresponding zero-range process is also degenerate, we cannot apply classical results (cf. [11, Section 5]) to derive its hydrodynamic limit, therefore our main theorem is not proved that way. It is also not straightforward to deduce a hydrodynamic limit for the AEP from the zero-range through the relation mentioned above.

Here follows an outline of the paper. We start in Section 2 by introducing notations, defining the model and stating our main results. Section 3 is devoted to mapping our exclusion process to a zero-range process. In Section 4 we use this transformation and we prove Theorem 2.4, which states that the system reaches its ergodic component in a sub-diffusive time scale. In Section 5 we give the rigorous proof of the hydrodynamic limit (Theorem 2.2) via a non-trivial adaptation of the entropy method, which involves two main ingredients: first, the *equivalence of ensembles*, proved in Section 6 (which also introduces the different reference measures we consider), and second, the *Replacement Lemma*, proved in Section 7.

2. MODEL AND RESULTS

2.1. Notations and conventions. Let us introduce notations and conventions that we use throughout the paper.

- N is an integer which plays the role of a scaling parameter and will go to infinity.
- For any finite set Λ we denote by $|\Lambda|$ its cardinality and by Λ^c its complement.
- We let $\mathbb{T}_N := \mathbb{Z}/N\mathbb{Z} = \{1, \dots, N\}$ be the discrete torus of size N , and $\mathbb{T} = [0, 1)$ be the 1-dimensional continuous torus.

- For any $\ell \in \mathbb{N}$ we set $B_\ell := \{-\ell, \dots, \ell\}$ the centered symmetric box of size $2\ell + 1$, which can be seen as either a subset of \mathbb{T}_N (if $2\ell + 1 \leq N$), or a subset of \mathbb{Z} . Similarly, we set $\Lambda_\ell := \{1, \dots, \ell\}$.
- We will consider particle systems on different state spaces $\{0, 1\}^E$, with E either the full discrete line \mathbb{Z} , or the discrete torus \mathbb{T}_N , or a finite box in \mathbb{Z} . To avoid any confusion, the elements of $\{0, 1\}^E$, which are named *configurations*, will be denoted: by ξ if $E = \mathbb{Z}$, by η if $E = \mathbb{T}_N$, and by σ, ς if E is a finite box. We will use the letters $\omega, \chi, \zeta, \Upsilon$ to denote configurations in yet different spaces.
- For any $x \in \mathbb{T}_N$ and configuration $\eta \in \{0, 1\}^{\mathbb{T}_N}$, we denote by $\eta(x) \in \{0, 1\}$ the occupation variable at site x , namely: $\eta(x) = 1$ if there is a particle at site x , and 0 otherwise. We treat similarly the configurations in different state spaces.
- For any measurable function $f : \{0, 1\}^{\mathbb{T}_N} \rightarrow \mathbb{R}$, and $x \in \mathbb{T}_N$, we denote by $\tau_x f$ the function obtained by translation as follows: $\tau_x f(\eta) := f(\tau_x \eta)$, where $(\tau_x \eta)(y) = \eta(x + y)$, for $y \in \mathbb{T}_N$. We treat similarly the other state spaces.
- For any $\Lambda \subset \mathbb{T}_N$ or $\Lambda \subset \mathbb{Z}$ and for any probability measure π on $\{0, 1\}^\Lambda$, we adopt the following notations:
 - (i) the configuration $\eta \in \{0, 1\}^{\mathbb{T}_N}$ (resp. $\xi \in \{0, 1\}^{\mathbb{Z}}$) restricted to Λ is denoted by $\eta|_\Lambda$ (resp. $\xi|_\Lambda$),
 - (ii) if $f : \{0, 1\}^\Lambda \rightarrow \mathbb{R}$ is a measurable function, $\pi(f)$ denotes the expectation of f w.r.t. the measure π ,
 - (iii) if $A \subset \{0, 1\}^\Lambda$ is a subset of all possible configurations, $\pi(\eta \in A)$ equivalently means $\pi(\mathbf{1}_{\{\eta \in A\}}) := \pi(A)$.
- For any sequence $(u_k)_{k \in \mathbb{N}}$, possibly depending on other parameters than the index k , we will denote $\mathcal{O}_k(u_k)$ (resp. $o_k(u_k)$) an arbitrary sequence $(v_k)_{k \in \mathbb{N}}$ such that there exists a constant $C > 0$ (resp. a vanishing sequence $(\varepsilon_k)_{k \in \mathbb{N}}$) – possibly depending on the other parameters – such that

$$v_k \leq C u_k \quad (\text{resp. } v_k \leq u_k \varepsilon_k) \quad \forall k \in \mathbb{N}.$$

2.2. The microscopic dynamics. The dynamics is as follows: on the periodic domain \mathbb{T}_N , we associate with each site a random Poissonian clock. When the clock at site x rings, if there is a particle sitting at this site x , that particle jumps to $x - 1$ or $x + 1$ if some local constraint is satisfied: a particle can jump to the right (resp. left) only if it has a particle to its left (resp. right).

More precisely, the infinitesimal generator ruling the evolution in time of this Markov process is given by \mathcal{L}_N , which acts on functions $f : \{0, 1\}^{\mathbb{T}_N} \rightarrow \mathbb{R}$ as

$$\mathcal{L}_N f(\eta) := \sum_{x \in \mathbb{T}_N} c_{x,x+1}(\eta) (f(\eta^{x,x+1}) - f(\eta)), \quad (2.1)$$

where the constraint and the exclusion rule are encoded in the rates $c_{x,x+1}$ as

$$c_{x,x+1}(\eta) = \eta(x-1)\eta(x)(1 - \eta(x+1)) + \eta(x+2)\eta(x+1)(1 - \eta(x)), \quad (2.2)$$

and $\eta^{x,y}$ denotes the configuration obtained from η by exchanging the states of sites x, y , namely $\eta^{x,y}(x) = \eta(y)$, $\eta^{x,y}(y) = \eta(x)$ and $\eta^{x,y}(z) = \eta(z)$ if $z \neq x, y$. Figure 2.1 below shows examples of jumps. Note that the dynamics conserves the total number of particles $\sum_{x \in \mathbb{T}_N} \eta(x)$.

This activated exclusion dynamics is *degenerate*, *gradient*, and *reversible*, as explained in the following three paragraphs.

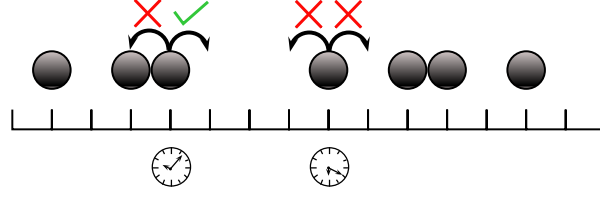


FIGURE 2.1. Allowed jumps are denoted by \checkmark . Forbidden jumps are denoted by \times .

2.2.1. *Degenerate dynamics.* Because the jump rates can vanish, the dynamics will be referred to as *degenerate*. For this reason, it is convenient to refine the classification of the configurations into transient/recurrent states as follows (this classification is fully justified in Section 3). Let \mathcal{H}_N^k be the hyperplane of configurations with k particles, with $k \in \{0, \dots, N\}$, namely:

$$\mathcal{H}_N^k := \left\{ \eta \in \{0, 1\}^{\mathbb{T}_N} : \sum_{x \in \mathbb{T}_N} \eta(x) = k \right\}.$$

Then,

- (1) if $k \leq \frac{N}{2}$, some configurations in \mathcal{H}_N^k are *blocked*, in the sense that no particle can jump because the constraint (2.2) is satisfied nowhere (they are *absorbing states* for the dynamics). Those are exactly the configurations in which all particles are isolated.

The other configurations in \mathcal{H}_N^k are not blocked but, starting from them, with probability one the process will arrive at a blocked configuration in a finite number of steps. We call them *transient bad* configurations. See Figure 2.2 below.

- (2) if $k > \frac{N}{2}$, the process will never reach an absorbing state (except in the trivial case $k = N$). In \mathcal{H}_N^k , there are configurations which are in the ergodic component (the recurrent states for the process, which in this case form an irreducible component); we call them *ergodic* configurations. They are the configurations in which empty sites are isolated.

Starting from the other configurations, which are called *transient good* configurations, the process enters the ergodic component after a finite number of steps a.s. See Figure 2.2 below.

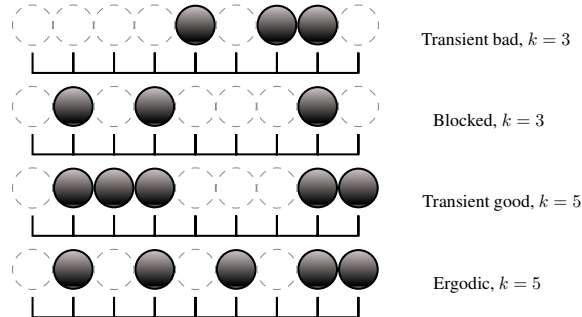


FIGURE 2.2. Four examples of configurations which belong to the different classes, for $N = 9$ sites.

Those observations (and similar ones for configurations in finite boxes) lead to the following definitions.

Definition 2.1. We denote by $\mathcal{E}_N \subset \{0, 1\}^{\mathbb{T}_N}$ the set of ergodic configurations on \mathbb{T}_N , namely

$$\mathcal{E}_N := \left\{ \eta \in \{0, 1\}^{\mathbb{T}_N} : \forall x \in \mathbb{T}_N, (\eta(x), \eta(x+1)) \neq (0, 0) \text{ and } \sum_{x \in \mathbb{T}_N} \eta(x) > \frac{N}{2} \right\}, \quad (2.3)$$

and by $\widehat{\mathcal{E}}_\Lambda$ the set of ergodic local configurations (which are actually restrictions of ergodic configurations) on a (finite) connected set $\Lambda \subset \mathbb{Z}$, namely

$$\widehat{\mathcal{E}}_\Lambda := \left\{ \sigma \in \{0, 1\}^\Lambda : \forall (x, x+1) \in \Lambda^2, (\sigma(x), \sigma(x+1)) \neq (0, 0) \right\}. \quad (2.4)$$

For $k > \frac{N}{2}$ we also let

$$\Omega_N^k := \mathcal{H}_N^k \cap \mathcal{E}_N, \quad (2.5)$$

be the set of ergodic configurations on \mathbb{T}_N which contain exactly k particles.

We conclude this paragraph with the following result:

Lemma 2.1. For any $N \geq 1$ and any $k > \frac{N}{2}$, Ω_N^k is an irreducible component for our Markov process.

Proof. Let us change the point of view and move the zeros around instead of the particles. Then it is clear that in a configuration in \mathcal{E}_N , a zero can jump as long as it remains at distance at least two from the others, and every allowed jump is reversible. Consequently, it is enough to show that from every configuration in Ω_N^k , one can reach the configuration $\circ \bullet \circ \bullet \dots \bullet \bullet$ where the $N - k$ zeros start alternating with particles until there are none left. Let $\eta \in \Omega_N^k$ and number its zeros from left to right (the “left-most” site being $1 \in \mathbb{T}_N$). In that order, pull each of them as much to the left as possible. That way we reach either the desired configuration, or the same shifted one step to the right (in case $\eta(N) = 0$): $\bullet \circ \bullet \circ \bullet \dots \bullet$. In the second case, iterating the process brings us back to the desired configuration. \square

2.2.2. Gradient system. We introduce the *instantaneous currents*, which are defined for any configuration η and any site x as

$$j_{x,x+1}(\eta) = c_{x,x+1}(\eta)(\eta(x+1) - \eta(x)),$$

and satisfy $\mathcal{L}_N(\eta(x)) = j_{x-1,x}(\eta) - j_{x,x+1}(\eta)$. One can easily check that our model is *gradient*, which is to say that these currents can be written as discrete gradients

$$j_{x,x+1}(\eta) = \tau_x h(\eta) - \tau_{x+1} h(\eta), \quad \text{for any } x \in \mathbb{T}_N,$$

with the function h given by

$$h(\eta) := \eta(-1)\eta(0) + \eta(0)\eta(1) - \eta(-1)\eta(0)\eta(1). \quad (2.6)$$

This function h plays a fundamental role in the derivation of the hydrodynamic limit of our process.

2.2.3. *Reversible measures.* The uniform measures on Ω_N^k , denoted below by π_N^k (with $k \in \{0, \dots, N\}$) are invariant for the Markov process induced by the infinitesimal generator \mathcal{L}_N and satisfy the *detailed balance condition*, as detailed in (6.1). When $k/N \rightarrow \rho \in (\frac{1}{2}, 1)$, these measures locally converge to an infinite volume grand canonical measure π_ρ on $\{0, 1\}^{\mathbb{Z}}$ for which an explicit formula can be derived. The measures π_ρ are not product, and all relevant canonical and grand canonical measures will be thoroughly investigated in Section 6.

2.3. **Main results.** We are now ready to state the main result of this article. Fix an initial smooth profile $\rho_0 : \mathbb{T} \rightarrow (\frac{1}{2}, 1]$, and consider the non-homogeneous product measure on $\{0, 1\}^{\mathbb{T}^N}$ fitting ρ_0 , defined as

$$\mu_N(\eta) := \prod_{k \in \mathbb{T}^N} \left(\rho_0\left(\frac{k}{N}\right) \eta(k) + (1 - \rho_0\left(\frac{k}{N}\right)) (1 - \eta(k)) \right). \quad (2.7)$$

Let $\{\eta_t : t \geq 0\}$ denote the Markov process driven by the *accelerated* infinitesimal generator $N^2 \mathcal{L}_N$ (cf. (2.1)) starting from the initial measure μ_N . Fix $T > 0$ and denote by \mathbb{P}_{μ_N} the probability measure on the Skorokhod path space $\mathcal{D}([0, T], \{0, 1\}^{\mathbb{T}^N})$ corresponding to this dynamics. We denote by \mathbb{E}_{μ_N} the corresponding expectation. Note that, even though it is not explicit in the notation, \mathbb{P} , \mathbb{E} and η_t strongly depend on N : through the size of the state space but also through the diffusive time scaling.

Theorem 2.2 (Hydrodynamic limit). *For any $t \in [0, T]$, any $\delta > 0$ and any continuous test function $\varphi : \mathbb{T} \rightarrow \mathbb{R}$, we have*

$$\lim_{N \rightarrow \infty} \mathbb{P}_{\mu_N} \left[\left| \frac{1}{N} \sum_{x \in \mathbb{T}^N} \varphi\left(\frac{x}{N}\right) \eta_t(x) - \int_{\mathbb{T}} \varphi(u) \rho(t, u) du \right| > \delta \right] = 0 \quad (2.8)$$

where $\rho(t, u)$ is the unique smooth solution of the hydrodynamic equation

$$\partial_t \rho = \Delta \left(\frac{2\rho - 1}{\rho} \right), \quad \rho(0, \cdot) = \rho_0(\cdot) : \mathbb{T} \rightarrow (\tfrac{1}{2}, 1]. \quad (2.9)$$

Remark 2.3. *The fact that there is a unique smooth solution to (2.9), provided that the initial profile satisfies $\rho_0 > \frac{1}{2}$, is quite standard in PDE theory, and can be found for instance in [19, Section 3.1]. The equation (2.9) belongs to the family of quasilinear parabolic problems of the form*

$$\partial_t \rho = \partial_u (A(\rho, \partial_u \rho)),$$

where, in our case, $A(p, q) := q/p^2$. The classical theory, which gives smoothness of solutions and maximum principles, cannot be used for equation (2.9), because the latter is not uniformly parabolic². However, the classical results can be used if the initial condition is non-degenerate. For instance, if ρ_0 satisfies $\frac{1}{2} + \varepsilon \leq \rho_0 \leq 1$, then one can apply the usual theory with such data, choosing $A(p, q) = q/p^2$ if $\frac{1}{2} + \varepsilon \leq p \leq 1$, and extending the function A as a linear function of p and q for p near $\frac{1}{2}$ or 1, making a smooth connection at the point $p = \frac{1}{2} + \varepsilon$. With this choice, the degeneracy has been eliminated, and therefore there exists a unique classical (smooth) solution, which satisfies the same bounds $\frac{1}{2} + \varepsilon \leq \rho \leq 1$. This means that ρ actually never takes values in the critical region, and therefore one gets a classical solution for (2.9).

²We say that A is uniformly parabolic if there exist constants $0 < c_1 < c_2 < \infty$ such that $c_1 \leq \frac{\partial A}{\partial q}(p, q) \leq c_2$, uniformly in (p, q) .

There are two main difficulties to prove Theorem 2.2. The first one lies in the fact that the dynamics is *a priori* not ergodic. We now state our second main result, which will be proved in Section 4. It states that the accelerated system reaches its ergodic component at a macroscopic time t_N of order $(\log N)^\alpha/N^2$ for some $\alpha > 0$. Therefore, for any macroscopic time $t > 0$ and for any N large enough such that $t_N < t$, the configuration η_t belongs to the ergodic component with very high probability. This is the one of the main novelties of this work.

Theorem 2.4 (Transience time for the exclusion process with absorption). *Letting $\ell_N = (\log N)^8$ and $t_N = \ell_N^4/N^2$, we have*

$$\lim_{N \rightarrow \infty} \mathbb{P}_{\mu_N} \left(\eta_{t_N} \notin \mathcal{E}_N \right) = 0.$$

The second main difficulty to prove Theorem 2.2 comes from the nature of the invariant measures of the process: we investigate in Section 6 the canonical and grand canonical measures for the generator \mathcal{L}_N , which only charge the ergodic component \mathcal{E}_N where all empty sites are isolated. These measures are therefore not product, and significant work is required to show the various properties that are needed to apply the classical *entropy method*.

The rest of the article is organised as follows. We describe in Section 3 a correspondence between our activated exclusion dynamics and a particular zero-range dynamics, which will be crucial to prove Theorem 2.4 in Section 4. We then give in Section 5 a general outline of the classical strategy to prove Theorem 2.2. We thoroughly study in Section 6 the invariant measures for the activated exclusion dynamics, and use the properties we derive to prove the Replacement Lemma in Section 7.

3. CORRESPONDENCE WITH THE ZERO-RANGE MODEL

Following [2], our exclusion-type dynamics can be mapped to a *zero-range model* in a way which we describe below.

First, we map the initial configuration of the exclusion dynamics (EX) to a zero-range (ZR) configuration on a reduced torus: with $\eta \in \mathcal{H}_N^k$, $k < N$, we associate a configuration $\omega \in \mathbb{N}^{\mathbb{T}_{N-k}}$ as follows. Look for the first empty site to the left of or at site 1 and label it 1. Moving to the right, label all empty sites in η from 1 to $N-k$. Then define $\omega(i)$ as the number of particles between the i -th and $(i+1)$ -th (with $N-k+1$ identified with 1) empty sites in η (see Figure 3.1). For any exclusion configuration η , let us denote by $\Pi(\eta)$ the corresponding zero-range configuration. Note that this mapping is not one-to-one. For instance, in the example considered in Figure 3.1, shifting η one site to the left does not change ω .

Now we couple the EX process $\{\eta_t : t \geq 0\}$ (generated by $N^2 \mathcal{L}_N$) with a ZR process $\{\omega_t : t \geq 0\}$. With η_0 , let us associate $\omega_0 := \Pi(\eta_0)$ in the way described above. Then whenever a particle jumps in the process η_t , a particle in the corresponding pile in ω_t jumps in the same direction (as shown in Figure 3.1)³. In particular, a jump of a particle to an empty site is allowed in the EX process iff the corresponding pile in the ZR process has at

³Note that, in order to make this informal description rigorous, we need to keep track of the correspondence between blocks of particles in the EX process and piles in the ZR process: it is not true that ω_t is obtained from η_t through the same mapping Π as ω_0 from η_0 (this can be checked in the example of Figure 3.1). One way of understanding the process is to tag the empty site labeled 1 (*i.e.* when a particle jumps there, the label moves to the site that particle previously occupied). Then the mapping from η_t to ω_t is the one described above, with the exception that the empty site with label 1 need not be the first to the left of the origin.

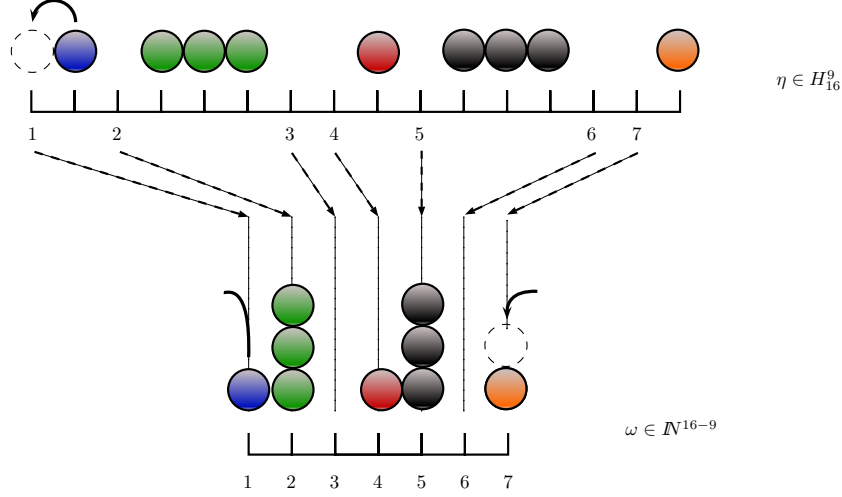


FIGURE 3.1. TOP: configuration in the exclusion dynamics, with $N = 16$ sites and $k = 9$ particles. BOTTOM: the associated zero-range configuration. The arrows depict a jump in the EX process together with its effect on the ZR process; the dashed circles represent the positions of the blue particles after the jump.

least two particles. Then, one can easily check that $\{\omega_t : t \geq 0\}$ is a Markov process with infinitesimal generator $N^2 \mathcal{L}_{N-k}^{ZR}$, with \mathcal{L}_{N-k}^{ZR} acting on functions $f : \mathbb{N}^{\mathbb{T}_{N-k}} \rightarrow \mathbb{R}$ as follows:

$$\mathcal{L}_{N-k}^{ZR} f(\omega) := \sum_{x \in \mathbb{T}_{N-k}} \sum_{z=\pm 1} \mathbf{1}_{\{\omega(x) \geq 2\}} (f(\omega^{x,x+z}) - f(\omega)), \quad (3.1)$$

where

$$\omega^{x,y}(z) := \begin{cases} \omega(x) - 1; & z = x, \\ \omega(y) + 1; & z = y, \\ \omega(z); & z \notin \{x, y\}. \end{cases}$$

Let us classify the possible states of this ZR process as we have already done in Section 2.2.1 for the EX process. In this setting, it is fairly clear that for $\omega \in \mathbb{N}^{\mathbb{T}_{N-k}}$ with k particles,

- (1) if $k \leq N - k$ (there are fewer particles than sites), either $\omega(x) \leq 1$ for all x (in which case the configuration is blocked), or this situation will be reached after a finite number of jumps a.s.
- (2) if $k > N - k$, by the pigeonhole principle there will always be at least one particle allowed to jump. The ergodic configurations are those where $\omega(x) \geq 1$ for all x , and the transient good the other ones, where there exists a site x such that $\omega(x) = 0$.

This translates immediately into the classification we claimed for the EX configurations in Section 2.2.1⁴. In particular we will use in the proof of Theorem 2.4 the fact that for all $t \geq 0$, η_t is in its ergodic component iff ω_t is in its ergodic component.

⁴Note that, even though the reverse mapping from ZR to EX is only defined up to the position of the empty site with label 1, since the properties we consider are translation invariant we can safely transfer them from one setting to the other.

4. PROOF OF THEOREM 2.4: TRANSCIENCE TIME TO REACH THE ERGODIC COMPONENT

In this section, we prove that after a time of order $(\log N)^\alpha/N^2$, with high probability, the exclusion process has reached the ergodic component. This result relies strongly on the correspondence between our activated exclusion process and the zero-range dynamics presented in Section 3. In Section 4.1, we state the main result of this section (Proposition 4.1), which is an estimate of the transience time for the zero-range process assuming that one starts from what we call a *regular* configuration, and use it to prove Theorem 2.4. Section 4.2 is dedicated to proving Proposition 4.1. Finally, in Section 4.3, we prove a technical lemma that states that, starting the exclusion process from the smooth product measure (2.7), the probability for the corresponding zero-range configuration to be regular is close to 1.

4.1. Ergodic component of the zero-range dynamics. Let us fix an integer $\ell \geq 0$, and set $\bar{\Lambda}_\ell := \Lambda_\ell \cup \{0, \ell + 1\} = \{0, \dots, \ell + 1\}$. For any given integer $K > \ell$, any zero-range configuration $\omega \in \mathbb{N}^{\mathbb{T}_K}$ on \mathbb{T}_K , and any $x \in \mathbb{T}_K$, we define the zero-range configuration $\bar{\omega}^{\ell, x} \in \mathbb{N}^{\bar{\Lambda}_\ell}$ as

$$\bar{\omega}^{\ell, x}(y) := \begin{cases} \omega(x+y) & \text{if } 1 \leq y \leq \ell, \\ 0 & \text{if } y = 0, \ell + 1. \end{cases} \quad (4.1)$$

which is the configuration in the box of size ℓ to the right of x in \mathbb{T}_K , with empty sites at the boundaries. We also let

$$n_\ell(\omega) := \sum_{y=1}^{\ell} \omega(y) \quad (4.2)$$

the initial number of particles in Λ_ℓ , and let

$$Z_\ell(\omega) := \sum_{y=1}^{\ell} y\omega(y) \quad (4.3)$$

the cumulated distance of the particles in $\bar{\Lambda}_\ell$ to the first site in the configuration ω .

Fix $\delta \in (0, 1)$. Throughout the proof, we will when convenient omit the integer part $\lfloor \cdot \rfloor$, and for example simply write $(1 + \delta)\ell$ instead of $\lfloor (1 + \delta)\ell \rfloor$.

For any configuration $\omega \in \mathbb{N}^{\mathbb{T}_K}$ and $x \in \mathbb{T}_K$, we denote by $\omega^{\ell, x}$ the configuration obtained from $\bar{\omega}^{\ell, x}$ by keeping at most the $(1 + \delta)\ell$ left-most particles and destroying all the other particles. More precisely, let us denote by $1 \leq x_1 \leq x_2 \leq \dots \leq x_{n_\ell(\bar{\omega}^{\ell, x})} \leq \ell$ the positions of the particles in $\bar{\omega}^{\ell, x}$, we then have

$$\bar{\omega}^{\ell, x}(y) = \sum_{k=1}^{n_\ell(\bar{\omega}^{\ell, x})} \mathbf{1}_{\{x_k=y\}},$$

and we let

$$\omega^{\ell, x}(y) = \sum_{k=1}^{\min\{(1+\delta)\ell, n_\ell(\bar{\omega}^{\ell, x})\}} \mathbf{1}_{\{x_k=y\}}. \quad (4.4)$$

We denote by

$$A_\ell := A_\ell(\delta) = \left\{ \omega \in \mathbb{N}^{\bar{\Lambda}_\ell} : \omega(0) = \omega(\ell + 1) = 0, \right.$$

$$n_\ell(\omega) = (1 + \delta)\ell \quad \text{and} \quad Z_\ell(\omega) \leq \left(1 + \frac{\delta}{2}\right)(1 + \delta) \frac{\ell(\ell + 1)}{2}, \quad (4.5)$$

the set of configurations on $\bar{\Lambda}_\ell$ for which Z_ℓ is not much bigger than it would be if the $(1 + \delta)\ell$ particles were spread evenly across Λ_ℓ . We also set, for any $K \geq \ell$

$$B_{K,\ell} := B_{K,\ell}(\delta) = \left\{ \omega \in \mathbb{N}^{\mathbb{T}^K} : \omega^{\ell,x} \in A_\ell \quad \forall x \in \mathbb{T}_K \right\}. \quad (4.6)$$

We call δ -regular configurations the elements of $B_{K,\ell}(\delta)$. In other words, since in $\omega^{\ell,x}$ we delete all but $(1 + \delta)\ell$ particles, $B_{K,\ell}$ is the set of zero-range configurations on \mathbb{T}_K such that in each box $\{x + 1, \dots, x + \ell\}$, there are at least $(1 + \delta)\ell$ particles, and those particles are not placed abnormally to the right of $\{x + 1, \dots, x + \ell\}$.

For a given initial configuration $\omega_0 \in \mathbb{N}^{\mathbb{T}^K}$, we denote by $\mathbb{Q}_{K,\omega_0}^{ZR}$ the distribution of the process $\{\omega_t : t \geq 0\}$ on the space of trajectories $\mathcal{D}([0, T], \mathbb{N}^{\mathbb{T}^K})$, starting from ω_0 and driven by the generator $N^2 \mathcal{L}_K^{ZR}$ introduced in (3.1). We start by stating that, assuming that the initial configuration ω_0 is δ -regular (i.e. is in $B_{K,\ell}(\delta)$), then the probability that the zero-range process starting from ω_0 has not reached the ergodic component at a time of order ℓ^4 is exponentially small.

Proposition 4.1 (Transience time for the zero-range process started from $B_{K,\ell}$). *For any integer $\ell \geq 0$, we define $T_\ell = \ell^4$. For any $\varepsilon < \frac{1}{2}$, and any constant $\delta > 0$, there exists $\bar{\ell} = \bar{\ell}(\varepsilon, \delta)$, such that for any integer $K > \ell \geq \bar{\ell}$, and any $N \in \mathbb{N}$*

$$\sup_{\omega_0 \in B_{K,\ell}(\delta)} \mathbb{Q}_{K,\omega_0}^{ZR} \left(\exists x \in \mathbb{T}_K, \omega_{T_\ell/N^2}(x) = 0 \right) \leq K \exp(-\ell^\varepsilon).$$

We now state that, while pairing the exclusion process and the zero-range process, if η_0 is started close to a smooth product measure, then with high probability the associated zero-range configuration ω_0 is in $B_{K,\ell}$ for some well chosen K, ℓ . Recall that $\{\eta_t : t \geq 0\}$ denotes the Markov process driven by the infinitesimal generator $N^2 \mathcal{L}_N$, whose law is denoted by \mathbb{P}_{μ_N} , where μ_N is the initial measure for the process defined in (2.7). Further recall that for any exclusion configuration η , $\Pi(\eta)$ denotes the corresponding zero-range configuration. We also denote by $K(\eta) = \sum_{x \in \mathbb{T}_N} (1 - \eta(x))$ the number of empty sites in η , i.e. the size of the associated zero-range configuration $\Pi(\eta)$ (as explained in Section 3).

Lemma 4.2. *Assume that the initial density profile ρ_0 is not identically equal to 1. Then, letting $\ell_N = (\log N)^8$, there exists a constant $\delta > 0$ such that*

$$\liminf_{N \rightarrow \infty} \mu_N \left(K(\eta) > \ell_N \quad \text{and} \quad \Pi(\eta) \in B_{K(\eta), \ell_N}(\delta) \right) = 1.$$

We prove this last lemma in Section 4.3. Before proving both Proposition 4.1 and Lemma 4.2, however, we conclude the proof of Theorem 2.4.

Proof of Theorem 2.4. This theorem is a consequence of Proposition 4.1 and Lemma 4.2. First recall (from Section 3) that

$$\eta_t \notin \mathcal{E}_N \Leftrightarrow \exists x \in \mathbb{T}_{K(\eta_0)}, \Pi(\eta_0)_t(x) = 0.$$

Letting $\ell_N = (\log N)^8$, and $t_N = \ell_N^4/N^2$ as in Theorem 2.4, we now write according to Lemma 4.2,

$$\begin{aligned} & \mathbb{P}_{\mu_N} \left(\eta_{t_N} \notin \mathcal{E}_N \right) \\ &= \mathbb{P}_{\mu_N} \left(\exists x \in \mathbb{T}_{K(\eta_0)}, \Pi(\eta_0)_{t_N}(x) = 0 \right) \end{aligned}$$

$$= \mathbb{P}_{\mu_N} \left(\exists x \in \mathbb{T}_{K(\eta_0)}, \Pi(\eta_0)_{I_N}(x) = 0 \mid K(\eta_0) > \ell_N \text{ and } \Pi(\eta_0) \in B_{K(\eta_0), \ell_N} \right) + o_N(1).$$

According to Proposition 4.1, for any $\varepsilon < \frac{1}{2}$, and any $\ell_N \geq \bar{\ell}(\varepsilon, \delta)$, the probability above is less than

$$\mathbb{E}_{\mu_N} \left[K(\eta_0) \mid K(\eta_0) > \ell_N \text{ and } \Pi(\eta_0) \in B_{K(\eta_0), \ell_N} \right] \exp(-\ell_N^\varepsilon).$$

Since $K(\eta)$ is the number of empty sites in η , it is less than N , and we now obtain for any $\varepsilon < \frac{1}{2}$, and any N large enough such that $\ell_N > \bar{\ell}(\varepsilon, \delta)$,

$$\mathbb{P}_{\mu_N} \left(\eta_{I_N} \notin \mathcal{E}_N \right) \leq N \exp(-\ell_N^\varepsilon) + o_N(1).$$

Since $\ell_N = (\log N)^\delta$, fixing $\varepsilon \in (\frac{1}{8}, \frac{1}{2})$ concludes the proof. \square

4.2. Estimation of the transience time. We now prove Proposition 4.1. Throughout this subsection, we assume that $\delta > 0$ is a fixed constant and ℓ is a fixed positive integer. Moreover, in order to avoid heavy notations, we work throughout this section with *unrescaled processes*. Namely, let $\{\tilde{\omega}_t : t \geq 0\}$ be the Markov process started from ω_0 and generated by \mathcal{L}_K^{ZR} . A version of this process is given by $\tilde{\omega}_t = \omega_{t/N^2}$, and therefore, in order to prove Proposition 4.1, we are reduced to prove that:

$$\sup_{\omega_0 \in B_{K, \ell}(\delta)} \mathbb{Q}_{K, \omega_0}^{ZR} \left(\exists x \in \mathbb{T}_K, \tilde{\omega}_{T_\ell}(x) = 0 \right) \leq K \exp(-\ell^\varepsilon).$$

For the sake of clarity, the proof will be divided into several lemmas which we prove at the end of this section. The principle of the proof is the following: fix $x \in \mathbb{T}_K$. By waiting a long time (of order ℓ^4), we make sure that with high probability each particle initially present in $\{x+1, \dots, x+\ell\}$ will either

- have encountered an empty site in the box $\{x+1, \dots, x+\ell\}$, and gotten stuck there forever,
- or have left the box $\{x+1, \dots, x+\ell\}$ at some point.

At most ℓ particles fall into the first case, therefore, if the box initially contains at least $(1+\delta)\ell$ particles, $\delta\ell$ particles or more will have left the box $\Lambda^{\ell, x} := \{x+1, \dots, x+\ell\}$ at least once before $T_\ell = \ell^4$. In order for the site x to remain empty up until T_ℓ , all of these particles must have left by the other boundary $x+\ell$, and therefore, particles in $\Lambda^{\ell, x}$ would have performed an abnormally large number of steps to the right. The probability of the last case occurring decays faster than $\exp(-\ell^\varepsilon)$ for any $\varepsilon < \frac{1}{2}$.

Our purpose is now to make this argument rigorous. Let us first write that

$$\mathbb{Q}_{K, \omega_0}^{ZR} \left(\exists x \in \mathbb{T}_K, \tilde{\omega}_{T_\ell}(x) = 0 \right) \leq K \max_{x \in \mathbb{T}_K} \mathbb{Q}_{K, \omega_0}^{ZR} \left(\tilde{\omega}_{T_\ell}(x) = 0 \right). \quad (4.7)$$

We are going to prove that the probability on the right hand side is, for any x , for any $\varepsilon < \frac{1}{2}$ and any ℓ large enough, smaller than $\exp(-\ell^\varepsilon)$.

Recall that we introduced the domain $\bar{\Lambda}_\ell = \{0, \dots, \ell+1\}$, and recall definitions (4.5) and (4.6) of A_ℓ and $B_{K, \ell}$, respectively. Let us fix $x \in \mathbb{T}_K$ and $\omega_0 \in B_{K, \ell}$. We introduce an auxiliary process $\{\chi_t(y)_{y \in \bar{\Lambda}_\ell} : t \geq 0\}$ (which will depend on ℓ and x). This process is started from $\omega_0^{\ell, x}$ (defined above in (4.4)), and therefore, since $\omega_0 \in B_{K, \ell}$, we know that $\chi_0 = \omega_0^{\ell, x} \in A_\ell$ contains exactly $(1+\delta)\ell$ particles located in Λ_ℓ . The process χ_t is driven

by the (stuck) zero-range generator of $\tilde{\omega}_t$ restricted to $\bar{\Lambda}_\ell$, with the exception that particles reaching either $y = 0$ or $y = \ell + 1$ never leave, namely the generator \mathcal{L}_ℓ^{SZR} defined as

$$\begin{aligned} \mathcal{L}_\ell^{SZR} f(\chi) := & \mathbf{1}_{\{\chi(\ell) \geq 2\}} (f(\chi^{\ell, \ell+1}) - f(\chi)) + \mathbf{1}_{\{\chi(1) \geq 2\}} (f(\chi^{1,0}) - f(\chi)) \\ & + \sum_{\substack{y,z \in \Lambda_\ell \\ |y-z|=1}} \mathbf{1}_{\{\chi(y) \geq 2\}} (f(\chi^{y,z}) - f(\chi)), \end{aligned}$$

for any $f : \mathbb{N}^{\bar{\Lambda}_\ell} \rightarrow \mathbb{R}$. Both sites $y = 0$ and $y = \ell + 1$ are initially empty and play the role of cemetery states where particles become stuck.

Let $\mathbb{Q}_{\chi_0}^{SZR}$ be the distribution of the process χ_t started from $\chi_0 = \omega_0^{\ell, x}$ and driven by \mathcal{L}_ℓ^{SZR} . We denote by

$$T_\chi = \inf \{t \geq 0 : \chi_t(y) \leq 1, \forall y \in \Lambda_\ell\} \quad (4.8)$$

the time at which all the particles have either left Λ_ℓ or been stuck in a hole. Then, by monotonicity of the event $\{\tilde{\omega}_{T_\ell}(x) = 0\}$ w.r.t. the configurations, we have for any $x \in \mathbb{T}_K$

$$\mathbb{Q}_{K, \omega_0}^{ZR}(\tilde{\omega}_{T_\ell}(x) = 0) \leq \mathbb{Q}_{\chi_0}^{SZR}(\chi_{T_\ell}(0) = 0).$$

Since $\chi_0 = \omega_0^{\ell, x} \in A_\ell$, Lemma 4.3 stated below yields, for any $\varepsilon < \frac{1}{2}$ and any $\ell \geq \bar{\ell}(\varepsilon, \delta)$, that

$$\mathbb{Q}_{K, \omega_0}^{ZR}(\tilde{\omega}_{T_\ell}(x) = 0) \leq \exp(-\ell^\varepsilon)$$

as wanted. Equation (4.7) then concludes the proof.

Lemma 4.3. *For any $\varepsilon < \frac{1}{2}$, there exists $\bar{\ell} := \bar{\ell}(\varepsilon, \delta)$, such that for any $\ell \geq \bar{\ell}$,*

$$\sup_{\chi_0 \in A_\ell} \mathbb{Q}_{\chi_0}^{SZR}(\chi_{T_\ell}(0) = 0) \leq \exp(-\ell^\varepsilon). \quad (4.9)$$

Proof of Lemma 4.3. This result immediately follows from Lemmas 4.4 and 4.5 below.

Lemma 4.4. *For any $\varepsilon < 1$, there exists $\bar{\ell} := \bar{\ell}(\varepsilon, \delta)$ such that for any $\ell \geq \bar{\ell}$,*

$$\sup_{\chi_0 \in A_\ell} \mathbb{Q}_{\chi_0}^{SZR}(T_\chi > T_\ell) \leq \exp(-\ell^\varepsilon).$$

Lemma 4.5. *For any $\varepsilon < \frac{1}{2}$, there exists $\bar{\ell} := \bar{\ell}(\varepsilon, \delta)$ such that for any $\ell \geq \bar{\ell}$,*

$$\sup_{\chi_0 \in A_\ell} \mathbb{Q}_{\chi_0}^{SZR}(\chi_{T_\chi}(0) = 0) \leq \exp(-\ell^\varepsilon).$$

In the next two paragraphs we prove Lemmas 4.4 and 4.5. □

4.2.1. *Proof of Lemma 4.4.* To prove Lemma 4.4, we couple our zero-range process χ_t on $\bar{\Lambda}_\ell$ with two auxiliary processes: $\{\{\zeta_t(y)\}_{y \in \bar{\Lambda}_\ell} : t \geq 0\}$ and $\{\{\Upsilon_t(y)\}_{y \in \bar{\Lambda}_\ell} : t \geq 0\}$, both starting from the same initial state as χ_t . Let us describe the dynamics of these processes.

The auxiliary process ζ_t is also a zero-range process, but with rate 1: no particle gets stuck inside Λ_ℓ , instead each particle moves freely until it reaches sites $y = 0$ or $y = \ell + 1$ where it remains stuck. Its generator \mathcal{L}_ℓ^{FZR} is given by

$$\mathcal{L}_\ell^{FZR} f(\zeta) := f(\zeta^{\ell, \ell+1}) - f(\zeta) + f(\zeta^{1,0}) - f(\zeta) + \sum_{\substack{y,z \in \Lambda_\ell \\ |y-z|=1}} (f(\zeta^{y,z}) - f(\zeta)).$$

We denote by $\mathbb{Q}_{\zeta_0}^{FZR}$ the distribution of the process ζ_t started from $\zeta_0 = \chi_0$ and driven by the generator \mathcal{L}_ℓ^{FZR} above and we denote by

$$T_\zeta = \inf \{t \geq 0 : \zeta_t(y) = 0, \forall y \in \Lambda_\ell\}$$

the time at which all the particles have left Λ_ℓ .

Recall that, since $\chi_0 \in A_\ell$, it contains $(1 + \delta)\ell$ particles in Λ_ℓ . In the second auxiliary process Υ_t the particles in Λ_ℓ perform independent simple random walks with jumps occurring at rate $1/((1 + \delta)\ell)$, and the particles get stuck forever once they have left Λ_ℓ . Its generator \mathcal{L}_ℓ^{IRW} is given by

$$\begin{aligned} \mathcal{L}_\ell^{IRW} f(\Upsilon) := & \frac{\Upsilon(\ell)}{(1 + \delta)\ell} (f(\Upsilon^{\ell, \ell+1}) - f(\Upsilon)) + \frac{\Upsilon(1)}{(1 + \delta)\ell} (f(\Upsilon^{1,0}) - f(\Upsilon)) \\ & + \sum_{\substack{y, z \in \Lambda_\ell \\ |y-z|=1}} \frac{\Upsilon(y)}{(1 + \delta)\ell} (f(\Upsilon^{y,z}) - f(\Upsilon)). \end{aligned}$$

We denote by $\mathbb{Q}_{\Upsilon_0}^{IRW}$ the distribution of the process Υ_t started from $\Upsilon_0 = \chi_0$ and driven by the generator \mathcal{L}_ℓ^{IRW} above, and we denote by

$$T_\Upsilon = \inf \{t \geq 0 : \Upsilon_t(y) = 0, \forall y \in \Lambda_\ell\}$$

the time at which all the particles have left Λ_ℓ .

Given these two new processes, we state the following three lemmas, which put together prove Lemma 4.4.

Lemma 4.6. *For any $\varepsilon < 1$, there exists $\bar{\ell} := \bar{\ell}(\varepsilon, \delta)$, such that for any $\ell \geq \bar{\ell}$,*

$$\sup_{\Upsilon_0 \in A_\ell} \mathbb{Q}_{\Upsilon_0}^{SZR} (T_\Upsilon > T_\ell) \leq \exp(-\ell^\varepsilon).$$

Lemma 4.7. *For any $\ell \geq 1$, and assuming that $\zeta_0 = \Upsilon_0 \in A_\ell$,*

$$\mathbb{Q}_{\zeta_0}^{FZR} (T_\zeta > T_\ell) \leq \mathbb{Q}_{\Upsilon_0}^{IRW} (T_\Upsilon > T_\ell).$$

Lemma 4.8. *For any $\ell \geq 1$, and assuming that $\chi_0 = \zeta_0 \in \mathbb{N}^{\Lambda_\ell}$,*

$$\mathbb{Q}_{\chi_0}^{SZR} (T_\chi > T_\ell) \leq \mathbb{Q}_{\zeta_0}^{FZR} (T_\zeta > T_\ell).$$

The proofs of Lemmas 4.7 and 4.8 are based on coupling arguments. Lemma 4.6 uses classic estimates on simple random walks, which are stated and proved in Appendix A. Lemma 4.4 follows immediately from these three lemmas.

4.2.2. Proof of Lemma 4.6. Consider a single particle performing a symmetric random walk, jumping left and right at the same rate $1/((1 + \delta)\ell)$. Thanks to the estimate (A.1) given in Appendix A, regardless of its initial position in Λ_ℓ the probability for this particle to remain in Λ_ℓ until time $T_\ell = \ell^4$ is bounded, for any $\ell > \bar{\ell}(\delta)$ by $e^{-C\ell}$ for some positive constant $C = C(\delta)$. Since $\Upsilon_0 \in A_\ell$, there are initially $(1 + \delta)\ell$ such particles located in Λ_ℓ . Moreover, since in Υ these particles move independently, by definition of T_Υ , it is straightforward to obtain as wanted, for any $\Upsilon_0 \in A_\ell$,

$$\mathbb{Q}_{\Upsilon_0}^{IRW} (T_\Upsilon > T_\ell) \leq 1 - (1 - e^{-C\ell})^{(1+\delta)\ell} \sim (1 + \delta)\ell e^{-C\ell}.$$

Given $\varepsilon < 1$, choosing $\bar{\ell}$ large enough such that $(1 - (1 - e^{-C\bar{\ell}})^{(1+\delta)\bar{\ell}}) e^{\bar{\ell}^\varepsilon} \leq 1$ proves the lemma.

4.2.3. *Proof of Lemma 4.7.* Fix an initial configuration $\zeta_0 = \Upsilon_0 \in A_\ell$. The difference between the processes Υ and ζ is that in the second one the jump rate of a given particle depends on the current configuration. However, since initially the total number of particles in Λ_ℓ is $(1 + \delta)\ell$, there are certainly never more than $(1 + \delta)\ell$ particles on a given site of Λ_ℓ , and therefore the particles in ζ are faster than those in Υ . For completeness we give a full construction of a coupling between the two processes.

The particles' trajectories will be the same for both ζ and Υ , but the order and times of the jumps will change. Let us start from a configuration $\zeta_0 = \Upsilon_0$ with $(1 + \delta)\ell$ particles in Λ_ℓ . For any $1 \leq p \leq (1 + \delta)\ell$, denote by x_p the initial position of the p -th particle in ζ_0 . We define the random trajectory $X_p(0) = x_p, \dots, X_p(k_p) \in \{0, \ell + 1\}$ of the p -th particle until it leaves Λ_ℓ , where k_p is the random number of jumps necessary to do so. In other words, we define $X_p(0) = x_p$ and for each $k \geq 1$, we let

$$X_p(k) = \begin{cases} X_p(k-1) + 1 & \text{with prob. } \frac{1}{2}, \\ X_p(k-1) - 1 & \text{with prob. } \frac{1}{2}, \end{cases}$$

and we define $k_p = \inf \{k \geq 1 : X_p(k) \in \{0, \ell + 1\}\}$. We finally denote by

$$\bar{X} = (X_p(k))_{1 \leq p \leq (1+\delta)\ell, 0 \leq k \leq k_p} \quad \text{and} \quad \bar{k} = (k_p)_{1 \leq p \leq (1+\delta)\ell}$$

the family of these trajectories and their lengths.

Given \bar{k}, \bar{X} , let us now define the order and times of the jumps in each process. Let

$$K = \sum_{p=1}^{(1+\delta)\ell} k_p$$

be the total number of jumps to perform in the configuration in order to reach a frozen state, which is the same for ζ and Υ since they start from the same configuration. We are going to define update times $t_0^\Upsilon = 0 < t_1^\Upsilon < \dots < t_K^\Upsilon = T_\Upsilon$ and $t_0^\zeta = 0 < t_1^\zeta < \dots < t_K^\zeta = T_\zeta$ for both processes, and build ζ and Υ in such a way that for any $i \leq K$,

$$t_i^\zeta \leq t_i^\Upsilon.$$

At time t_i^Υ , (resp. t_i^ζ) we are going to choose a particle p_i^Υ (resp. p_i^ζ) which we will move in Υ (resp. ζ) according to its predetermined trajectory $X_{p_i^\Upsilon}$ (resp. $X_{p_i^\zeta}$).

Assume that both ζ and Υ are defined respectively until time t_i^ζ and t_i^Υ . We denote by $n_i^\Upsilon \in \{1, \dots, (1 + \delta)\ell\}$ the number of *particles* still in Λ_ℓ at time t_i^Υ and by $n_i^\zeta \in \{1, \dots, \ell\}$ the number of *sites* in Λ_ℓ still occupied by at least one particle in ζ at time t_i^ζ . We then sample two exponential times τ_{i+1}^ζ and τ_{i+1}^Υ , with respective parameters n_i^ζ and $n_i^\Upsilon / ((1 + \delta)\ell)$. For any $i < K$, in Λ_ℓ there is at least one occupied site in ζ , therefore $n_i^\zeta \geq 1$. Furthermore, there are in Λ_ℓ at most $(1 + \delta)\ell$ particles in Υ , so that $n_i^\Upsilon / ((1 + \delta)\ell) \leq 1$. It is therefore possible to couple τ_{i+1}^ζ and τ_{i+1}^Υ in such a way that

$$\tau_{i+1}^\zeta \leq \tau_{i+1}^\Upsilon. \tag{4.10}$$

We then let

$$t_{i+1}^\Upsilon = t_i^\Upsilon + \tau_{i+1}^\Upsilon \leq t_i^\zeta + \tau_{i+1}^\zeta = t_{i+1}^\zeta.$$

We leave Υ (resp. ζ) unchanged in the time segment $(t_i^\Upsilon, t_{i+1}^\Upsilon)$ (resp. $(t_i^\zeta, t_{i+1}^\zeta)$). In order to close the construction, we now need to build $\zeta_{t_{i+1}^\zeta}$ and $\Upsilon_{t_{i+1}^\Upsilon}$. To do so, we need to choose

the particle to move, and then the chosen particle will move according to the pre-chosen trajectories \bar{X} , as we explained above.

To choose the index of the particle to move in ζ , we choose one site uniformly among the occupied sites in $\zeta_{t_i^\zeta}$, and then choose the particle to update uniformly among the particles present at this site. We then update the position of the particles according to the trajectories stored in \bar{X} . To choose the particle to move in Υ , we simply choose the particle uniformly among the particles in $\Upsilon_{t_i^\Upsilon}$, and update the position of this particle according to the fixed trajectories \bar{X} . After the final time $t_K^\zeta = T_\zeta$ (resp. $t_K^\Upsilon = T_\Upsilon$), the process ζ (resp. Υ) remains frozen.

It is not difficult to see that under this construction, ζ and Υ are respectively distributed according to $\mathbb{Q}_{\zeta_0}^{FZR}$ and $\mathbb{Q}_{\Upsilon_0}^{IRW}$. Furthermore, by construction, the update times for each step satisfy (4.10), therefore

$$t_K^\zeta = T_\zeta \leq t_K^\Upsilon = T_\Upsilon,$$

thus concluding the proof of Lemma 4.7.

4.2.4. *Proof of Lemma 4.8.* Let us couple χ and ζ through the following graphical construction. Attach to every site in Λ_ℓ independently a Poisson process of parameter 1 and think of it as a sequence of clock rings signalling possible update times for the process. Decorate the clock rings with independent $\text{Ber}(\frac{1}{2})$ variables (also independent from the Poisson processes). In a given configuration of the process χ , we call *excess particle* any particle which is not alone on a site. The process χ (resp. ζ) can be constructed by saying that when a clock rings at site x , one excess particle at x (resp. one particle present at x) jumps right or left depending on the corresponding Bernoulli variable. Now in order to couple the two processes in a way to have $T_\chi \leq T_\zeta$, we turn ζ into a *two-color process*. In the initial configuration, in ζ , one particle per site is coloured red (corresponding to the particles that will remain stuck in χ) and the others blue. We build the processes in such a way that the following property is preserved:

$$\text{The number of excess particles in } \chi \text{ is equal to the number of blue particles in } \zeta. \quad (4.11)$$

When a clock rings in a configuration satisfying this condition, three cases arise.

- If the ring occurs on a site with no particle in either process, nothing happens.
- If the ring occurs on a site with only red particles in ζ , one of them jumps according to the Bernoulli variable and nothing happens in χ (where there is either no particle or a single – and therefore blocked – one).
- If the ring occurs on a site with excess particles in χ , one of them jumps according to the Bernoulli variable. In ζ , one of the blue particles makes the same jump. If the excess particle in χ becomes stuck after the jump, then color the corresponding particle in ζ red.

It is immediate to check that property (4.11) is preserved through any of these transitions. Also, χ has the desired law, as does the color-blind version of ζ . Moreover, with this description, T_χ is the time at which all particles in ζ have either turned red or left the box, which clearly happens before they all leave the box. This ends the proof of Lemma 4.8.

Now it only remains to prove Lemma 4.5.

4.2.5. *Proof of Lemma 4.5.* For any $\chi_0 \in A_\ell$, we are going to estimate

$$\mathbb{Q}_{\chi_0}^{SZR}(\chi_{T_\chi}(0) = 0).$$

The core idea to estimate the probability above is that in order for each particle that did not get stuck to leave Λ_ℓ at $x = \ell + 1$ rather than at $x = 0$, an unusual number of particle jumps to the right must have occurred, which happens with probability decaying exponentially fast in $\ell^{1/2}$.

For any realization of χ , we denote by $L = L(\chi)$ the random number of jumps occurring before T_χ , i.e. the number of particle jumps needed to reach a frozen state:

$$L(\chi) = \#\{0 \leq t \leq T_\chi : \chi_{t-} \neq \chi_t\}.$$

As before, since $\chi_0 \in A_\ell$, there are $(1 + \delta)\ell$ particles initially in Λ_ℓ , and each one of those typically requires $\mathcal{O}_\ell(\ell^2)$ jumps to either leave or get stuck. For any $\varepsilon < \frac{1}{2}$, we claim that there must exist $\bar{\ell}(\varepsilon, \delta)$ such that for any $\ell \geq \bar{\ell}$

$$\mathbb{Q}_{\chi_0}^{SZR} \left(L(\chi) > \ell^{3+1/2} \right) \leq \exp(-\ell^\varepsilon). \quad (4.12)$$

To prove (4.12), we use the same coupling as in the proof of Lemma 4.8, where the process χ is compared with the process ζ where particles do not get stuck in empty sites. Then, denote by

$$L(\zeta) = \#\{t \leq T_\zeta : \zeta_{t-} \neq \zeta_t\},$$

the random number of particle jumps which are necessary for each particle in ζ to leave Λ_ℓ . Using the coupling described in the proof of Lemma 4.8, we obtain that

$$L(\chi) \leq L(\zeta),$$

and thus we need to estimate $\mathbb{Q}_{\zeta_0}^{FZR} \left(L(\zeta) > \ell^{3+1/2} \right)$.

Since in ζ , the number L_i of jumps which are necessary for the i -th particle to leave Λ_ℓ does not depend on the other trajectories (even though the time to leave does), we can now write (using the fact that $\zeta_0 = \chi_0$ is in A_ℓ)

$$\begin{aligned} \mathbb{Q}_{\zeta_0}^{FZR} \left(L(\zeta) = \sum_{i=1}^{(1+\delta)\ell} L_i > \ell^{3+1/2} \right) &\leq \mathbb{Q}_{\zeta_0}^{FZR} \left(\exists i \in \{1, \dots, (1+\delta)\ell\}, L_i > \frac{\ell^{2+1/2}}{(1+\delta)} \right) \\ &\leq (1+\delta)\ell \sup_{i \in \{1, \dots, (1+\delta)\ell\}} \mathbb{Q}_{\zeta_0}^{FZR} \left(L_i > \frac{\ell^{2+1/2}}{(1+\delta)} \right). \end{aligned}$$

As a consequence of equation (A.2) given in Lemma A.1, regardless of the starting point of each particle, there exists a universal constant $C > 0$ such that for any ℓ , and any $1 \leq i \leq (1+\delta)\ell$,

$$\mathbb{Q}_{\zeta_0}^{FZR} \left(L_i > \frac{\ell^{2+1/2}}{(1+\delta)} \right) \leq \exp(-C\ell^{1/2}),$$

thus proving (4.12) for any $\varepsilon < \frac{1}{2}$. Given (4.12), in order to prove Lemma 4.5, we now only need to consider the realizations of χ for which $L(\chi) \leq \ell^{3+1/2}$, and therefore Lemma 4.5 is a consequence of Lemma 4.9 below.

Lemma 4.9. *For any $\varepsilon < \frac{1}{2}$, there exists $\bar{\ell}(\varepsilon, \delta)$, such that for any $\ell \geq \bar{\ell}$*

$$\sup_{\chi_0 \in A_\ell} \mathbb{Q}_{\chi_0}^{SZR} \left(\chi_{T_\chi}(0) = 0 \mid L(\chi) \leq \ell^{3+1/2} \right) \leq \exp(-\ell^\varepsilon). \quad (4.13)$$

4.2.6. *Proof of Lemma 4.9.* To prove this last result, we introduce a discrete time random walker $(X_k)_{k \leq L(\chi)}$ evolving on \mathbb{Z} , starting from $X_0 = 0$ and jumping each time a particle of χ present in Λ_ℓ jumps, in the same direction as the particle. In other words, if the k -th particle jump in χ is to the left (resp. to the right), then $X_k = X_{k-1} - 1$ (resp. $X_k = X_{k-1} + 1$).

Let us define

$$Z(t, \chi) := \sum_{y=0}^{\ell+1} y \chi_t(y),$$

and let $0 < \tau_1 < \dots < \tau_{L(\chi)} = T_\chi$ be the successive jump times in χ . Then, an elementary computation yields that

$$Z(\tau_k, \chi) - Z(0, \chi) = X_k, \quad (4.14)$$

because each jump to the right (resp. left) increases (resp. decreases) $Z(t, \chi)$ by 1. All the jumps occurring in χ being symmetric, the law of $(X_k)_{k \leq L(\chi)}$ is that of a symmetric random walker in discrete time. To avoid burdensome notations, we now assume that the random walk X is defined for any time $k \in \mathbb{N}$ by completing the random walk after $L(\chi)$ independently from χ .

Assuming that no particle in χ reaches $x = 0$, we know that after T_χ , each particle must have either gotten stuck in Λ_ℓ , or reached the site $x = \ell + 1$. Note that in order to minimize $Z(T_\chi, \chi)$ in this case, there must be exactly one particle per site in Λ_ℓ , and all the remaining particles must be at site $\ell + 1$. In that case, since the initial number of particles in $\bar{\Lambda}_\ell$ is $(1 + \delta)\ell$,

$$Z(T_\chi, \chi) = (2\delta + 1) \frac{\ell(\ell + 1)}{2}.$$

We can therefore write

$$\chi_{T_\chi}(0) = 0 \implies Z(T_\chi, \chi) \geq (2\delta + 1) \frac{\ell(\ell + 1)}{2}. \quad (4.15)$$

Recall that $\chi_0 \in A_\ell$ and therefore by definition

$$Z(0, \chi) \leq \left(1 + \frac{\delta}{2}\right) (1 + \delta) \frac{\ell(\ell + 1)}{2}. \quad (4.16)$$

Thanks to (4.15) and (4.16), we obtain that

$$(\chi_0 \in A_\ell \text{ and } \chi_{T_\chi}(0) = 0) \implies Z(T_\chi, \chi) - Z(0, \chi) \geq (\delta - \delta^2) \frac{\ell(\ell + 1)}{4} \geq (\delta - \delta^2) \frac{\ell^2}{4}.$$

Recall that we assumed $0 < \delta < 1$. Letting $C = C(\delta) := (\delta - \delta^2)/4 > 0$, we can finally write for any $\chi_0 \in A_\ell$

$$\begin{aligned} \mathbb{Q}_{\chi_0}^{SZR} \left(\chi_{T_\chi}(0) = 0 \mid L(\chi) \leq \ell^{3+1/2} \right) \\ \leq \mathbb{Q}_{\chi_0}^{SZR} \left(Z(T_\chi, \chi) - Z(0, \chi) \geq C\ell^2 \mid L(\chi) \leq \ell^{3+1/2} \right). \end{aligned} \quad (4.17)$$

Now, recall that by equation (4.14), $Z(T_\chi, \chi) - Z(0, \chi) = X_{L(\chi)}$, and let \mathbb{P}^0 denote the distribution of a discrete time symmetric random walk on \mathbb{Z} started from 0. We get

$$\mathbb{Q}_{\chi_0}^{SZR} \left(Z(T_\chi, \chi) - Z(0, \chi) \geq C\ell^2 \mid L(\chi) \leq \ell^{3+1/2} \right) \leq \mathbb{P}^0 \left(\sup_{0 \leq k \leq \ell^{3+1/2}} X_k \geq C\ell^2 \right).$$

Since X_k is a discrete time symmetric random walk, we use Doob's inequality on the positive sub-martingale $e^{\lambda X_k}$ for $\lambda = C\ell^{-\frac{3}{2}}$. This yields

$$\mathbb{P}^0 \left(\sup_{0 \leq k \leq \ell^{3+1/2}} X_k \geq C\ell^2 \right) \leq \frac{\exp \left[\ell^{3+1/2} \left(\frac{C^2 \ell^{-3}}{2} + \mathcal{O}_\ell(\ell^{-6}) \right) \right]}{\exp \left(C^2 \sqrt{\ell} \right)}.$$

Therefore, for any $\varepsilon < \frac{1}{2}$, there exists an integer $\bar{\ell}(\varepsilon, \delta)$ such that for any $\ell \geq \bar{\ell}$

$$\mathbb{P}^0 \left(\sup_{0 \leq k \leq \ell^{3+1/2}} X_k \geq C\ell^2 \right) \leq \exp(-\ell^\varepsilon),$$

which concludes the proof of Lemma 4.9.

4.3. Probability of the regular configurations. In this section, we prove Lemma 4.2, namely:

$$\lim_{N \rightarrow \infty} \mu_N \left(K(\eta) > \ell_N \quad \text{and} \quad \Pi(\eta) \in B_{K(\eta), \ell_N} \right) = 1,$$

where we have already defined $\ell_N = (\log N)^8$. We start by proving that

$$\lim_{N \rightarrow \infty} \mu_N \left(K(\eta) > \ell_N \right) = 1. \quad (4.18)$$

Recall that $K(\eta)$ is the number of empty sites in the exclusion configuration η , and that $\rho_0(\cdot) \in (\frac{1}{2}, 1]$, and ρ_0 is smooth. We assume that ρ_0 is not identically equal to 1, in which case the model is trivial. Then, since ρ_0 is smooth, there exist $\alpha > 0$, $u \in \mathbb{T}$, and $\varepsilon > 0$ such that uniformly in $[u, u + \varepsilon]$, we have $\rho_0 < 1 - \alpha$. Since under μ_N , the state of each site is chosen independently, we can use a standard large deviation estimate: for any $c < \alpha\varepsilon$, there exists a constant $C(c)$ such that

$$\mu_N \left(K(\eta) \leq cN \right) \leq e^{-cN},$$

which proves (4.18).

To close the proof of Lemma 4.2, we now prove that

$$\lim_{N \rightarrow \infty} \mu_N \left(K(\eta) > \ell_N \quad \text{and} \quad \Pi(\eta) \notin B_{K(\eta), \ell_N} \right) = 0. \quad (4.19)$$

Because of the random nature of the number $K(\eta)$ of sites of $\Pi(\eta)$, and of the dependency between $K(\eta)$ and the number of particles in $\Pi(\eta)$, this proof, whose principle is rather simple, becomes quite technical. We first treat the case where the profile ρ_0 is homogeneous with density $\bar{\rho} \in (\frac{1}{2}, 1)$, and then prove that the case where ρ_0 is not constant follows, thus concluding the proof of Lemma 4.2. Since we will need to differentiate the cases depending on the initial profile, we introduce explicitly ρ_0 in our notation, and denote by μ_{N, ρ_0} the initial product measure fitting the macroscopic profile ρ_0 (defined in (2.7)).

Lemma 4.10. *For any $\bar{\rho} \in (\frac{1}{2}, \frac{2}{3})$, there exists $\delta = \delta(\bar{\rho}) \in (0, 1)$ such that*

$$\lim_{N \rightarrow \infty} \mu_{N, \bar{\rho}} \left(K(\eta) > \ell_N \quad \text{and} \quad \Pi(\eta) \notin B_{K(\eta), \ell_N} \right) = 0$$

where $\mu_{N, \bar{\rho}}$ denotes the measure associated with the constant profile $\rho_0 \equiv \bar{\rho}$.

Corollary 4.11. *Fix a non-constant smooth function $\rho_0 : \mathbb{T} \rightarrow (\frac{1}{2}, 1]$, then, there exists $\delta \in (0, 1)$ such that*

$$\lim_{N \rightarrow \infty} \mu_{N, \rho_0} \left(K(\eta) > \ell_N \quad \text{and} \quad \Pi(\eta) \notin B_{K(\eta), \ell_N} \right) = 0.$$

We assume for now Lemma 4.10 in order to prove Corollary 4.11.

Proof of Corollary 4.11. Fix $\bar{\rho} \in \left(\frac{1}{2}, \frac{2}{3} \wedge \min_{u \in \mathbb{T}} \rho_0(u)\right)$.

We claim that in order to prove Corollary 4.11, it is enough to build on the same probability space an increasing family of configurations η^λ on \mathbb{T}_N , for $\lambda \in [0, 1]$, with joint distribution ν_N , such that η^0 has law $\mu_{N, \bar{\rho}}$, η^1 has law μ_{N, ρ_0} and

$$\nu_N \left(\Pi(\eta^1) \notin B_{K(\eta^1), \ell_N} \mid \Pi(\eta^0) \in B_{K(\eta^0), \ell_N}, K(\eta^0) > \ell_N \text{ and } K(\eta^1) > \ell_N \right) = 0. \quad (4.20)$$

Indeed, since η^0 is distributed according to $\mu_{N, \bar{\rho}}$, according to Lemma 4.10 and equation (4.18) (which holds for any density profile not identically equal to 1), the probability of the conditioning event goes to 1 as $N \rightarrow \infty$, therefore (4.20) proves Corollary 4.11.

To build the family η^λ , and prove (4.20), consider a family of N i.i.d. uniform $\mathcal{U}([0, 1])$ variables U_1, \dots, U_N . Define for any $\lambda \in [0, 1]$ the function

$$\rho^\lambda = \bar{\rho} + \lambda(\rho_0 - \bar{\rho}),$$

which is an interpolation between $\bar{\rho}$ and ρ_0 . We then write for any $x \in \mathbb{T}_N$

$$\eta^\lambda(x) = \mathbf{1}_{\{U_x \leq \rho^\lambda(x/N)\}}.$$

Then, as wanted, $\eta^0 \sim \mu_{N, \bar{\rho}}$ and $\eta^1 \sim \mu_{N, \rho_0}$. We now prove that (4.20) holds.

As λ goes from 0 to 1, particles are added in η^0 until it becomes η^1 . Fix an exclusion configuration η with an empty site at x , and consider the configuration $\eta + \mathbf{1}_x$ where a particle has been added at site x (namely, $\mathbf{1}_x$ is the configuration such that $\mathbf{1}_x(x) = 1$ and $\mathbf{1}_x(y) = 0$ for any $y \neq x$). Since x is empty in η , it is associated with a site \tilde{x} in $\Pi(\eta)$. Then, the associated zero-range configuration $\Pi(\eta + \mathbf{1}_x)$ is obtained by deleting the site \tilde{x} in $\Pi(\eta)$, and putting all the particles that were present at \tilde{x} in $\Pi(\eta)$ on the previous site $\tilde{x} - 1$. Another particle is then added on the site $\tilde{x} - 1$ to represent the empty site that is now a particle. We follow this construction, throughout the transition from η^0 to η^1 . Assuming that $\Pi(\eta^0) \in B_{K(\eta^0), \ell_N}$, then in $\Pi(\eta_0)$, all the boxes of size ℓ_N already contain $(1 + \delta)\ell_N$ particles. Then, as λ goes from 0 to 1, we add particles to η^0 , and therefore delete sites in $\Pi(\eta^0)$, while placing all the particles on the deleted site on the left neighbor of the deleted site. In each box $\{x + 1, \dots, x + \ell_N\}$ in the smaller torus, the cumulated distance between the $(1 + \delta)\ell_N$ left-most particles to x therefore diminishes as we go from $\Pi(\eta^0)$ to $\Pi(\eta^1)$. This construction proves equation (4.20), and concludes the proof of the corollary. \square

Proof of Lemma 4.10. In this proof, we write μ_N for $\mu_{N, \bar{\rho}}$. One technical issue is that even though η is now translation invariant under μ_N , $\Pi(\eta)$ is not. To solve this, we map the exclusion configuration η to a zero-range configuration on Λ_ℓ starting from an arbitrary point in \mathbb{T}_N .

Fix $x \in \mathbb{T}_N$ and a positive integer ℓ such that $(2 + \delta)\ell < N$. As represented in Figure 4.1, let $\eta_{(x, \ell)}$ be the exclusion configuration on $\{0, 1\}^{\mathbb{T}_N}$ obtained from η by keeping at most the first $(1 + \delta)\ell$ particles in $\{x + 1, x + 2, \dots, x\}$ (in cyclic order). We define $\Pi_{(x, \ell)}(\eta) \in \mathbb{N}^{\bar{\Lambda}_\ell}$ as the zero-range configuration of size ℓ in which

- $\Pi_{(x, \ell)}(\eta)(0) = \Pi_{(x, \ell)}(\eta)(\ell + 1) = 0$,
- $\Pi_{(x, \ell)}(\eta)(y)$ is the number of particles between the y -th and $(y + 1)$ -th zero in $\eta_{(x, \ell)}$, starting as if⁵ the first zero were at x .

⁵Due to the condition $(2 + \delta)\ell < N$, $\Pi_{(x, \ell)}(\eta)$ actually does not depend on $\eta(x)$.

Note that, if $K(\eta) > \ell$, for any $y \in \mathbb{T}_{K(\eta)}$, there exists $x \in \mathbb{T}_N$ such that

$$\Pi(\eta)^{\ell,y} = \Pi_{(x,\ell)}(\eta) \in \mathbb{N}^{\Lambda_\ell} \quad (4.21)$$

(recall that for any zero-range configuration ω , the configuration $\omega^{\ell,y}$ has been defined in (4.4)). With (4.21), we can write, for N large enough to have $(2 + \delta)\ell_N < N$

$$\begin{aligned} \mu_N \left(K(\eta) > \ell_N \text{ and } \Pi(\eta) \notin B_{K(\eta), \ell_N} \right) & \\ & \leq \mu_N \left(K(\eta) > \ell_N \text{ and } \exists y \in \mathbb{T}_{K(\eta)}, \Pi(\eta)^{\ell_N, y} \notin A_{\ell_N} \right) \\ & \leq \mu_N \left(K(\eta) > \ell_N \text{ and } \exists x \in \mathbb{T}_N, \Pi_{(x, \ell_N)}(\eta) \notin A_{\ell_N} \right) \\ & \leq N \max_{x \in \mathbb{T}_N} \mu_N \left(\Pi_{(x, \ell_N)}(\eta) \notin A_{\ell_N} \right) \\ & = N \mu_N \left(\Pi_{(1, \ell_N)}(\eta) \notin A_{\ell_N} \right), \end{aligned}$$

because μ_N is translation invariant.

We now argue that we can estimate the last probability via standard properties of i.i.d. sequences of geometric random variables.

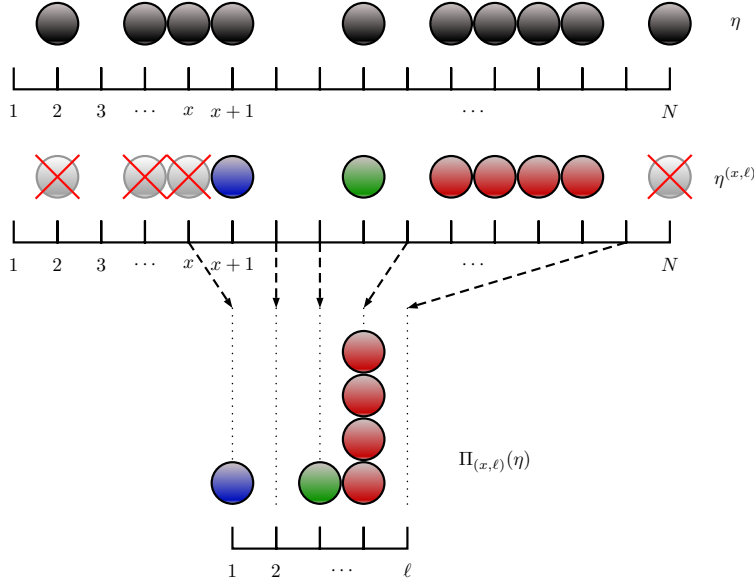


FIGURE 4.1. Representation of η , $\eta_{(x,\ell)}$, $\Pi_{(x,\ell)}(\eta)$, for $\ell = 5$ and $(1 + \delta)\ell = 6$. Starting from η , we obtain $\eta_{(x,\ell)}$ by keeping the first $(1 + \delta)\ell$ particles to the right of x (x excluded). We then obtain $\Pi_{(x,\ell)}(\eta)$ by taking the zero-range configuration associated with the first ℓ empty sites in $\eta_{(x,\ell)}$, where x is considered to be the first empty site.

Let us uncover one by one the variables given by the states of η on sites $1, 2, \dots, (2 + \delta)\ell_N$. Since $(2 + \delta)\ell_N < N$, they are i.i.d. Then $\{n_{\ell_N}(\Pi_{(1, \ell_N)}(\eta)) < (1 + \delta)\ell_N\}$ is the event that in this process we see $\ell_N + 1$ zeros before seeing $(1 + \delta)\ell_N$ particles and therefore

$$\mu_N \left(n_{\ell_N}(\Pi_{(1, \ell_N)}(\eta)) < (1 + \delta)\ell_N \right) = P \left(\sum_{k=1}^{\ell_N} G_k < (1 + \delta)\ell_N \right), \quad (4.22)$$

where the G_k are i.i.d. geometric random variables of parameter $\bar{\rho}$, namely for $k \in \mathbb{N}$, $P(G_1 = k) = (1 - \bar{\rho})\bar{\rho}^k$.

For the same reason, $\Pi_{(1, \ell_N)}$ is dominated by a configuration given by i.i.d. $\text{Geom}(\bar{\rho})$ variables on Λ_{ℓ_N} . Consequently,

$$\begin{aligned} \mu_N \left[Z_{\ell_N}(\Pi_{(1, \ell_N)}(\eta)) > \left(1 + \frac{\delta}{2}\right) (1 + \delta) \frac{\ell_N(\ell_N + 1)}{2} \right] \\ \leq P \left[\sum_{k=1}^{\ell_N} k G_k > \left(1 + \frac{\delta}{2}\right) (1 + \delta) \frac{\ell_N(\ell_N + 1)}{2} \right]. \end{aligned} \quad (4.23)$$

Now it only remains to show that the right hand side in (4.22) and (4.23) is $o_N(1/N)$ for a well-chosen $\delta(\bar{\rho})$. We choose $\delta \in (0, 1)$ such that

$$1 + \delta < \frac{\bar{\rho}}{1 - \bar{\rho}} < (1 + \delta) \left(1 + \frac{\delta}{2}\right). \quad (4.24)$$

This is possible because $\bar{\rho} \in (\frac{1}{2}, \frac{2}{3})$. By standard arguments, there exists $C = C(\delta, \bar{\rho}) > 0$ such that for N large enough

$$P \left(\sum_{k=1}^{\ell_N} G_k < (1 + \delta)\ell_N \right) \leq e^{-C\ell_N}. \quad (4.25)$$

To estimate (4.23), we crop the variables G_k so that they remain bounded, and then apply the Azuma-Hoeffding inequality to $\sum_{k=1}^{\ell_N} k G_k$. For any integer p_N to be determined later on, and for any $k = 1, \dots, \ell_N$, we define

$$\widehat{G}_k = G_k \wedge p_N.$$

For any given $k = 1, \dots, \ell_N$, we can write

$$P(\widehat{G}_k \neq G_k) = \sum_{j=p_N+1}^{\infty} (1 - \bar{\rho})\bar{\rho}^j \leq \bar{\rho}^{p_N+1}.$$

Thus, letting $p_N = \ell_N^{1/4}$, we obtain by a simple union bound

$$P(\exists k \leq \ell_N, \widehat{G}_k \neq G_k) \leq \ell_N \bar{\rho}^{p_N+1} = \mathcal{O}_N \left(\ell_N e^{-C\ell_N^{1/4}} \right). \quad (4.26)$$

Let us define the martingale

$$M_K = \sum_{k=1}^K k (\widehat{G}_k - E[\widehat{G}_k]), \quad \text{for } K = 1, \dots, \ell_N,$$

with $M_0 = 0$ and $M_K = M_{\ell_N}$ for $K \geq \ell_N$. Note that $E[\widehat{G}_k] = \frac{\bar{\rho}}{1 - \bar{\rho}} + \mathcal{O}_N(p_N \bar{\rho}^{p_N})$. For any $\alpha > \frac{\bar{\rho}}{1 - \bar{\rho}}$,

$$\begin{aligned} P \left(\sum_{k=1}^{\ell_N} k G_k > \alpha \frac{\ell_N(\ell_N + 1)}{2} \right) &\leq P(\exists k \leq \ell_N, \widehat{G}_k \neq G_k) \\ &\quad + P \left(M_{\ell_N} > \left(\alpha - E[\widehat{G}_1] \right) \frac{\ell_N(\ell_N + 1)}{2} \right). \end{aligned}$$

Recall that we defined $\ell_N = (\log N)^8$. To prove that the right hand side of (4.23) is $o_N(1/N)$, it is therefore sufficient to prove that for $\beta > 0$ there exists $C(\beta) > 0$ such that

$$P\left(M_{\ell_N} > \frac{\beta \ell_N (\ell_N + 1)}{2}\right) = \mathcal{O}_N(e^{-C\ell_N^{1/2}}).$$

By definition, the martingale $(M_K)_K$ has increments bounded a.s. as follows:

$$|M_K - M_{K-1}| \leq K p_N.$$

Therefore we can apply the Azuma-Hoeffding inequality to obtain that

$$P\left(M_{\ell_N} > \frac{\beta \ell_N (\ell_N + 1)}{2}\right) \leq \exp\left(-\frac{\beta^2 \ell_N^2 (\ell_N + 1)^2}{8 \sum_{K=1}^{\ell_N} K^2 p_N^2}\right) = \mathcal{O}_N\left(e^{-\frac{3}{8}\beta^2 \sqrt{\ell_N}}\right),$$

which concludes the proof. \square

5. PROOF OF THEOREM 2.2: HYDRODYNAMIC LIMITS

Our proof of the hydrodynamic limit relies on the classical *entropy method* (which is explained in details for instance in [11, Chapter 5]), but requires adaptations to solve the ergodicity issue, and to account for the non-product invariant measures. In Section 5.1, we describe the main steps of the proof in the form of two distinct results:

- the first one states that after reaching the ergodic component, the macroscopic density profile of the system is very close to the initial density profile ρ_0 . This result is proved in Section 5.2;
- the second result states that starting from the ergodic component, the hydrodynamic limit holds, and is proved in Section 5.3. The latter requires the classical Replacement Lemma, which, because of the shape of our canonical and grand canonical measures, requires significant work. The proof of the Replacement Lemma is therefore postponed to Section 7, after we have proved the relevant properties for the canonical measures in Section 6.

5.1. Sketch of the proof. Recall that we defined in Theorem 2.4 the time $t_N := T_{\ell_N}/N^2$ with $T_{\ell_N} = (\log N)^{32}$. We denote by

$$\tau_N = \inf\{t > 0 : \eta_t \in \mathcal{E}_N\},$$

the hitting time of the ergodic component for the process rescaled diffusively in time. Recall that \mathcal{E}_N has been defined in (2.3). We proved in Theorem 2.4 that

$$\limsup_{N \rightarrow \infty} \mathbb{P}_{\mu_N}(\tau_N > t_N) = 0.$$

Let $\tilde{\mu}_N$ denote the distribution of η_{t_N} at time t_N , conditioned to have already reached the ergodic component, namely

$$\tilde{\mu}_N(\eta') = \mathbb{P}_{\mu_N}(\eta_{t_N} = \eta' \mid \tau_N \leq t_N), \quad \eta' \in \mathcal{E}_N. \quad (5.1)$$

Proof of Theorem 2.2. Fix a time $t > t_N$ and $\delta > 0$. We bound from above the probability in (2.8) by

$$\mathbb{P}_{\mu_N}\left(\left|\frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi\left(\frac{x}{N}\right) \eta_t(x) - \int_{\mathbb{T}} \varphi(u) \rho(t, u) du\right| > \delta \mid \tau_N \leq t_N\right) + \mathbb{P}_{\mu_N}(\tau_N > t_N).$$

As already pointed out, the second term above vanishes as a consequence of Theorem 2.4. The first term can be rewritten for any $t > t_N$, by the Markov property, as

$$\mathbb{P}_{\tilde{\mu}_N} \left(\left| \frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi \left(\frac{x}{N} \right) \eta_{t-t_N}(x) - \int_{\mathbb{T}} \varphi(u) \rho(t, u) du \right| > \delta \right) \quad (5.2)$$

$$\leq \mathbb{P}_{\tilde{\mu}_N} \left(\left| \frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi \left(\frac{x}{N} \right) (\eta_t(x) - \eta_{t-t_N}(x)) \right| > \frac{\delta}{3} \right) \quad (5.3)$$

$$+ \mathbb{P}_{\tilde{\mu}_N} \left(\left| \frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi \left(\frac{x}{N} \right) \eta_t(x) - \int_{\mathbb{T}} \varphi(u) \rho(t, u) du \right| > \frac{\delta}{3} \right). \quad (5.4)$$

The first term in the right hand side (5.3) vanishes as a consequence of the Markov inequality and Lemma 5.1 stated below, by taking for ν_N the distribution under $\mathbb{P}_{\tilde{\mu}_N}$ of η_{t-t_N} . The second term (5.4) vanishes as well according to Proposition 5.2 stated below.

Fix $t > 0$, then for some N large enough, we have $t > t_N$, so that the bound (5.2) indeed holds. This proves the vanishing limit (2.8) for any positive time t . Since it is trivially true for $t = 0$ as well, this proves Theorem 2.2. \square

Lemma 5.1 (Empirical measure at the transience time). *For any smooth test function $\varphi : \mathbb{T} \rightarrow \mathbb{R}$, and any initial measure ν_N on $\{0, 1\}^{\mathbb{T}_N}$, we have*

$$\lim_{N \rightarrow \infty} \mathbb{E}_{\nu_N} \left[\left| \frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi \left(\frac{x}{N} \right) (\eta_{t_N}(x) - \eta_0(x)) \right| \right] = 0.$$

Proposition 5.2 (Hydrodynamic limit starting from the ergodic component). *For any $t \in [0, T]$, any $\delta > 0$ and any smooth test function $\varphi : \mathbb{T} \rightarrow \mathbb{R}$, we have*

$$\lim_{N \rightarrow \infty} \mathbb{P}_{\tilde{\mu}_N} \left(\left| \frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi \left(\frac{x}{N} \right) \eta_t(x) - \int_{\mathbb{T}} \varphi(u) \rho(t, u) du \right| > \delta \right) = 0 \quad (5.5)$$

where $\rho(t, u)$ is the unique solution of the hydrodynamic equation (2.9).

Before proving these two results, we give the following corollary of Lemma 5.1, which states that at time t_N , the macroscopic density profile has not changed. This will be useful to prove Proposition 5.2 in Section 5.3.

Corollary 5.3 (Macroscopic profile after the transience time). *For any smooth function $\varphi : \mathbb{T} \rightarrow \mathbb{R}$,*

$$\lim_{N \rightarrow \infty} \mathbb{E}_{\mu_N} \left[\left| \frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi \left(\frac{x}{N} \right) \eta_{t_N}(x) - \int_{\mathbb{T}} \varphi(u) \rho_0(u) du \right| \right] = 0.$$

This corollary is an immediate consequence of the definition of μ_N , together with the law of large numbers, and Lemma 5.1.

In the last two paragraphs of this section we prove Lemma 5.1 and Proposition 5.2.

5.2. Proof of Lemma 5.1. This proof is pretty elementary, since in a time $t_N = o_N(N)$, no particle in the system has time to travel on a microscopic distance much larger than $T_N = (\log N)^{32}$, and in particular no particle travels on a macroscopic (of order N) distance. Let us denote by $k = k(\eta) = \sum_{x \in \mathbb{T}_N} \eta_0(x)$ the random number of particles initially in \mathbb{T}_N . Let us also denote by $x_1(t), \dots, x_k(t) \in \mathbb{T}_N$ the *microscopic* positions of the particles at time t . The main result we need to prove is the following: for some well chosen $\beta > 0$

$$\mathbb{P}_{\nu_N} \left(\sup_{0 \leq i \leq k} |x_i(0) - x_i(t_N)| \geq T_N^\beta \right) \xrightarrow{N \rightarrow \infty} 0. \quad (5.6)$$

Before briefly proving this bound, let us show that Lemma 5.1 follows from it. By reorganizing the terms, we can rewrite the expectation in Lemma 5.1 as

$$\begin{aligned} & \mathbb{E}_{\nu_N} \left[\left| \frac{1}{N} \sum_{i=1}^k \left(\varphi \left(\frac{x_i(t_N)}{N} \right) - \varphi \left(\frac{x_i(0)}{N} \right) \right) \right| \right] \\ & \leq \sup_{\substack{x, y \in \mathbb{T}_N \\ |x-y| \leq T_N^\beta}} \left| \varphi \left(\frac{x}{N} \right) - \varphi \left(\frac{y}{N} \right) \right| + 2 \|\varphi\|_\infty \mathbb{P}_{\nu_N} \left(\sup_{1 \leq i \leq k} |x_i(t_N) - x_i(0)| \geq T_N^\beta \right). \end{aligned}$$

In the bound above, the first term vanishes because φ is a smooth function and because $T_N^\beta/N = o_N(1)$, whereas the second term vanishes according to (5.6).

We therefore only need to prove (5.6) to conclude the proof of Lemma 5.1. Since there are at most N particles in the system, the probability in (5.6) is less than

$$N \sum_{j \leq N} \mathbb{P}_{\nu_N}(k(\eta) = j) \sup_{i \leq j} \mathbb{P}_{\nu_N} \left(|x_i(0) - x_i(t_N)| \geq T_N^\beta \mid k(\eta) = j \right).$$

Since the particles jump to a neighboring site at rate at most N^2 , the probability above is less than the probability that a random walker jumping to its neighbors at rate N^2 jumped $X \geq T_N^\beta$ times in a time t_N , so that we only have to prove that

$$N \mathbb{P}(X \geq T_N^\beta) \xrightarrow{N \rightarrow \infty} 0. \quad (5.7)$$

This number of jumps X is distributed according to a Poisson law with mean T_N , so that by a standard argument of large deviations, one obtains that from some constant C ,

$$\mathbb{P}(X \geq T_N^\beta) \leq e^{-CT_N^{\beta-1}}.$$

We now only need to choose β large enough to obtain (5.7) and then conclude the proof of Lemma 5.1. \square

5.3. Proof of Proposition 5.2. In this section, we prove Proposition 5.2 by adapting the classical *entropy method* (cf. [11, Chapter 5]). Let us consider the process $\{(\eta_t(x))_{x \in \mathbb{T}_N} : t \in [0, T]\}$ started from the ergodic component under the measure $\tilde{\mu}_N$ defined in (5.1). For $s \in [0, T]$, we denote by $\langle \cdot, m_s^N \rangle$ the integral on $\{0, 1\}^{\mathbb{T}_N}$ w.r.t. the empirical measure at time s denoted by m_s^N , namely

$$\langle \psi, m_s^N \rangle = \frac{1}{N} \sum_{x \in \mathbb{T}_N} \psi \left(\frac{x}{N} \right) \eta_s(x), \quad \text{with } \psi : \mathbb{T} \rightarrow \mathbb{R}.$$

Fix a smooth function

$$\begin{aligned} \varphi : [0, T] \times \mathbb{T} & \rightarrow \mathbb{R} \\ (s, u) & \mapsto \varphi_s(u). \end{aligned}$$

The evolution of the process between time 0 and time T is described by the following Dynkin's formula:

$$\begin{aligned} \langle \varphi_T, m_T^N \rangle & = \langle \varphi_0, m_0^N \rangle + \int_0^T ds \frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi_s \left(\frac{x}{N} \right) N^2 \mathcal{L}_N \eta_s(x) \\ & \quad + \int_0^T ds \langle \partial_s \varphi_s, m_s^N \rangle + M_T^N(\varphi), \end{aligned}$$

where $(M_T^N(\varphi))$ is a $\mathbb{P}_{\tilde{\mu}_N}$ -martingale which vanishes⁶ in \mathbf{L}^2 , as $N \rightarrow \infty$. Furthermore, by summation by parts, we can also rewrite (thanks to the gradient property of our model, see Section 2.2.2) the first integrand above as

$$\frac{1}{N} \sum_{x \in \mathbb{T}_N} \varphi_s \left(\frac{x}{N} \right) N^2 \mathcal{L}_N \eta(x) = \frac{1}{N} \sum_{x \in \mathbb{T}_N} \Delta_x^N \varphi_s \tau_x h(\eta), \quad (5.8)$$

where h was defined in (2.6) and $\Delta^N \varphi$ denotes the discrete approximation of the continuous Laplacian defined as usual as

$$\Delta_x^N \varphi := N^2 \left(\varphi \left(\frac{x+1}{N} \right) + \varphi \left(\frac{x-1}{N} \right) - 2\varphi \left(\frac{x}{N} \right) \right).$$

The next step consists in replacing in the right hand side of (5.8) the local function $\tau_x h$ by its expectation with respect to the invariant measure parametrized by the empirical density in a mesoscopic box around x . For that purpose, let us define the empirical density of the configuration η in the box B_ℓ (of size $2\ell + 1$ symmetric around 0) as

$$\rho^\ell(0)(\eta) := \frac{1}{2\ell + 1} \sum_{y \in B_\ell} \eta(y),$$

and around x as $\rho^\ell(x)(\eta) := \rho^\ell(0)(\tau_x \eta)$. In what follows we will also use a simplified notation when looking at the macroscopic time s :

$$\rho_s^\ell(x) := \rho^\ell(x)(\eta_s).$$

The replacement in (5.8) is then done thanks to the *Replacement Lemma*, which in our case is stated as follows:

Lemma 5.4 (Replacement Lemma). *Recall that h has been defined in (2.6), and denote $F(\rho) = \max\{0, (2\rho - 1)/\rho\}$. We have for any $T > 0$*

$$\limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \mathbb{E}_{\tilde{\mu}_N} \left[\int_0^T ds \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x |V_{\varepsilon N}(h, \eta_s)| \right] = 0, \quad (5.9)$$

where, for any $k \in \mathbb{N}$,

$$V_k(h, \eta) := \frac{1}{2k + 1} \sum_{y \in B_k} \tau_y h(\eta) - F(\rho^k(0)(\eta)). \quad (5.10)$$

The standard proof of this lemma (see for example [11, Lemma 5.7 p.97 and Chapter 5, Section 3]) has to be adapted very carefully to our case. In particular, the grand canonical measures are different in the periodic and non-periodic setting, so that the former cannot be easily restricted to finite volume. Since the proof is quite long and technical, we postpone it to Section 7.

Thanks to Lemma 5.4, since φ is smooth, and since the martingale $(M_t^N(\varphi))$ vanishes in probability, we can write for any $\delta > 0$

$$\limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \mathbb{P}_{\tilde{\mu}_N} \left(\left| \langle \varphi_T, m_T^N \rangle - \langle \varphi_0, m_0^N \rangle - \int_0^T ds \left[\langle \partial_s \varphi_s, m_s^N \rangle + \frac{1}{N} \sum_{x \in \mathbb{T}_N} \Delta_x^N \varphi_s F(\rho_s^{\varepsilon N}(x)) \right] \right| > \delta \right) = 0. \quad (5.11)$$

⁶To see this, one can easily compute its quadratic variation, and shows that its expectation vanishes. See [11, Chapter 4, Page 76] for instance.

The quantity inside absolute values is in fact a function of the empirical measure process $\{m_s^N ; s \in [0, T]\}$: for any measure m on \mathbb{T} and any $u \in \mathbb{T}$, we set

$$\tilde{F}_u^\varepsilon(m) = F(\langle \delta_u^\varepsilon, m \rangle),$$

where δ_u^ε is defined as:

$$\delta_u^\varepsilon(v) = \frac{1}{2\varepsilon} \mathbf{1}_{\{u-\varepsilon \leq v \leq u+\varepsilon\}}, \quad \text{for any } v \in \mathbb{T}.$$

Note that, when $u = x/N$, since $\langle \delta_u^\varepsilon, m_s^N \rangle = \rho_s^{\varepsilon N}(x) + o_N(1)$, from (6.7) we have (uniformly in $s \in [0, T]$):

$$\tilde{F}_{x/N}^\varepsilon(m_s^N) = \max \left\{ 0, \frac{2\rho_s^{\varepsilon N}(x) - 1}{\rho_s^{\varepsilon N}(x)} \right\} + o_N(1).$$

Let us denote by Q_N the distribution of the empirical measure process $\{m_s^N ; s \in [0, T]\}$ when $\{\eta_s ; s \in [0, T]\}$ is distributed according to $\mathbb{P}_{\tilde{\mu}_N}$. We can now rewrite the previous identity (5.11) as

$$\limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} Q_N \left(\left| \langle \varphi_T, m_T^N \rangle - \langle \varphi_0, m_0^N \rangle - \int_0^T ds \left[\langle \partial_s \varphi_s, m_s^N \rangle + \frac{1}{N} \sum_{x \in \mathbb{T}_N} \Delta_x^N \varphi_s \tilde{F}_{x/N}^\varepsilon(m_s^N) \right] \right| > \delta \right) = 0. \quad (5.12)$$

It is straightforward, using the classical tools, to prove that the sequence (Q_N) is relatively compact for the weak topology so that for any of its limit points Q^* , we can write

$$\limsup_{\varepsilon \rightarrow 0} Q^* \left(\left| \langle \varphi_T, m_T \rangle - \langle \varphi_0, m_0 \rangle - \int_0^T ds \left[\langle \partial_s \varphi_s, m_s \rangle + \int_{\mathbb{T}} du \Delta \varphi_s(u) \tilde{F}_u^\varepsilon(m_s) \right] \right| > \delta \right) = 0.$$

Since at most one particle is allowed at any site, one can prove that any limit point Q^* is concentrated on measures m which admit at any time s a density ρ_s w.r.t. the Lebesgue measure on \mathbb{T} , so that we can let ε go to 0 in the last equality, to obtain thanks to Corollary 5.3 the following: for any $\delta > 0$,

$$Q^* \left(\left| \int_{\mathbb{T}} \varphi_T(u) \rho_T(u) du - \int_{\mathbb{T}} \varphi_0(u) \rho_0(u) du - \int_0^T ds \int_{\mathbb{T}} du \left[\partial_s \varphi_s(u) \rho_s(u) + \Delta \varphi_s(u) \max \left\{ 0, \frac{2\rho_s(u) - 1}{\rho_s(u)} \right\} \right] \right| > \delta \right) = 0.$$

Since a maximum principle holds for the hydrodynamic equation (2.9), letting δ go to 0 proves that Q^* is concentrated on trajectories whose density w.r.t. the Lebesgue measure is a weak solution of (2.9). The uniqueness of such solutions concludes the proof.

6. INVARIANT MEASURES AND EQUIVALENCE OF ENSEMBLES

6.1. Canonical measures for the exclusion dynamics in periodic setting.

Definition 6.1. For $k > N/2$, π_N^k is the uniform measure on Ω_N^k (recall (2.5)).

We know from [7] that π_N^k is the invariant measure for the exclusion process on \mathbb{T}_N with k particles. It is translation invariant and satisfies the *detailed balance condition*: for any $x \in \mathbb{T}_N$ and $\eta \in \Omega_N^k$,

$$\begin{aligned} & \pi_N^k(\eta)\eta(x-1)\eta(x)(1-\eta(x+1)) \\ &= \pi_N^k(\eta)\eta(x-1)\eta(x)(1-\eta(x+1))\eta(x+2) \\ &= \pi_N^k(\eta^{x,x+1})\eta^{x,x+1}(x-1)(1-\eta^{x,x+1}(x))\eta^{x,x+1}(x+1)\eta(x+2) \\ &= \pi_N^k(\eta^{x,x+1})(1-\eta^{x,x+1}(x))\eta^{x,x+1}(x+1)\eta(x+2), \end{aligned} \quad (6.1)$$

where we used the fact that the holes are isolated on Ω_N^k in the first and last equalities. Let us now characterize the marginals of π_N^k .

Lemma 6.1. *Let $k \in \{1, \dots, N-1\}$ and $m = N-k$. We have the identity*

$$|\Omega_N^k| = \binom{k}{m} + \binom{k-1}{m-1} = \frac{N}{k} \binom{k}{N-k}$$

Furthermore, fix $\ell \leq N$ and a local ergodic configuration $\sigma \in \widehat{\mathcal{E}}_{\Lambda_\ell}$. We define

- $p := \sum_{x \in \Lambda_\ell} \sigma(x) \in \{0, \dots, \ell\}$ its number of particles,
- $z := \ell - p$ its number of holes,

and, assume $p \leq k$ and $z \leq m$. Then we have

$$\left| \left\{ \eta \in \Omega_N^k : \eta|_{\Lambda_\ell} = \sigma \right\} \right| = \binom{k-p+\sigma(1)+\sigma(\ell)-1}{m-z}.$$

where by convention $\binom{n}{n'} = 0$ if $n' > n$.

Proof of Lemma 6.1. The proof is contained in [7]. We reproduce it here for completeness.

Split the configurations in Ω_N^k into those with a particle at 1 and those with a hole at 1. The number of the first kind is given by the number of ways one can insert the m holes between two particles. Since there are k particles which are situated on the torus, there are k possible positions for the holes and we must choose m of those. When there is a hole at 0, it remains to insert $m-1$ holes in the remaining $k-1$ positions.

To compute the cardinality of the second set, we consider that after fixing the configuration σ on Λ_ℓ , it remains to insert $m-z$ holes. The number of available positions is given by the remaining number of particles $k-p$ minus one if $\sigma(1) = \sigma(\ell) = 0$ (in that case the constraint imposes particles at $\ell+1$ and $0 = N$), and similarly for the other cases. \square

6.2. Grand canonical measures.

6.2.1. In infinite volume.

Definition 6.2. *Define the grand canonical measure π_ρ as the translation invariant measure on $\{0, 1\}^{\mathbb{Z}}$ such that*

- for $\rho > \frac{1}{2}$, $\ell \geq 1$,

$$\pi_\rho \left(\xi|_{\Lambda_\ell} = \sigma \right) = \mathbf{1}_{\{\sigma \in \widehat{\mathcal{E}}_{\Lambda_\ell}\}} (1-\rho) \left(\frac{1-\rho}{\rho} \right)^{\ell-1-p} \left(\frac{2\rho-1}{\rho} \right)^{2p-\ell+1-\sigma(1)-\sigma(\ell)}, \quad (6.2)$$

where $p = \sum_{x \in \Lambda_\ell} \sigma(x) \in \{0, \dots, \ell\}$ is the number of particles in σ .

- for $\rho \leq \frac{1}{2}$,

$$\pi_\rho = \frac{1}{2} \delta_{\circ\bullet} + \frac{1}{2} \delta_{\bullet\circ}, \quad (6.3)$$

where $\circ\bullet$ (resp. $\bullet\circ$) is the configuration in which there is a particle at x iff x is odd (resp. even), and δ_η is the usual Dirac measure concentrated on the configuration η .

Note that, by periodizing the configurations, we can see the measures π_N^k as measures on $\{0, 1\}^{\mathbb{Z}}$. In that case, since the state space is compact, the sequence $(\pi_N^k)_N$ is tight. If $k/N \rightarrow \rho > \frac{1}{2}$, there is a unique limit point. Using Lemma 6.1, one can check that this limit point is the measure π_ρ defined in Definition 6.2.

The formula (6.2) can be rewritten as

$$\pi_\rho \left(\xi_{|\Lambda_\ell} = \sigma \right) = \mathbf{1}_{\{\sigma \in \widehat{\mathcal{E}}_{\Lambda_\ell}\}} \kappa \alpha^\rho \beta^\ell \gamma^{\sigma(1)+\sigma(\ell)}, \quad (6.4)$$

where the four constants $\kappa, \alpha, \beta, \gamma$ all depend on $\rho \in (\frac{1}{2}, 1)$, and are given by

$$\kappa(\rho) := 2\rho - 1, \quad \alpha(\rho) := \frac{(2\rho - 1)^2}{\rho(1 - \rho)}, \quad \beta(\rho) := \frac{1 - \rho}{2\rho - 1} \quad \text{and} \quad \gamma(\rho) := \frac{\rho}{2\rho - 1}. \quad (6.5)$$

Whenever no confusion arises, we will omit their dependence on ρ .

Remark 6.2. Fix $\rho > \frac{1}{2}$. Note that $\pi_\rho(\xi(0)) = \rho$. Moreover, if one assumes $\ell \geq 1$ and takes $\sigma = (0, 1, \dots, 1)$, by applying the above formula one finds that

$$\pi_\rho \left(\xi_{|\Lambda_\ell} = (0, 1, \dots, 1) \right) = (1 - \rho) \left(\frac{2\rho - 1}{\rho} \right)^{\ell - 2}, \quad (6.6)$$

which, since $\frac{1 - \rho}{\rho} = 1 - \frac{2\rho - 1}{\rho}$, is nothing but

$$\pi_\rho(\xi(0) = 0) \times \mathbb{P} \left(\text{Geom} \left(\frac{1 - \rho}{\rho} \right) \geq \ell - 2 \right).$$

In fact, we can understand π_ρ as the measure which puts a hole in some position with probability $1 - \rho$, then puts one plus a random geometric number of parameter $\frac{1 - \rho}{\rho}$ particles on its right, then a hole, then starts again and so on.

The properties of the (non-product) measure π_ρ will be further investigated in later sections. In particular, we prove in Section 6.3 that for any density $\rho \in (\frac{1}{2}, 1)$, under π_ρ , the correlations between two boxes at a distance of order ℓ decay as $\exp(-C\ell)$. We also prove the *equivalence of ensembles*, stated in Corollary 6.10, which is a key ingredient to prove the Replacement Lemma, namely Lemma 5.4.

Remark 6.3. As a consequence of equation (6.2), for any $\rho \in (\frac{1}{2}, 1)$, there exist two constants $C = C(\rho)$ and $a = a(\rho) < 1$, such that for any configuration $\sigma \in \{0, 1\}^{\Lambda_\ell}$, $\pi_\rho(\xi_{|\Lambda_\ell} = \sigma) \leq Ca^\ell$.

Remark 6.4. Recall that we introduced the local function h in (2.6) as

$$h(\eta) = \eta(-1)\eta(0) + \eta(0)\eta(1) - \eta(-1)\eta(0)\eta(1),$$

and that we introduced in the Replacement Lemma 5.4 the function $F(\rho) = \min\{0, (2\rho - 1)/\rho\}$. Then, one easily checks that for any $\rho \in [\frac{1}{2}, 1]$

$$\pi_\rho(h) = F(\rho), \quad (6.7)$$

which is, as expected, the quantity appearing in the Laplacian in (2.9).

For $\rho > \frac{1}{2}$, we now introduce a periodic variant of the grand canonical measure π_ρ , which is defined on \mathbb{T}_N , is translation invariant, and is locally close in the limit $N \rightarrow \infty$ to π_ρ .

6.2.2. *In the periodic setting.* Let us fix $\rho > \frac{1}{2}$. For any periodic configuration $\eta \in \{0, 1\}^{\mathbb{T}_N}$, let us first define

$$\tilde{v}_{\rho,N}(\eta) := \frac{1}{N} \sum_{x \in \mathbb{T}_N} \pi_\rho \left(\xi_{\{1, \dots, N\}} = \tau_x \eta \right).$$

Note that, since to define the measure above we break the periodicity of the configuration η , one can find a non-ergodic periodic configuration $\eta \notin \mathcal{E}_N$ such that $\tilde{v}_{\rho,N}(\eta) > 0$, because, for some x , the non-periodic configuration $\sigma^x := \tau_x \eta$ belongs to $\widehat{\mathcal{E}}_{\Lambda_N}$. We therefore need to restrict $\tilde{v}_{\rho,N}$ to the periodic ergodic component \mathcal{E}_N as follows. For completeness, we also give a definition of the periodic grand canonical measure for densities below $\frac{1}{2}$.

Definition 6.3. For $\rho > \frac{1}{2}$, $\eta \in \mathbb{T}_N$,

$$v_{\rho,N}(\eta) := \frac{\mathbf{1}_{\{\eta \in \mathcal{E}_N\}}}{\tilde{v}_{\rho,N}(\mathcal{E}_N)} \tilde{v}_{\rho,N}(\eta). \quad (6.8)$$

For $\rho \leq \frac{1}{2}$: if N is even, $v_{\rho,N}$ is the uniform measure on the set of the two alternate configurations and if N is odd, $v_{\rho,N}$ is the uniform measure on the set (of cardinal N) of configurations with $(N-1)/2$ isolated holes.

We let the reader check the following straightforward properties:

- the probability measure $v_{\rho,N}$ is translation invariant;
- the probability measure $v_{\rho,N}$ is invariant for the dynamics generated by \mathcal{L}_N , since it satisfies the detailed balance condition (see (6.1));
- for $\rho > \frac{1}{2}$, we prove in Lemma 6.7 below that $\tilde{v}_{\rho,N}(\mathcal{E}_N) = \rho(2-\rho) + \mathcal{O}_N(e^{-CN})$. In particular, according to Remark 6.3, there exist two constants C and $0 < a < 1$ depending only on ρ such that for any $N \geq N_0(\rho)$ and any configuration $\eta \in \{0, 1\}^{\mathbb{T}_N}$

$$v_{\rho,N}(\eta) \leq Ca^N; \quad (6.9)$$

- for $\eta, \eta' \in \mathcal{E}_N$ two configurations with the same number of particles, we have $v_{\rho,N}(\eta) = v_{\rho,N}(\eta')$. In particular, for any $\rho \in (\frac{1}{2}, 1)$, any integer k satisfying $\frac{N}{2} < k < N$, we have $v_{\rho,N}(\cdot \mid \eta \in \Omega_N^k) = \pi_N^k(\cdot)$.

The rest of this section is dedicated to stating and proving several properties of the grand canonical measures (GCM) defined above. In Section 6.3, we prove that for any non-degenerate $\rho \in (\frac{1}{2}, 1]$, the correlations under the grand canonical measures π_ρ of two boxes at distance ℓ decay as $\exp(-C\ell)$ regardless of the size of the boxes. In Section 6.4, we show that the measure of the ergodic set $\tilde{v}_{\rho,N}(\mathcal{E}_N)$ converges exponentially fast to $\rho(2-\rho)$. In Section 6.5, we show that as N diverges, the periodic GCM $v_{\rho,N}$ locally resembles the infinite volume GCM π_ρ . Finally, the main result of this section is the so-called *equivalence of ensembles*, stated in Corollary 6.10 and proved in Section 6.7. These results are quite technical, we therefore elected to present a rather detailed proof in order to be as clear as possible.

6.3. **Decay of correlations for the GCM π_ρ .** We start by investigating the correlations of the infinite volume measure π_ρ . Let us define

$$P_\ell := \pi_\rho \left(\xi(0) = \xi(\ell) = 1 \right).$$

We first prove that the two-points correlations under the invariant measure π_ρ decay exponentially, which in particular yields a law of large numbers for the measure π_ρ .

Lemma 6.5. *For any $\rho \in (\frac{1}{2}, 1)$, there exist two constants $C = C(\rho) > 0$ and $q = q(\rho) > 0$ such that*

$$|P_\ell - \rho^2| \leq qe^{-C\ell}.$$

Proof of Lemma 6.5. To prove this result, we use equation (6.4). Let us recall the notation $\Lambda_\ell = \{1, \dots, \ell\}$ and rewrite

$$\begin{aligned} P_\ell &= \sum_{\sigma \in \{0,1\}^{\Lambda_{\ell-1}}} \pi_\rho \left(\xi(0) = 1, \xi_{|\Lambda_{\ell-1}} = \sigma, \xi(\ell) = 1 \right) \\ &= \sum_{\sigma \in \{0,1\}^{\Lambda_{\ell-2}}} \pi_\rho \left(\xi(0) = 1, \xi_{|\Lambda_{\ell-2}} = \sigma, \xi(\ell-1) = 1, \xi(\ell) = 1 \right) \\ &\quad + \sum_{\sigma \in \{0,1\}^{\Lambda_{\ell-2}}} \pi_\rho \left(\xi(0) = 1, \xi_{|\Lambda_{\ell-2}} = \sigma, \xi(\ell-1) = 0, \xi(\ell) = 1 \right) \\ &= \alpha(\rho)\beta(\rho) \sum_{\sigma \in \{0,1\}^{\Lambda_{\ell-2}}} \pi_\rho \left(\xi(0) = 1, \xi_{|\Lambda_{\ell-2}} = \sigma, \xi(\ell-1) = 1 \right) \\ &\quad + \alpha(\rho)\beta(\rho)\gamma(\rho) \sum_{\sigma \in \{0,1\}^{\Lambda_{\ell-2}}} \pi_\rho \left(\xi(0) = 1, \xi_{|\Lambda_{\ell-2}} = \sigma, \xi(\ell-1) = 0 \right) \\ &= \frac{2\rho-1}{\rho} \left[P_{\ell-1} + \frac{\rho}{2\rho-1} (\rho - P_{\ell-1}) \right]. \end{aligned}$$

This in turn yields

$$P_\ell = \rho^2 + (P_1 - \rho^2) \left(1 - \frac{1}{\rho}\right)^{\ell-1},$$

with $P_1 = \pi_\rho(\xi(0)\xi(1)) = 2\rho - 1$. This proves Lemma 6.5, since $\rho \in (\frac{1}{2}, 1)$. \square

We now use this lemma to state that the correlations under π_ρ of two boxes at distance ℓ from one another decays as $e^{-C\ell}$.

Corollary 6.6 (Correlation decay). *Fix $k \geq 1$, and shorten*

$$A = A(k) := \{-k+1, \dots, 0\} = \tau_{-k}\Lambda_k \quad \text{and} \quad B = B(\ell, k) = \{\ell+1, \dots, \ell+k\} = \tau_\ell\Lambda_k.$$

For any two configurations σ_1, σ_2 in $\{0, 1\}^A, \{0, 1\}^B$, any $\rho \in (\frac{1}{2}, 1)$, there exists $C := C(\rho) > 0$ such that

$$\frac{\pi_\rho(\xi_{|A} = \sigma_1, \xi_{|B} = \sigma_2)}{\pi_\rho(\xi_{|A} = \sigma_1)\pi_\rho(\xi_{|B} = \sigma_2)} = 1 + \mathcal{O}_\ell(e^{-C\ell}),$$

where the $\mathcal{O}_\ell(e^{-C\ell})$ depends on ρ , but can be bounded uniformly in k and σ_1, σ_2 .

Proof of Corollary 6.6. To prove this result, first notice that for any two neighboring sets $A_1 = \{a, \dots, b\}, A_2 = \{b+1, \dots, c\}$, and any given two configurations σ_1, σ_2 on these sets, we can write thanks to the explicit formula (6.4)

$$\pi_\rho \left(\xi_{|A_1} = \sigma_1, \xi_{|A_2} = \sigma_2 \right) = \frac{\mathbf{1}_{\{\sigma_1(b) + \sigma_2(b+1) \geq 1\}}}{\kappa \gamma^{\sigma_1(b) + \sigma_2(b+1)}} \pi_\rho \left(\xi_{|A_1} = \sigma_1 \right) \pi_\rho \left(\xi_{|A_2} = \sigma_2 \right).$$

Note in particular that if the configuration (σ_1, σ_2) is not ergodic, then either σ_1 is not, or σ_2 is not, or $\sigma_1(b) + \sigma_2(b+1) = 0$, so that the right hand side above vanishes.

Getting back to our problem, summing over all the possible configurations σ in Λ_ℓ , and using the identity above (since A, Λ_ℓ and B are neighboring sets), we can write for any σ_1, σ_2 respectively in $\{0, 1\}^A, \{0, 1\}^B$,

$$\begin{aligned} \pi_\rho \left(\xi|_A = \sigma_1, \xi|_B = \sigma_2 \right) &= \pi_\rho \left(\xi|_A = \sigma_1 \right) \pi_\rho \left(\xi|_B = \sigma_2 \right) \\ &\times \sum_{\sigma \in \{0,1\}^{\Lambda_\ell}} \frac{\mathbf{1}_{\{\sigma_1(0)+\sigma(1) \geq 1\}} \mathbf{1}_{\{\sigma(\ell)+\sigma_2(\ell+1) \geq 1\}}}{\kappa^2 \gamma^{\sigma_1(0)+\sigma(1)+\sigma(\ell)+\sigma_2(\ell+1)}} \pi_\rho \left(\xi|_{\Lambda_\ell} = \sigma \right). \end{aligned}$$

Thanks to this identity, proving Corollary 6.6 is a matter of verification.

Start for example by assuming that $\sigma_1(0) = \sigma_2(\ell+1) = 1$, both indicator functions in the sum of the right hand side are always 1, and since $\kappa\gamma = \rho$, we can rewrite the last identity as

$$\pi_\rho \left(\xi|_A = \sigma_1, \xi|_B = \sigma_2 \right) = \frac{\pi_\rho \left(\xi|_A = \sigma_1 \right) \pi_\rho \left(\xi|_B = \sigma_2 \right)}{\rho^2} \pi_\rho \left(\gamma^{-\xi(1)-\xi(\ell)} \right).$$

Elementary computations and Lemma 6.5 yield $\pi_\rho \left(\gamma^{-\xi(1)-\xi(\ell)} \right) = \rho^2 + \mathcal{O}_\ell(e^{-C\ell})$, thus concluding this case. Since it is a little different, we now consider the case $\sigma_1(0) = \sigma_2(\ell+1) = 0$. In this case, for the indicator functions not to vanish, we must have a particle at each extremity of Λ_ℓ , so that

$$\begin{aligned} \pi_\rho \left(\xi|_A = \sigma_1, \xi|_B = \sigma_2 \right) &= \pi_\rho \left(\xi|_A = \sigma_1 \right) \pi_\rho \left(\xi|_B = \sigma_2 \right) \\ &\times \sum_{\sigma \in \{0,1\}^{\Lambda_\ell \setminus \{1,\ell\}}} \frac{1}{\kappa^2 \gamma^2} \pi_\rho \left(\xi|_{\Lambda_\ell} = (1, \sigma, 1) \right) \\ &= \pi_\rho \left(\xi|_A = \sigma_1 \right) \pi_\rho \left(\xi|_B = \sigma_2 \right) \frac{\pi_\rho \left(\xi(1) = \xi(\ell) = 1 \right)}{\rho^2}. \end{aligned}$$

Therefore, Lemma 6.5 proves this case as well. The last case $\sigma_1(0) \neq \sigma_2(\ell+1)$ is treated in the same way. \square

6.4. Measure of the ergodic set. We now prove that the measure of the ergodic set $\tilde{\nu}_{\rho,N}(\mathcal{E}_N)$ is exponentially close to $\rho(2-\rho)$.

Lemma 6.7. *Fix $\rho \in (\frac{1}{2}, 1]$. Then there exists $C = C(\rho) > 0$ such that*

$$\tilde{\nu}_{\rho,N}(\mathcal{E}_N) = \pi_\rho \left((\xi(1), \xi(N)) \neq (0,0) \right) + \mathcal{O}_N(e^{-CN}) = \rho(2-\rho) + \mathcal{O}_N(e^{-CN}). \quad (6.10)$$

Note that $\rho(2-\rho) > \frac{3}{4}$.

Proof. For any $1 \leq x \leq N$ and any configuration ξ in Λ_N , denote by ξ^x the *periodic* configuration on \mathbb{T}_N obtained by injecting Λ_N in \mathbb{T}_N , and then translating by $x \in \mathbb{T}_N$. By definition,

$$\begin{aligned} \tilde{\nu}_{\rho,N}(\mathcal{E}_N) &= \sum_{\eta \in \mathcal{E}_N} \frac{1}{N} \sum_{x \in \mathbb{T}_N} \pi_\rho \left(\xi|_{\{1,\dots,N\}} = \tau_x \eta \right) \\ &= \frac{1}{N} \sum_{x \in \mathbb{T}_N} \pi_\rho \left(\xi^x|_{\{1,\dots,N\}} \in \mathcal{E}_N \right) \\ &= \pi_\rho \left((\xi(1), \xi(N)) \neq (0,0), \sum_{y=1}^N \xi(y) > \frac{N}{2} \right) \end{aligned}$$

$$\begin{aligned}
&= \pi_\rho \left((\xi(1), \xi(N)) \neq (0, 0) \right) \\
&\quad - 2\pi_\rho \left(\forall i \in \{1, \dots, N\}, \xi(i) = 0 \text{ iff } i \text{ even and } \xi(1)\xi(N) \neq 0 \right),
\end{aligned}$$

where the last equality follows from the fact that a.s. under π_ρ holes are isolated. What this says is that the only way for $\xi_{\{1, \dots, N\}}$ not to be in \mathcal{E}_N is to be one of the two configurations alternating holes and particles (which may only happen if N is even), or to have $\xi(1) = \xi(N) = 0$. Therefore, using the notations introduced in Sections 6.2 and 6.3 and recalling Remark 6.3, there exists $a < 1$ such that

$$\begin{aligned}
\tilde{v}_{\rho, N}(\mathcal{E}_N) &= \pi_\rho \left(\xi(1)(1 - \xi(N)) + \xi(N)(1 - \xi(1)) + \xi(1)\xi(N) \right) + \mathcal{O}_N(a^N) \\
&= P_{N-1} + 2(\rho - P_{N-1}) + \mathcal{O}_N(a^N)
\end{aligned}$$

The conclusion follows from Lemma 6.5. \square

6.5. Local convergence of the periodic GCM $v_{\rho, N}$ to the infinite volume GCM π_ρ . In this section, we show that as N diverges, the periodic GCM $v_{\rho, N}$ is locally close to the infinite volume GCM π_ρ . More precisely, we have the following result:

Lemma 6.8. *Let $f : \{0, 1\}^{\mathbb{Z}} \rightarrow \mathbb{R}$ be a local function and assume without any loss of generality that its support is included in B_ℓ , for some fixed $\ell \in \mathbb{N}$. Then for any sequence (a_N) such that $\log N = o_N(a_N)$ and $a_N = o_N(N)$, there exists $C = C(\ell, \rho) > 0$ such that*

$$\left| v_{\rho, N}(f) - \pi_\rho(f) \right| \leq C \|f\|_\infty \frac{a_N}{N}. \quad (6.11)$$

Furthermore, if $\rho > \frac{1}{2}$, for any integer $\ell \geq 1$, and any two configurations σ_1 and σ_2 in $\{0, 1\}^{B_\ell}$,

$$\begin{aligned}
\limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \sup_{\sqrt{N} < |y| < \varepsilon N} \left| v_{\rho, N} \left(\xi_{|B_\ell} = \sigma_1, \xi_{|\tau_y B_\ell} = \sigma_2 \right) \right. \\
\left. - \pi_\rho \left(\xi_{|B_\ell} = \sigma_1 \right) \pi_\rho \left(\xi_{|B_\ell} = \sigma_2 \right) \right| = 0, \quad (6.12)
\end{aligned}$$

uniformly in σ_1, σ_2 .

Proof of Lemma 6.8. We start by proving (6.11) for $\rho > \frac{1}{2}$, (6.12) is proved using the same tools. Fix a local function f with support included in B_ℓ , for some fixed $\ell \in \mathbb{N}$. One can write f as a linear combination of functions of the form $f(\xi) = \mathbf{1}_{\{\xi_{|B_\ell} = \sigma\}}$, where $\sigma \in \{0, 1\}^{B_\ell}$, therefore we only need to prove the result for functions of this form. For such a function, by definition

$$\begin{aligned}
v_{\rho, N}(f) &= \sum_{\substack{\eta \in \mathcal{E}_N \\ \eta_{|B_\ell} = \sigma}} v_{\rho, N}(\eta) \\
&= \frac{1}{\tilde{v}_{\rho, N}(\mathcal{E}_N)} \sum_{\substack{\eta \in \mathcal{E}_N \\ \eta_{|B_\ell} = \sigma}} \frac{1}{N} \sum_{x=1}^N \pi_\rho \left(\xi_{\{1, \dots, N\}} = \tau_x \eta \right).
\end{aligned}$$

To proceed, we use the decoupling property of π_ρ (cf. Corollary 6.6). We divide the last sum that appears above into two sums:

- one over $A_N := \{a_N, \dots, N - a_N\}$,
- and the other one over the complement, namely $\{1, \dots, N\} \setminus A_N$.

From now on we assume N sufficiently large such that $\ell < a_N$. Let us treat the first sum: for any $x \in A_N$,

$$\begin{aligned} & \sum_{\substack{\eta \in \mathcal{E}_N \\ \eta|_{B_\ell} = \sigma}} \pi_\rho \left(\xi|_{\{1, \dots, N\}} = \tau_x \eta \right) \\ &= \pi_\rho \left(\xi|_{\{x-\ell, \dots, x+\ell\}} = \sigma \quad \text{and} \quad (\xi(1), \xi(N)) \neq (0, 0) \right) + \mathcal{O}_N(e^{-CN}) \\ &= \pi_\rho \left(\xi|_{\{x-\ell, \dots, x+\ell\}} = \sigma \right) \pi_\rho \left((\xi(1), \xi(N)) \neq (0, 0) \right) + \mathcal{O}_N(e^{-CN} + e^{-Ca_N}) \\ &= \pi_\rho \left(\xi|_{\{x-\ell, \dots, x+\ell\}} = \sigma \right) \tilde{\nu}_{\rho, N}(\mathcal{E}_N) + \mathcal{O}_N(e^{-CN} + e^{-Ca_N}), \end{aligned}$$

where $C = C(\rho, \ell) > 0$. In the identity above, the term $\mathcal{O}_N(e^{-CN})$ comes from the fact that alternate configurations like $(1, 0, \dots, 1, 0)$ and $(0, 1, \dots, 0, 1)$ have positive probability under π_ρ but do not appear in \mathcal{E}_N , which excludes configurations with density $\frac{1}{2}$ or less. The first equality follows from the isolation of holes in π_ρ (recall the proof of Lemma 6.7) and the last equality from Corollary 6.6. Indeed, for any $x \in A_N$, the points belonging to $\{x - \ell, \dots, x + \ell\}$ are at distance at least $a_N - \ell - 1$ from both points 1 and N . Therefore, one can write the first sum over A_N as (changing the constants in the exponentials)

$$\begin{aligned} & \frac{1}{N \tilde{\nu}_{\rho, N}(\mathcal{E}_N)} \sum_{x \in A_N} \sum_{\substack{\eta \in \mathcal{E}_N \\ \eta|_{B_\ell} = \sigma}} \pi_\rho \left(\xi|_{\{1, \dots, N\}} = \tau_x \eta \right) \\ &= \frac{1}{N} \sum_{x \in A_N} \pi_\rho \left(\xi|_{\{x-\ell, \dots, x+\ell\}} = \sigma \right) + \mathcal{O}_N(e^{-CN} + e^{-Ca_N}) \\ &= \pi_\rho(f) + \mathcal{O}_N \left(e^{-CN} + e^{-Ca_N} + \frac{a_N}{N} \right). \end{aligned} \tag{6.13}$$

Note that we used the fact that $\tilde{\nu}_{\rho, N}(\mathcal{E}_N)$ is bounded away from 0 (from Lemma 6.7) and the translation invariance of π_ρ .

Let us now turn to the second sum over $\{1, \dots, N\} \setminus A_N$: for any x , we can bound

$$\sum_{\substack{\eta \in \mathcal{E}_N \\ \eta|_{B_\ell} = \sigma}} \pi_\rho \left(\xi|_{\{1, \dots, N\}} = \tau_x \eta \right) \leq \sum_{\eta \in \mathcal{E}_N} \pi_\rho \left(\xi|_{\{1, \dots, N\}} = \tau_x \eta \right) = \tilde{\nu}_{\rho, N}(\mathcal{E}_N).$$

Consequently, the second sum can be easily bounded as follows:

$$\frac{1}{N \tilde{\nu}_{\rho, N}(\mathcal{E}_N)} \sum_{x \in \{1, \dots, N\} \setminus A_N} \sum_{\substack{\eta \in \mathcal{E}_N \\ \eta|_{B_\ell} = \sigma}} \pi_\rho \left(\xi|_{\{1, \dots, N\}} = \tau_x \eta \right) \leq \frac{2a_N}{N}.$$

This proves (6.11).

We simply sketch the proof of the second identity (6.12), which relies on the same steps, together with the decoupling property for π_ρ (cf. Corollary 6.6). For any given $\varepsilon > 0$ (which will tend to 0), we can set $a_N = 3\varepsilon N$, and repeat the same proof as before. We obtain for any fixed configurations σ^1 , and σ^2 on B_ℓ , and for any $2\ell + 2 < |y| < \varepsilon N$, that

$$\nu_{\rho, N} \left(\xi|_{B_\ell} = \sigma^1, \xi|_{\tau_y B_\ell} = \sigma^2 \right) = \pi_\rho \left(\xi|_{B_\ell} = \sigma^1, \xi|_{\tau_y B_\ell} = \sigma^2 \right) + o_\varepsilon(1) + o_N(1).$$

We can then use the decay of correlations proved for π_ρ in Corollary 6.6 and translation invariance of π_ρ , to obtain that for any fixed configurations σ^1, σ^2 on B_ℓ ,

$$\limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \sup_{\sqrt{N} < |y| < \varepsilon N} \left| v_{\rho, N} \left(\xi_{|B_\ell} = \sigma^1, \xi_{|\complement_y B_\ell} = \sigma^2 \right) - \pi_\rho \left(\xi_{|B_\ell} = \sigma^1 \right) \pi_\rho \left(\xi_{|B_\ell} = \sigma^2 \right) \right| = 0,$$

uniformly in σ^1, σ^2 as wanted.

To conclude the proof of Lemma 6.8, recall that for $\rho \leq \frac{1}{2}$, we defined $v_{\rho, N}$ as the uniform measure on the set of the two alternate configurations if N is even; and on the set of configurations with $(N-1)/2$ isolated holes if N is odd. With this observation, (6.11) is a matter of simple verification. \square

6.6. Measures on a finite box. We now define the projections of the measures $\pi_\rho, v_{\rho, N}$ on a finite box B_ℓ .

Recall the explicit expression (6.4) for the GCM π_ρ , which straightforwardly translates to B_ℓ instead of Λ_ℓ , by translation invariance. For any $\ell \in \mathbb{N}$, let us define the probability measure $\widehat{\pi}_\rho^\ell$ on $\widehat{\mathcal{E}}_{B_\ell}$ as

$$\widehat{\pi}_\rho^\ell(\sigma) := \pi_\rho \left(\xi_{|B_\ell} = \sigma \right) = \kappa(\rho) \alpha(\rho)^p \beta(\rho)^{2\ell+1} \gamma(\rho)^{\sigma(-\ell) + \sigma(\ell)}, \quad (6.14)$$

where $p = \sum_{x \in B_\ell} \sigma(x)$. Then, for any integer $j \in \{\ell, \dots, 2\ell+1\}$ we denote by $\widehat{\pi}_\rho^{\ell, j}$ the measure $\widehat{\pi}_\rho^\ell$ conditioned on having j particles. Namely, if we denote by $\widehat{\Omega}_\ell^j$ the hyperplane

$$\widehat{\Omega}_\ell^j := \left\{ \sigma \in \widehat{\mathcal{E}}_{B_\ell} : \sum_{x \in B_\ell} \sigma(x) = j \right\},$$

then $\widehat{\pi}_\rho^{\ell, j}$ is the probability measure on $\widehat{\Omega}_\ell^j$ such that

$$\widehat{\pi}_\rho^{\ell, j}(\sigma) = \frac{\widehat{\pi}_\rho^\ell(\sigma)}{\widehat{\pi}_\rho^\ell(\widehat{\Omega}_\ell^j)}, \quad \text{for any } \sigma \in \widehat{\Omega}_\ell^j. \quad (6.15)$$

Similarly, we define the probability measure $\widehat{v}_{\rho, N}^\ell$ on $\widehat{\mathcal{E}}_{B_\ell}$ as

$$\widehat{v}_{\rho, N}^\ell(\sigma) = v_{\rho, N}(\eta_{|B_\ell} = \sigma). \quad (6.16)$$

Note in particular that according to Lemma 6.8, for any sequence (a_N) such that $\log N = o_N(a_N)$ and $a_N = o_N(N)$, one can write

$$\sup_{\sigma \in \widehat{\mathcal{E}}_{B_\ell}} \left| \widehat{v}_{\rho, N}^\ell(\sigma) - \widehat{\pi}_\rho^\ell(\sigma) \right| \leq \frac{C(\ell, \rho) a_N}{N}. \quad (6.17)$$

6.7. Equivalence of ensembles. We now state and prove the main result for this section, namely the equivalence of ensembles, which states that far enough from the extremities $-\ell$ and ℓ , the canonical measure $\widehat{\pi}_\rho^{\ell, j}$ defined in the previous paragraph is locally close to the GCM $\pi_{j/(2\ell+1)}$ with parameter the empirical density $j/(2\ell+1)$.

Let $k \in \mathbb{N}$ be fixed, define $B_k(x) = \{x-k, \dots, x+k\}$ (in particular $B_k = B_k(0)$) and

$$\rho_\ell = \rho_\ell(j) := \frac{j}{2\ell+1},$$

the density in B_ℓ under $\widehat{\pi}_\rho^{\ell, j}$. Furthermore, let us introduce

$$E_\ell = \{\ell, \dots, 2\ell+1\}, \quad \text{and} \quad E_\ell(\delta) = \{\lceil \ell(1+\delta) \rceil, \dots, \lfloor (1-\delta)(2\ell+1) \rfloor\},$$

which are respectively the set of possible particle numbers j in B_ℓ , and the same set, excluding the densities $\rho_\ell(j)$ which are at distance less than δ to the critical densities $\frac{1}{2}$ and 1. Note that $E_\ell = E_\ell(0)$ and $E_\ell(\delta) \subset E_\ell$ for any $\delta > 0$.

Proposition 6.9. *Fix a configuration $\sigma \in \{0, 1\}^{B_k}$. Then, for any $\rho \in (\frac{1}{2}, 1)$, any $\delta_1, \delta_2 > 0$, there exists a constant $C_0 := C_0(k, \sigma, \delta_1, \delta_2)$ such that for any $\ell > 0$,*

$$\max_{\substack{j \in E_\ell(\delta_1) \\ x \in B_{(1-\delta_2)\ell-k}}} \left| \widehat{\pi}_\rho^{\ell,j}(\varsigma_{B_k(x)} = \sigma) - \pi_{\rho_\ell(j)}(\xi_{B_k} = \sigma) \right| \leq \frac{C_0}{\ell^{1/4}}. \quad (6.18)$$

Furthermore, for some constant $C_1 := C_1(k, \delta_2)$,

$$\limsup_{\ell \rightarrow \infty} \max_{\substack{j \in E_\ell \setminus E_\ell(\delta_1) \\ x \in B_{(1-\delta_2)\ell-k}}} \left| \widehat{\pi}_\rho^{\ell,j}(\varsigma_{B_k(x)} = \sigma) - \pi_{\rho_\ell(j)}(\xi_{B_k} = \sigma) \right| \leq C_1 \delta_1. \quad (6.19)$$

Note that $B_k(x)$ splits B_ℓ in three parts: a left cluster, $B_k(x)$ itself, and a right cluster. The condition $x \in B_{(1-\delta_2)\ell-k}$ ensures that both left and right clusters are of size at least $\delta_2\ell$. Because we need an estimate which is uniform in the number of particles and in the positions at which the box is taken, the proof of this result is quite technical. To prove the Replacement Lemma 5.4, we will use the following corollary of Proposition 6.9.

Corollary 6.10 (Equivalence of ensembles). *For any local function $f : \{0, 1\}^{\mathbb{Z}} \rightarrow \mathbb{R}$, any $\rho \in (\frac{1}{2}, 1]$, and any $\delta > 0$, we have*

$$\lim_{\ell \rightarrow \infty} \max_{\substack{j \in E_\ell \\ x \in B_{(1-\delta)\ell}}} \left| \widehat{\pi}_\rho^{\ell,j}(\tau_x f) - \pi_{\rho_\ell(j)}(f) \right| = 0. \quad (6.20)$$

We start by briefly proving that Proposition 6.9 yields Corollary 6.10.

Proof of Corollary 6.10. Fix $\rho \in (\frac{1}{2}, 1)$ and a local function f . There exists an integer k such that f only depends on sites in B_k . Then, if we choose $\delta_2 = \delta/2$, for any $\ell > 2k/\delta$, we have $B_{(1-\delta)\ell} \subset B_{(1-\delta_2)\ell-k}$. Then, for ℓ large enough, we can use the triangular inequality and Proposition 6.9 to write for any $\delta_1 > 0$, $x \in B_{(1-\delta)\ell}$, $j \in E_\ell$,

$$\begin{aligned} & \left| \widehat{\pi}_\rho^{\ell,j}(\tau_x f) - \pi_{\rho_\ell(j)}(f) \right| \\ & \leq 2^{2k+1} \|f\|_\infty \max_{\substack{j \in E_\ell \\ \sigma \in \{0,1\}^{B_k}}} \left| \widehat{\pi}_\rho^{\ell,j}(\varsigma_{B_k(x)} = \sigma) - \pi_{\rho_\ell(j)}(\xi_{B_k} = \sigma) \right| \\ & \leq 2^{2k+1} \|f\|_\infty \left(\ell^{-1/4} \max_{\sigma \in \{0,1\}^{B_k}} \{C_0(k, \sigma, \delta_1, \delta_2)\} + (1 + o_\ell(1)) \delta_1 C_1(k, \delta_2) \right). \end{aligned}$$

The last estimate has been obtained after writing $E_\ell = E_\ell(\delta_1) \cup (E_\ell \setminus E_\ell(\delta_1))$ and applying both results of Proposition 6.9. Since k is fixed, we then let $\ell \rightarrow \infty$, then $\delta_1 \rightarrow 0$ to prove Corollary 6.10. \square

We now prove Proposition 6.9.

Proof of Proposition 6.9. We start by the case where the fixed density $\rho_\ell(j)$ is close to the extreme values $\frac{1}{2}$ or 1 (i.e. $j \in E_\ell \setminus E_\ell(\delta_1)$).

PROOF OF (6.19). We will only detail the case $\rho_\ell(j) \geq 1 - \delta_1$, since the case $\rho_\ell(j) \leq (1 + \delta_1)/2$ is treated in the same way. Denote by $\mathbf{1}_k$ the constant configuration on B_k with one particle at each site, and no empty site. To prove (6.19), it is sufficient to show that for some constant $C_2 := C_2(k, \delta_2)$, we have for any $j \in E_\ell \setminus E_\ell(\delta_1)$ satisfying $\rho_\ell(j) \geq 1 - \delta_1$,

$$\pi_{\rho_\ell(j)}(\xi_{|B_k} = \mathbf{1}_k) \geq 1 - C_2 \delta_1,$$

and for any $x \in B_{(1-\delta_2)\ell-k}$,

$$\widehat{\pi}_\rho^{\ell,j}(\varsigma_{|B_k(x)} = \mathbf{1}_k) \geq 1 - C_2 \delta_1, \quad (6.21)$$

and then let $C_1(k, \delta_2) := 2C_2(k, \delta_2)$. The first bound is an immediate consequence of the explicit formula (6.2) for π_ρ . From now on, for $\varepsilon_1, \varepsilon_2 \in \{0, 1\}$, we denote by

$$\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell,j}(\cdot) = \widehat{\pi}_\rho^{\ell,j}(\cdot \mid \zeta(-\ell) = \varepsilon_1, \zeta(\ell) = \varepsilon_2),$$

the measure $\widehat{\pi}_\rho^{\ell,j}$ conditioned to be in the state ε_1 (resp. ε_2) at site $-\ell$ (resp. ℓ). Then, to prove (6.21), it is clearly sufficient to prove for any $\varepsilon_1, \varepsilon_2 \in \{0, 1\}$

$$\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell,j}(\varsigma_{|B_k(x)} = \mathbf{1}_k) \geq 1 - C_2 \delta_1.$$

Furthermore, to prove the latter, it is sufficient to prove that for some constant $C(k, \delta_2)$, for any $\sigma \neq \mathbf{1}_k \in \{0, 1\}^{B_k}$

$$\frac{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell,j}(\varsigma_{|B_k(x)} = \sigma)}{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell,j}(\varsigma_{|B_k(x)} = \mathbf{1}_k)} \leq C(k, \delta_2) \delta_1.$$

Indeed, summing this bound above over $\sigma \in \{0, 1\}^{B_k} \setminus \{\mathbf{1}_k\}$, yields

$$\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell,j}(\varsigma_{|B_k(x)} = \mathbf{1}_k) \geq \frac{1}{1 + 2^{|B_k|} C(k, \delta_2) \delta_1} \geq 1 - C_2(k, \delta_2) \delta_1,$$

as wanted. We therefore state the result we need as a separate lemma, the proof of which requires elements of the proof of (6.18).

Lemma 6.11. *There exists a constant $C = C(k, \delta_2)$ such that for any non-full configuration $\sigma \in \{0, 1\}^{B_k} \setminus \{\mathbf{1}_k\}$, any $j \geq (1 - \delta_1)(2\ell + 1)$ and any $\varepsilon_1, \varepsilon_2 \in \{0, 1\}$*

$$\limsup_{\ell \rightarrow \infty} \frac{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell,j}(\varsigma_{|B_k(x)} = \sigma)}{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell,j}(\varsigma_{|B_k(x)} = \mathbf{1}_k)} \leq C(k, \delta_2) \delta_1. \quad (6.22)$$

The same is true if $j \leq (1 + \delta_1)\ell$, and with $\{\varsigma_{|B_k(x)} = \mathbf{1}_k\}$ replaced by $\{\varsigma_{|B_k(x)} \in \mathbf{1}/\mathbf{2}_k\}$, where $\mathbf{1}/\mathbf{2}_k$ represents the set which contains the two alternate configurations $\circ\bullet$ and $\bullet\circ$ with respectively $k, k+1$ particles.

In order to introduce the notations and estimates needed to prove this result, we postpone the proof of Lemma 6.11 to the end of the section, and for now prove the first estimate (6.18):

PROOF OF (6.18). Fix some particle number $j \in E_\ell(\delta_1)$, we drop to simplify our notation the dependency in j and simply write $\rho_\ell \in [\frac{1}{2} + \delta_1, 1 - \delta_1]$ for the density in B_ℓ . For any $\ell > \ell_0$, (where ℓ_0 only depends on k) we must have for $\sigma \in \widehat{\mathcal{O}}_{B_k}$

$$\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell,j}(\varsigma_{|B_k(x)} = \sigma) > 0 \quad \text{and} \quad \pi_{\rho_\ell}(\xi_{|B_k} = \sigma) > 0. \quad (6.23)$$

Indeed, the first statement holds because, since ρ_ℓ is bounded away from 1 and $\frac{1}{2}$, for any $j \in E_\ell(\delta_1)$, one will always be able to find a configuration ς in Ω_ℓ^j such that $\varsigma_{|B_k(x)} = \sigma$. The second statement holds because ρ_ℓ is neither $\frac{1}{2}$ nor 1.

From (6.23) we can write

$$\begin{aligned}
& \left| \widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma) - \pi_{\rho_\ell}(\xi_{|B_k} = \sigma) \right| \\
& \leq \left| \frac{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma) - \pi_{\rho_\ell}(\xi_{|B_k} = \sigma)}{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma) \pi_{\rho_\ell}(\xi_{|B_k} = \sigma)} \right| \widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma) \\
& \leq \left| \frac{1}{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma)} - \frac{1}{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma)} \right| \widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma) \\
& = \left| \sum_{\sigma' \in \{0,1\}^{B_k}} \frac{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma')}{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma)} - \frac{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma')}{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma)} \right| \widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma) \\
& \leq 2^{2k+1} \max_{\sigma' \in \{0,1\}^{B_k}} \left| \frac{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma')}{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma)} - \frac{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma')}{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma)} \right| \widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma).
\end{aligned}$$

The interest of this elementary computation is that the quotient of the $\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma)$'s is far easier to estimate than the $\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma)$'s themselves, because the normalizing constants

$$\pi_\rho \left(\sum_{x \in B_\ell} \xi(x) = j, \xi(-\ell) = \varepsilon_1, \xi(\ell) = \varepsilon_2 \right)$$

cancel out. The first estimate (6.18) is therefore a consequence of Lemma 6.12 below, by defining

$$C_0(k, \sigma, \delta_1, \delta_2) = 2^{2k+1} \max_{\sigma' \in \{0,1\}^{B_k}} C_3(k, \sigma, \sigma', \delta_1, \delta_2),$$

thus proving Proposition 6.9. \square

Lemma 6.12. *Fix $k \in \mathbb{N}$ as well as two configurations $\sigma, \sigma' \in \{0,1\}^{B_k}$. For any $\rho \in (\frac{1}{2}, 1)$, any $\varepsilon_1, \varepsilon_2 \in \{0,1\}$, and any positive δ_1, δ_2 , there exists a constant $C_3 := C_3(k, \sigma, \sigma', \delta_1, \delta_2)$ such that*

$$\max_{\substack{j \in E_\ell(\delta_1) \\ x \in B_{(1-\delta_2)\ell-k}}} \left| \frac{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma')}{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma)} - \frac{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma')}{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma)} \right| \leq \frac{C_3}{\widehat{\pi}_{\rho, \varepsilon_1, \varepsilon_2}^{\ell, j}(\varsigma_{|B_k(x)} = \sigma) \ell^{1/4}}. \quad (6.24)$$

Proof of Lemma 6.12. Denote by $p(\sigma)$ the number of particles in σ . Then, thanks to the explicit formula (6.2) for π_ρ , we can rewrite the second term in the left hand side of (6.24) as

$$\frac{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma')}{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma)} = \alpha(\rho_\ell)^{p(\sigma) - p(\sigma')} \gamma(\rho_\ell)^{\sigma'(-k) + \sigma'(k) - \sigma(-k) - \sigma(k)}, \quad (6.25)$$

where α and γ were introduced before in (6.5), recall

$$\alpha(\rho) = \frac{(2\rho - 1)^2}{\rho(1 - \rho)} \quad \text{and} \quad \gamma(\rho) = \frac{\rho}{2\rho - 1}.$$

We now estimate the first part in the left hand side of (6.24). Once again, we will only write the case $\varepsilon_1 = \varepsilon_2 = 1$, the others can be treated similarly. Recall that $x \in B_{(1-\delta_2)\ell-k}$, and define

$$\ell_1 := \ell_1(x, \ell, k) = \ell - k + x \geq \delta_2 \ell, \quad \ell_2 := \ell_1(x, \ell, k) = \ell - k - x \geq \delta_2 \ell,$$

which are the respective sizes of the clusters to the left and right of $B_k(x)$ in B_ℓ . We denote by n the number of particles in the cluster $\{-\ell, \dots, -\ell + \ell_1 - 1\}$ to the left of $B_k(x)$ in B_ℓ , which yields the decomposition

$$\frac{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma')}{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma)} = \frac{\sum_{n=1}^j a_n(\sigma')}{\sum_{n=1}^j a_n(\sigma)}. \quad (6.26)$$

where

$$a_n(\sigma) = a_n(\sigma, j, \ell_1, \ell_2) := \binom{n-1 + \sigma(-k)}{\ell_1 - n} \binom{j-n-p(\sigma)-1 + \sigma(k)}{\ell_2 - (j-n-p(\sigma))}. \quad (6.27)$$

The first binomial is the number of ergodic configuration on the left cluster with n particles and compatible with the left boundary of σ , whereas the second binomial coefficient is the number of ergodic configurations on the right cluster compatible with the right boundary of σ . The quantity $a_n(\sigma)$ is therefore the number of ergodic configurations on B_ℓ with j particles, such that the configuration in $B_k(x)$ is given by σ , and such that the number of particles in $\{-\ell, \dots, -\ell + \ell_1 - 1\}$ (resp. $\{\ell - \ell_2 + 1, \dots, \ell\}$) is n (resp. $j - p(\sigma) - n$).

Note that in both sums in the right hand side of (6.26), some terms vanish (using the convention $\binom{a}{b} = 0$ if $b \notin \{0, \dots, a\}$) therefore we do not need to be too specific with the index limits. Let us shorten

$$S(\sigma) = \sum_{n=1}^j a_n(\sigma), \quad \text{and} \quad S^* = \sum_{n=1}^j a_n^*, \quad (6.28)$$

where

$$a_n^* = \binom{n}{\ell_1 - n} \binom{j-n}{\ell_2 - (j-n)}.$$

Given these notations, we are going to prove that Lemma 6.12 is a consequence of Lemma 6.13 below.

Lemma 6.13. *For any configuration $\sigma \in \{0, 1\}^{B_k}$, and any positive δ_1, δ_2 , there exists a constant $C_4 := C_4(k, \delta_1, \delta_2)$ such that for all $j \in E_\ell(\delta_1)$*

$$\max_{x \in B_{(1-\delta_2)\ell-k}} |S(\sigma) - F(\rho_\ell, \sigma)S^*| \leq \frac{C_4}{\ell^{1/4}} S^*$$

where F is the function

$$F(\rho, \sigma) := \frac{(2\rho - 1)^{2p(\sigma)+2-\sigma(-k)-\sigma(k)}}{\rho^{p(\sigma)+2-\sigma(-k)-\sigma(k)}(1-\rho)^{p(\sigma)}} = \frac{(2\rho - 1)^2}{\rho^2} \alpha(\rho)^{p(\sigma)} \gamma(\rho)^{\sigma(-k)+\sigma(k)}. \quad (6.29)$$

Before proving this lemma, we show that it is enough to conclude the proof of Lemma 6.12. For that purpose, write

$$\frac{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma')}{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma)} = \frac{S(\sigma')}{S(\sigma)} = \frac{F(\rho_\ell, \sigma')}{F(\rho_\ell, \sigma)} + \frac{F(\rho_\ell, \sigma)S^*S(\sigma') - F(\rho_\ell, \sigma')S^*S(\sigma)}{F(\rho_\ell, \sigma)S^*S(\sigma)},$$

so that, thanks to Lemma 6.13,

$$\left| \frac{S(\sigma')}{S(\sigma)} - \frac{F(\rho_\ell, \sigma')}{F(\rho_\ell, \sigma)} \right| \leq \frac{C_5}{\ell^{1/4}} \frac{S(\sigma') + S(\sigma)}{F(\rho_\ell, \sigma)S(\sigma)},$$

where we defined $C_5 = C_5(k, \delta_1, \delta_2) = 2C_4(k, \delta_1, \delta_2)$.

By definition of F and equation (6.25),

$$\frac{\pi_\rho(\xi_{|B_k} = \sigma')}{\pi_\rho(\xi_{|B_k} = \sigma)} = \frac{F(\rho, \sigma')}{F(\rho, \sigma)},$$

and since $j \in E_\ell(\delta_1)$, $F(\rho_\ell, \sigma)$ is bounded away from 0 by some constant $c_F := c_F(k, \delta_1) > 0$. The bound above can therefore be rewritten for any $j \in E_\ell(\delta_1)$

$$\begin{aligned} & \left| \frac{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma')}{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma)} - \frac{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma')}{\pi_{\rho_\ell}(\xi_{|B_k} = \sigma)} \right| \\ & \leq \frac{C_5}{c_F \ell^{1/4}} \left(1 + \frac{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma')}{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma)} \right) \leq \frac{C_3}{\ell^{1/4} \widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{|B_k(x)} = \sigma)}, \end{aligned}$$

where $C_3 := C_3(k, \delta_1, \delta_2) = 2C_5/c_F$, which proves Lemma 6.12. \square

We now prove Lemma 6.13.

Proof of Lemma 6.13. Let us define

$$\begin{aligned} n^\sigma &= n - 1 + \sigma(-k), & j^\sigma &= j - p(\sigma) - 2 + \sigma(k) + \sigma(-k), \\ \ell_1^\sigma &= \ell_1 - 1 + \sigma(-k), & \ell_2^\sigma &= \ell_2 - 1 + \sigma(k), \end{aligned}$$

in order to rewrite

$$a_n(\sigma) = \binom{n^\sigma}{\ell_1^\sigma - n^\sigma} \binom{j^\sigma - n^\sigma}{\ell_2^\sigma - (j^\sigma - n^\sigma)}.$$

Note that the difference between n^σ (resp. j^σ , ℓ_1^σ , ℓ_2^σ) and n (resp. j , ℓ_1 , ℓ_2) is of order k . Define

$$n_*^\sigma = \frac{j^\sigma \ell_1^\sigma}{\ell_1^\sigma + \ell_2^\sigma},$$

which satisfies

$$|n_*^\sigma - \ell_1 \rho_\ell| \leq C(k, \delta_2). \quad (6.30)$$

Denote by $I_\ell(\sigma)$ the set of n 's in $\{0, \dots, j\}$ such that $|n^\sigma - n_*^\sigma| \leq \ell^{3/4}$. We can now write by triangular inequality

$$|S(\sigma) - F(\rho_\ell, \sigma)S^*| \leq \sum_{n \in I_\ell(\sigma)} \left| \frac{a_n(\sigma)}{a_n^*} - F(\rho_\ell, \sigma) \right| a_n^* + \sum_{n \notin I_\ell(\sigma)} (a_n(\sigma) + F(\rho_\ell, \sigma)a_n^*). \quad (6.31)$$

Before carrying on, let us state the following lemma, that we will prove after concluding the proof of Lemma 6.13.

Lemma 6.14. *There exists a constant $C_7 := C_7(k, \delta_1, \delta_2)$ such that*

$$\max_{x \in B_{(1-\delta_2)\ell-k}} \max_{n \in I_\ell(\sigma)} \left| \frac{a_n(\sigma)}{a_n^*} - F(\rho_\ell, \sigma) \right| \leq \frac{C_7}{\ell^{1/4}}.$$

Moreover, letting $n_0 = \lfloor j\ell_1/(\ell_1 + \ell_2) \rfloor$, there also exist two constants $C_8 := C_8(k, \delta_1, \delta_2)$ and $C_9 := C_9(k, \delta_2)$ such that

$$\max_{x \in B_{(1-\delta_2)\ell-k}} \max_{n \notin I_\ell(\sigma)} \frac{\max\{a_n(\sigma), a_n^*\}}{a_{n_0}^*} \leq C_8 e^{-C_9 \sqrt{\ell}}. \quad (6.32)$$

According to Lemma 6.14 above, (6.31) becomes

$$\begin{aligned} |S(\sigma) - F(\rho_\ell, \sigma)S^*| &\leq \sum_{n \in I_\ell(\sigma)} \frac{C_7}{\ell^{1/4}} a_n^* + |I_\ell(\sigma)^c| (1 + F(\rho_\ell, \sigma)) C_8 e^{-C_9 \sqrt{\ell}} a_{n_0}^* \\ &\leq \left(\frac{C_7}{\ell^{1/4}} + (2\ell + 1) C_8' e^{-C_9 \sqrt{\ell}} \right) S^*, \end{aligned}$$

where we used that $a_{n_0}^* \leq S^*$ and that $(1 + F(\rho_\ell, \sigma)) C_8 \leq C_8'(k, \delta_1, \delta_2)$, which holds since by assumption $\rho_\ell \in [(1 + \delta_1)/2, 1 - \delta_1]$. We therefore obtain as wanted that for some constant $C_4 = C_4(k, \delta_1, \delta_2)$

$$|S(\sigma) - F(\rho_\ell, \sigma)S^*| \leq \frac{C_4 S^*}{\ell^{1/4}},$$

which proves Lemma 6.13. \square

We now prove Lemma 6.14.

Proof of Lemma 6.14. We start by proving the first identity. Shortening $\hat{n} = j - n$, and extracting the excess terms yield that

$$\frac{a_n(\sigma)}{a_n^*} = \frac{(2n - \ell_1)^{1 - \sigma(-k)}}{n^{1 - \sigma(-k)}} \frac{(2\hat{n} - \ell_2) \cdots (2\hat{n} - 2p + \sigma(k) - \ell_2)}{\hat{n} \cdots (\hat{n} - p + \sigma(k)) (\ell_2 + p - \hat{n}) \cdots (\ell_2 + 1 - \hat{n})}. \quad (6.33)$$

Denote by $x_n = n/\ell_1$ the density in the left cluster, and $g(x_n) = (j - n)/\ell_2 = (j - \ell_1 x_n)/\ell_2$, which is close to the density in the right cluster. We start by proving that there exists a constant $C_{10}(k, \delta_1, \delta_2)$ such that

$$\max_{n \in I_\ell(\sigma)} \left\{ |x_n - \rho_\ell| + |g(x_n) - \rho_\ell| \right\} \leq \frac{C_{10}}{\ell^{1/4}}. \quad (6.34)$$

To prove this bound, one merely has to split by the triangular inequality

$$|x_n - \rho_\ell| \leq \frac{|n - n^\sigma|}{\ell_1} + \frac{|n^\sigma - n_*^\sigma|}{\ell_1} + \left| \frac{n_*^\sigma}{\ell_1} - \rho_\ell \right|.$$

Since $n^\sigma = n - 1 + \sigma(-k)$, and thanks to (6.30) the first and third terms in the right hand side above are less than $C(k, \delta_2)/\ell$ because $\ell_1 \geq \delta_2 \ell$. The second term is also less than $\ell^{-1/4}/\delta_2$ by definition of $I_\ell(\sigma)$. The estimation of $|g(x_n) - \rho_\ell|$ is obtained analogously, thus proving (6.34).

Since $\ell_1, \ell_2 \geq \delta_2 \ell$, (6.33) and (6.34) yield after elementary computations that

$$\max_{n \in I_\ell(\sigma)} \left| \frac{a_n(\sigma)}{a_n^*} - F(\rho_\ell, \sigma) \right| \leq \frac{C_7(k, \delta_1, \delta_2)}{\ell^{1/4}},$$

where F is the function defined in (6.29), as wanted.

Before proving the second identity (6.32), we state that a_n is relatively small as soon as the density n/ℓ_1 in the first cluster is far away from the expected density ρ_ℓ .

Lemma 6.15. *For any $\delta_2 > 0$, there exists a constant $K_1 := K_1(\delta_2)$ such that for any $\ell \geq 1$, $\delta_2 \ell \leq \ell_1, \ell_2 \leq 2\ell + 1$, any $\ell + 1 \leq j \leq 2\ell + 1$, and any integer n such that*

$$\frac{\ell_1}{2} < n \leq \ell_1, \quad \frac{\ell_2}{2} < j - n \leq \ell_2,$$

we have

$$\binom{n}{\ell_1 - n} \binom{j - n}{\ell_2 - (j - n)} \leq K_* \exp\left(-K_1 \frac{(n - n_*)^2}{\ell}\right), \quad (6.35)$$

where $n_* = j\ell_1/(\ell_1 + \ell_2)$, and K_* is a large constant depending only on j , ℓ_1 , and ℓ_2 . Note that n_* is not a priori an integer. Furthermore, there exists a constant $K_2 := K_2(\delta_1, \delta_2)$ such that, letting $n_0 = \lfloor n_* \rfloor$

$$\binom{n_0}{\ell_1 - n_0} \binom{j - n_0}{\ell_2 - (j - n_0)} \geq \frac{K_*}{2^8 \ell^2} \exp\left(-\frac{K_2}{\ell}\right). \quad (6.36)$$

Before proving this lemma, we conclude the proof of (6.32) and Lemma 6.14. Recall that we introduced $n_0 = \lfloor j\ell_1/(\ell_1 + \ell_2) \rfloor$ the approximate index of the largest term in S^* . We are going to prove that for any $n \notin I_\ell(\sigma)$, we have

$$\frac{\max\{a_n(\sigma), a_n^*\}}{a_{n_0}^*} \leq C_8 \exp(-C_9 \sqrt{\ell}).$$

This bound is a direct consequence of Lemma 6.15, since according to (6.35), we have ⁷

$$a_n(\sigma) \leq K_*(\sigma) \exp\left(-K_1(\delta_2/2) \frac{(n^\sigma - n_*^\sigma)^2}{\ell}\right)$$

and

$$a_n^* \leq K_* \exp\left(-K_1(\delta_2) \frac{(n - n_*)^2}{\ell}\right),$$

and (6.36) yields

$$a_{n_0}^* \geq \frac{K_*}{2^8 \ell^2} \exp\left(-\frac{K_2(\delta_1, \delta_2)}{\ell}\right).$$

According to the proof of Lemma 6.15 the constants $K_*(\sigma)$ and K_* are respectively given by

$$K_*(\sigma) = \exp\left(-(\ell_1^\sigma + \ell_2^\sigma)H(x_*^\sigma)\right) \quad \text{and} \quad K_* = \exp\left(-(\ell_1 + \ell_2)H(x_*)\right)$$

with $x_*^\sigma = j^\sigma/(\ell_1^\sigma + \ell_2^\sigma)$ and $x_* = j/(\ell_1 + \ell_2)$, and the entropy functional H is given below by (6.37). In particular, $x_*^\sigma - x_* = \mathcal{O}_\ell(1/\ell)$, and both are bounded away from $\frac{1}{2}$ and 1 by some constant depending only on k and δ_1 , therefore $|H(x_*^\sigma) - H(x_*)| \leq C(k, \delta_1)/\ell$, and since $|\ell_1^\sigma + \ell_2^\sigma - (\ell_1 + \ell_2)| \leq 2$, we finally obtain $K_*(\sigma)/K_* \leq C(k, \delta_1)$. It is easily checked that for any $n \notin I_\ell(\sigma)$, we have $|n - n_*| \wedge |n^\sigma - n_*^\sigma| \geq c\ell^{3/4}$ for some constant $c(k) > 0$ depending only on k . We finally obtain

$$\begin{aligned} & \frac{\max\{a_n(\sigma), a_n^*\}}{a_{n_0}^*} \\ & \leq 2^8 \ell^2 \max\{1, C(k, \delta_1)\} \exp\left(\frac{K_2(\delta_1, \delta_2)}{\ell} - \min\{K_1(\delta_2), K_1(\delta_2/2)\} c^2 \sqrt{\ell}\right), \end{aligned}$$

which is what we wanted to prove. \square

We now prove the technical concentration lemma we just used.

Proof of Lemma 6.15. Sondow's bound [17], which is a quantitative extension of Stirling's formula, can be written for any integer $0 \leq k \leq \ell$ as

$$\frac{1}{4 \max(k, \ell - k)} \frac{\ell^\ell}{k^k (\ell - k)^{\ell - k}} \leq \binom{\ell}{k} \leq \frac{\ell^\ell}{k^k (\ell - k)^{\ell - k}},$$

which immediately yields

$$\frac{1}{4\ell} \frac{\ell^\ell}{k^k (\ell - k)^{\ell - k}} \leq \binom{\ell}{k} \leq \frac{\ell^\ell}{k^k (\ell - k)^{\ell - k}},$$

⁷The dependency of $K_1(\delta_2/2)$ comes from the crude bound $\ell_1^\sigma, \ell_2^\sigma \geq \delta_2 \ell/2$.

Using both parts of this bound yields

$$\begin{aligned} & \frac{1}{16j^2} \frac{n^n}{(\ell_1 - n)^{\ell_1 - n} (2n - \ell_1)^{2n - \ell_1}} \frac{(j - n)^{j - n}}{(\ell_2 - (j - n))^{\ell_2 - (j - n)} (2(j - n) - \ell_2)^{2(j - n) - \ell_2}} \\ & \leq \binom{n}{\ell_1 - n} \binom{j - n}{\ell_2 - (j - n)} \leq \\ & \frac{n^n}{(\ell_1 - n)^{\ell_1 - n} (2n - \ell_1)^{2n - \ell_1}} \frac{(j - n)^{j - n}}{(\ell_2 - (j - n))^{\ell_2 - (j - n)} (2(j - n) - \ell_2)^{2(j - n) - \ell_2}}. \end{aligned}$$

Denoting $x_n = n/\ell_1$, and $g(x_n) = (j - n)/\ell_2 = (j - \ell_1 x_n)/\ell_2$, and introducing the entropy functional

$$H(x) = (1 - x) \log\left(\frac{1 - x}{x}\right) + (2x - 1) \log\left(\frac{2x - 1}{x}\right), \quad (6.37)$$

the previous bound rewrites

$$\begin{aligned} & \frac{1}{16j^2} \exp(-\ell_1 H(x_n) - \ell_2 H(g(x_n))) \\ & \leq \binom{n}{\ell_1 - n} \binom{j - n}{\ell_2 - (j - n)} \leq \\ & \exp(-\ell_1 H(x_n) - \ell_2 H(g(x_n))). \quad (6.38) \end{aligned}$$

Let us shorten $\Phi = \ell_1 H + \ell_2 H \circ g$. We now take a look at the quantity inside the exponential. Since $g' = -\ell_1/\ell_2$,

$$\Phi'(u) = \ell_1 (H'(u) - H'(g(u))) = \ell_1 \int_{g(u)}^u dv H''(v).$$

Denote by

$$x_* = \frac{j}{\ell_1 + \ell_2} = \frac{n_*}{\ell_1}$$

the only fixed point for g , so that

$$\Phi(x) = \Phi(x_*) + \ell_1 \int_{x_*}^x du \int_{g(u)}^u dv H''(v),$$

where

$$H''(u) = \frac{4}{2u - 1} - \frac{1}{u} + \frac{1}{1 - u}$$

is positive and bounded from below on $(\frac{1}{2}, 1)$ by some fixed constant c_0 , and uniformly bounded from above on any segment $[a, b] \subset (\frac{1}{2}, 1)$. Define

$$C_H(x) = \max\{H''(u), u \in [x, x_*] \text{ or } [x_*, x]\}.$$

Since $u - g(u) = (\ell_1 + \ell_2)(u - x_*)/\ell_2$, it is then a matter of simple verification to check that

$$\Phi(x_*) + c_0 \frac{\ell_1(\ell_1 + \ell_2)(x - x_*)^2}{2\ell_2} \leq \Phi(x) \leq \Phi(x_*) + C_H(x) \frac{\ell_1(\ell_1 + \ell_2)(x - x_*)^2}{2\ell_2},$$

which can be rewritten

$$\Phi(x_*) + c_0 \frac{(n - n_*)^2}{2 \max(\ell_1, \ell_2)} \leq \Phi(x) \leq \Phi(x_*) + C_H(x) \frac{(n - n_*)^2}{\min(\ell_1, \ell_2)}.$$

We inject this bound in (6.38), to finally obtain that for any $[a, b] \subset (\frac{1}{2}, 1)$ and for any n ,

$$\begin{aligned} \frac{K_*}{16j^2} \exp\left(-C_H(n/\ell_1) \frac{(n-n_*)^2}{\min\{\ell_1, \ell_2\}}\right) \\ \leq \binom{n}{\ell_1-n} \binom{j-n}{\ell_2-(j-n)} \leq K_* \exp\left(-c_0 \frac{(n-n_*)^2}{2 \max\{\ell_1, \ell_2\}}\right), \end{aligned}$$

where we defined $K_* = \exp(-\Phi(x_*))$. Letting $n_0 = \lfloor n_* \rfloor$, and assuming that $\ell_1 \geq \delta_2 \ell$, one can verify that $C_H(n_0/\ell_1) \leq C(\delta_1, \delta_2)$, so that the left part of the last estimation rewrites

$$\binom{n_0}{\ell_1-n_0} \binom{j-n_0}{\ell_2-(j-n_0)} \geq \frac{K_*}{16j^2} \exp\left(-\frac{\delta_2 C(\delta_1, \delta_2)}{\ell}\right),$$

thus concluding the proof. \square

To conclude this section, we now give the proof of Lemma 6.11, which we postponed early in the section, and prove that the canonical measure $\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{B_k(x)} = \sigma)$ is of order δ_1 if j is larger than $(1-\delta_1)(2\ell+1)$.

Proof of Lemma 6.11. We will once again only prove (6.22) in the case where $\varepsilon_1 = \varepsilon_2 = 1$, the other cases being analogous. To do so, denote by $\mathbf{1}_k$ the full configuration on B_k , using the same notations as in (6.26), rewrite the left hand side of (6.22) as

$$\frac{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{B_k(x)} = \sigma)}{\widehat{\pi}_{\rho,1,1}^{\ell,j}(\mathfrak{S}_{B_k(x)} = \mathbf{1}_k)} = \frac{\sum_{n=1}^j a_n(\sigma)}{\sum_{n=1}^j a_n(\mathbf{1}_k)}. \quad (6.39)$$

We are going to show that assuming that the configuration σ is not full, there exists a constant $C(k, \delta_2)$ such that for any $1 \leq n \leq j$,

$$a_n(\sigma) \leq C(k, \delta_2) \delta_1 a_n(\mathbf{1}_k). \quad (6.40)$$

We first prove that if $\widehat{n} = j - n < (1 - 4\delta_1/\delta_2)\ell_2$ the two sides above vanish, because there can be no associated configuration. Indeed, by assumption,

$$j \geq (1 - \delta_1)(2\ell + 1).$$

Furthermore, we can safely assume that $n \leq \ell_1$ since if not, we are trying to put more particles than there are holes in the first box, and the binomial coefficients obviously vanish. Therefore,

$$\widehat{n} \geq (1 - \delta_1)(2\ell + 1) - \ell_1 = \ell_2 + (2k + 1) - \delta_1(2\ell + 1) \geq \ell_2 \left(1 - \frac{4\delta_1}{\delta_2}\right)$$

where we used that by assumption, $\ell_2 \geq \delta_2 \ell \geq \delta_2(2\ell + 1)/4$. In particular, for any n such that $\widehat{n}/\ell_2 < 1 - 4\delta_1/\delta_2$, (6.40) obviously holds since both sides vanish.

We now assume that $\widehat{n}/\ell_2 \geq 1 - 4\delta_1/\delta_2$. To prove (6.40) in this case, by extracting the excess terms, as we did in (6.33), we rewrite

$$\begin{aligned} a_n(\sigma) &= a_n(\mathbf{1}_k) \frac{(2n - \ell_1)^{1-\sigma(-k)}}{n^{1-\sigma(-k)}} \\ &\times \frac{(\widehat{n} - p(\sigma) + \sigma(k) - 1) \cdots (\widehat{n} - 2k)(\ell_2 + 2k + 1 - \widehat{n}) \cdots (\ell_2 + p(\sigma) + 1 - \widehat{n})}{(2\widehat{n} - 2p(\sigma) + \sigma(k) - \ell_2 - 1) \cdots (2\widehat{n} - 4k + 1 - \ell_2)}. \end{aligned}$$

The first part of the quotient is less than one since $(2n - \ell_1)/n \leq 1$. Divide each term in the second part of the right hand side by ℓ_2 , underestimate the powers in the numerator and use that $x_n \leq 1$ to obtain that

$$a_n(\sigma) \leq a_n(\mathbf{1}_k) \frac{x_n(1 - x_n + \varepsilon_\ell)}{(2x_n - 1 - \varepsilon_\ell)^{4k - 2p(\sigma)}},$$

where $\varepsilon_\ell = 4k/\ell_2$, and $x_n = \widehat{n}/\ell_2 \geq 1 - 4\delta_1/\delta_2$. The configuration σ contains at least one empty site, so that $2k + 1 \geq p + 1$. We now use that $x_n \geq 1 - C\delta_1$, therefore ε_ℓ is small w.r.t. x_n and $2x_n - 1$. Recall that ℓ goes to infinity before δ_1 goes to 0. We obtain as wanted

$$a_n(\sigma) \leq a_n(\mathbf{1}_k)(o_\ell(1) + C(k, \delta_2)\delta_1),$$

which concludes the proof. \square

7. PROOF OF THE REPLACEMENT LEMMA 5.4

7.1. Strategy of the proof. Recall that $\widetilde{\mu}_N$ has been defined in (5.1). Let $\widetilde{\mu}_s^N := \widetilde{\mu}_N e^{sN^2 \mathcal{L}_N}$ be the measure at macroscopic time s of the process started from $\widetilde{\mu}_N$. Fix $\rho \in (\frac{1}{2}, 1)$, and let us introduce, for any $s \geq 0$

$$f_s^N = \frac{d\widetilde{\mu}_s^N}{d\nu_{\rho,N}}, \quad \text{and} \quad \bar{f}_T^N = \frac{1}{T} \int_0^T f_s^N ds,$$

where $\nu_{\rho,N}$ is the translation invariant, periodic grand canonical measure defined in Definition 6.3. This density f_s^N is well defined: $\nu_{\rho,N}(\eta) = 0$ if $\eta \notin \mathcal{E}_N$, and $\widetilde{\mu}_N^s$ has support in \mathcal{E}_N by definition.

Denote \mathcal{D}_N the Dirichlet form, defined as

$$\begin{aligned} \mathcal{D}_N(g) &:= -\nu_{\rho,N}(\sqrt{g} \mathcal{L}_N \sqrt{g}) \\ &= \frac{1}{2} \nu_{\rho,N} \left(\sum_{x \in \mathbb{T}_N} \eta(x+2)\eta(x-1) (\sqrt{g}(\eta^{x,x+1}) - \sqrt{g}(\eta))^2 \right). \end{aligned}$$

It is straightforward to prove that \mathcal{D}_N thus defined is non-negative and convex. To obtain the second identity, we used that $\nu_{\rho,N}$ -a.s., any configuration such that $\eta(x) \neq \eta(x+1)$ and such that both η and $\eta^{x,x+1}$ are ergodic, we must have $\eta(x+2) = \eta(x-1) = 1$, because two empty sites cannot be neighbors under $\nu_{\rho,N}$. In Proposition 7.4 below we will prove that the Dirichlet form \mathcal{D}_N evaluated at \bar{f}_T^N satisfies:

$$\mathcal{D}_N(\bar{f}_T^N) \leq \frac{C}{N}. \quad (7.1)$$

Let us turn now to the proof of Lemma 5.4. We rewrite the expression under the limit in (5.9) as:

$$\nu_{\rho,N} \left(\bar{f}_T^N \times \frac{T}{N} \sum_{x \in \mathbb{T}_N} \tau_x |V_{\varepsilon N}(h, \cdot)| \right).$$

Therefore, in order to prove (5.9) it is enough to prove the following:

$$\limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \sup_{f: \mathcal{D}_N(f) \leq \frac{C}{N}} \nu_{\rho,N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x |V_{\varepsilon N}(h, \cdot)| \right) = 0, \quad (7.2)$$

where the supremum is taken over all densities f w.r.t. $\nu_{\rho,N}$ which satisfy $\mathcal{D}_N(f) \leq \frac{C}{N}$. Take $\ell \in \mathbb{N}$. Inside the absolute values in (7.2) we add and subtract

$$\frac{1}{2\varepsilon N + 1} \sum_{|y| \leq \varepsilon N} \left(\frac{1}{2\ell' + 1} \sum_{|z-y| \leq \ell'} \tau_z h(\eta) - \pi_{\rho^\ell(y)}(\eta)(h) \right),$$

where we shortened $\ell' = \ell - 1$, to make the sum measurable w.r.t. the sites in $\tau_y B_\ell$. Thanks to the triangular inequality, we are reduced to estimate three terms:

$$I_{\ell, \varepsilon N}^1 := \nu_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x \left| \frac{1}{2\varepsilon N + 1} \sum_{|y| \leq \varepsilon N} \left(\tau_y h - \frac{1}{2\ell' + 1} \sum_{|z-y| \leq \ell'} \tau_z h \right) \right| \right),$$

$$I_{\ell, \varepsilon N}^2 := \nu_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x \left| \frac{1}{2\ell' + 1} \sum_{|y| \leq \ell'} \tau_y h - \pi_{\rho^{\ell(0)}}(h) \right| \right), \quad (7.3)$$

$$I_{\ell, \varepsilon N}^3 := \nu_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x \left| \frac{1}{2\varepsilon N + 1} \sum_{|y| \leq \varepsilon N} \pi_{\rho^{\ell(y)}}(h) - \pi_{\rho^{\varepsilon N(0)}}(h) \right| \right). \quad (7.4)$$

We are going to show that each one of these terms vanishes, as first $N \rightarrow \infty$ then $\varepsilon \rightarrow 0$ and then $\ell \rightarrow \infty$. We treat them separately: the first one can be easily bounded as follows ; there exists $C > 0$ such that

$$I_{\ell, \varepsilon N}^1 \leq \frac{C\ell}{\varepsilon N} \nu_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x h \right) \leq \frac{C\ell}{\varepsilon N} \xrightarrow{N \rightarrow \infty} 0,$$

since h is bounded by 1. Let us slightly change our notation, and define

$$\tilde{V}_\ell(h, \eta) = \frac{1}{2\ell' + 1} \sum_{|y| \leq \ell'} \tau_y h(\eta) - \pi_{\rho^{\ell(0)(\eta)}}(h),$$

whose only difference with $V_\ell(h, \eta)$ is that the first average is taken on a slightly reduced box to make the sum measurable w.r.t. the sites in B_ℓ . Then, the second term (7.3) rewrites with this notation as

$$I_{\ell, \varepsilon N}^2 = \nu_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x |\tilde{V}_\ell(h, \cdot)| \right),$$

and it vanishes thanks to the *one-block estimate* stated and proved in Lemma 7.1 below. We finally bound the third one (7.4) as

$$I_{\ell, \varepsilon N}^3 \leq \sup_{|y| \leq \varepsilon N} \nu_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x \left| \pi_{\rho^{\varepsilon N(0)}}(h) - \pi_{\rho^{\ell(y)}}(h) \right| \right)$$

$$\leq 4 \sup_{|y| \leq \varepsilon N} \nu_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \left| \rho^{\varepsilon N}(x) - \rho^\ell(x+y) \right| \right), \quad (7.5)$$

the last inequality coming from the fact that $|\pi_{\rho_\star}(h) - \pi_\rho(h)| \leq 4|\rho_\star - \rho|$ for any $\rho, \rho_\star \in [0, 1]$ (after using the identity (6.7)). The left hand side of (7.5) vanishes according to the *two-blocks estimate* stated and proved in Lemma 7.3 below.

7.2. One-block estimate.

Lemma 7.1. *Fix $\rho \in (\frac{1}{2}, 1]$. For any finite constant $C > 0$,*

$$\limsup_{\ell \rightarrow \infty} \limsup_{N \rightarrow \infty} \sup_{f: \mathcal{D}_N(f) \leq \frac{C}{N}} \nu_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x |\tilde{V}_\ell(h, \cdot)| \right) = 0. \quad (7.6)$$

Proof. Following [11, Chapter 5, Section 4], we divide the proof into several steps listed below.

STEP 1 – REDUCTION TO MICROSCOPIC BLOCKS. Since the measure $\nu_{\rho,N}$ is translation invariant, we can rewrite the expectation that appears in (7.6) as

$$\nu_{\rho,N}\left(\tilde{f} \mid \tilde{V}_\ell(h, \cdot)\right), \quad \text{where } \tilde{f} := \frac{1}{N} \sum_{x \in \mathbb{T}_N} \tau_x f.$$

Note that $\tilde{V}_\ell(h, \eta)$ depends on η only through its values on the box $B_\ell = \{-\ell, \dots, \ell\} \subset \mathbb{T}_N$.

We now define the conditional expectation of \tilde{f} with respect to the coordinates on B_ℓ , namely:

$$\tilde{f}_\ell(\sigma) = \begin{cases} \frac{1}{\widehat{\nu}_{\rho,N}^\ell(\sigma)} \sum_{\substack{\eta \in \{0,1\}^{\mathbb{T}_N} \\ \eta|_{B_\ell} = \sigma}} \tilde{f}(\eta) \nu_{\rho,N}(\eta) & \text{for any } \sigma \in \widehat{\mathcal{E}}_{B_\ell}^\ell, \\ 0 & \text{else.} \end{cases}$$

Lemma 7.2. *The function \tilde{f}_ℓ is a density with respect to $\widehat{\nu}_{\rho,N}^\ell$, and*

$$\sup_{\substack{f \text{ density w.r.t. } \widehat{\nu}_{\rho,N}^\ell, \\ f|_{\widehat{\mathcal{E}}_{B_\ell}^c} = 0}} \|f\|_\infty \leq C(\ell, \rho). \quad (7.7)$$

Proof of Lemma 7.2. The first point is clear. For the second, notice that if f is a density w.r.t. $\widehat{\nu}_{\rho,N}^\ell$, then $\sum_{\sigma \in \widehat{\mathcal{E}}_{B_\ell}^\ell} \widehat{\nu}_{\rho,N}^\ell(\sigma) f(\sigma) = 1$, and therefore $f(\sigma) \leq \widehat{\nu}_{\rho,N}^\ell(\sigma)^{-1}$ for all $\sigma \in \widehat{\mathcal{E}}_{B_\ell}^\ell$.

By formula (6.14) and estimate (6.17), for any $\sigma \in \widehat{\mathcal{E}}_{B_\ell}^\ell$, we have

$$\widehat{\nu}_{\rho,N}^\ell(\sigma) \geq \kappa \min(\alpha^{2\ell+1}, 1) \beta^{2\ell+1} - \frac{C(\ell, \rho)}{\sqrt{N}} \geq C(\ell, \rho) > 0$$

for large enough N , and Lemma 7.2 is proved. \square

Going back to Lemma 7.1, we finally rewrite the expectation in (7.6) as

$$\widehat{\nu}_{\rho,N}^\ell\left(\tilde{f}_\ell \mid \tilde{V}_\ell(h, \cdot)\right). \quad (7.8)$$

As a consequence of Lemma 6.8 and Lemma 7.2, we can replace in the large N limit $\widehat{\nu}_{\rho,N}^\ell$ with $\widehat{\pi}_\rho^\ell$, and to prove (7.6) we are reduced to prove that

$$\limsup_{\ell \rightarrow \infty} \limsup_{N \rightarrow \infty} \sup_{f: \mathcal{D}_N(f) \leq \frac{C}{N}} \widehat{\pi}_\rho^\ell\left(\tilde{f}_\ell \mid \tilde{V}_\ell(h, \cdot)\right) = 0. \quad (7.9)$$

STEP 2 – ESTIMATES ON THE DIRICHLET FORM OF \tilde{f}_ℓ . Let us first introduce some notations: for any $y \in \mathbb{T}_N$, and positive function g , define

$$I_y(g, \eta) = \frac{1}{2} \eta(y+2) \eta(y-1) (\sqrt{g}(\eta^{y,y+1}) - \sqrt{g}(\eta))^2, \quad (7.10)$$

and consider two Dirichlet forms restricted to bonds in B_ℓ defined for all densities $f: \{0, 1\}^{B_\ell} \rightarrow \mathbb{R}_+$ w.r.t. $\widehat{\nu}_{\rho,N}^\ell$ which vanish on $\widehat{\mathcal{E}}_{B_\ell}^c$ by

$$\begin{aligned} \mathcal{D}_N^\ell(f) &:= \sum_{y: \{y,y+1\} \subset B_\ell} \widehat{\nu}_{\rho,N}^{\ell+1}(I_y(f, \cdot)) \\ \mathcal{D}^\ell(f) &:= \sum_{y: \{y,y+1\} \subset B_\ell} \widehat{\pi}_\rho^{\ell+1}(I_y(f, \cdot)). \end{aligned}$$

From Lemma 6.8 and Lemma 7.2, these two Dirichlet forms are pretty close as $N \rightarrow \infty$. Recalling the definition (7.1) we also have, for any density $f : \{0, 1\}^{\mathbb{T}^N} \rightarrow \mathbb{R}_+$,

$$\mathcal{D}_N(f) = \sum_{x \in \mathbb{T}^N} \nu_{\rho, N}(I_x(f, \cdot)).$$

The exact same proof as in [11, Step 3 p.86] shows that

$$\text{if } \mathcal{D}_N(f) \leq \frac{C}{N} \text{ then } \mathcal{D}_N^\ell(\tilde{f}_\ell) \leq \frac{C(\ell)}{N^2}, \quad (7.11)$$

where $C(\ell)$ is a positive constant. From Lemma 6.8, this implies

$$\mathcal{D}^\ell(\tilde{f}_\ell) \leq \frac{C'(\ell)}{\sqrt{N}},$$

where $C'(\ell)$ is another positive constant. For the sake of completeness we expose quickly the arguments which prove (7.11):

- first, by the property of the marginal distribution, we have:

$$\widehat{\nu}_{\rho, N}^{\ell+1}(I_y(\tilde{f}_\ell, \cdot)) = \nu_{\rho, N}(I_y(\tilde{f}_\ell, \cdot)),$$

for any y such that $\{y, y+1\} \subset B_\ell$, which implies

$$\mathcal{D}_N^\ell(\tilde{f}_\ell) = \sum_{y: \{y, y+1\} \subset B_\ell} \nu_{\rho, N}(I_y(\tilde{f}_\ell, \cdot));$$

- second, since \tilde{f} and $\nu_{\rho, N}$ are translation invariant,

$$\sum_{y: \{y, y+1\} \subset B_\ell} \nu_{\rho, N}(I_y(\tilde{f}_\ell, \cdot)) = \frac{2\ell}{N} \mathcal{D}_N(\tilde{f}_\ell);$$

- finally, by convexity of the Dirichlet form, $\mathcal{D}_N(\tilde{f}_\ell) \leq \mathcal{D}_N(f)$.

Putting all these statements together proves the claim (7.11). As a result, to prove (7.9) we are reduced to prove that

$$\limsup_{\ell \rightarrow \infty} \limsup_{N \rightarrow \infty} \sup_{f \in \mathcal{A}_{\ell, N}} \widehat{\pi}_\rho^\ell(f | \widetilde{V}_\ell(h, \cdot)) = 0, \quad (7.12)$$

where we denoted

$$\mathcal{A}_{\ell, N} = \left\{ f : \widehat{\mathcal{E}}_{B_\ell} \rightarrow \mathbb{R}_+ : \widehat{\nu}_{\rho, N}^\ell(f) = 1, \quad \mathcal{D}^\ell(f) \leq \frac{C'(\ell)}{\sqrt{N}} \right\}.$$

Note that for any $f \in \mathcal{A}_{\ell, N}$, we have $\|f\|_\infty \leq C(\ell, \rho)$, where

$$C(\ell, \rho)^{-1} = \liminf_{N \rightarrow \infty} \inf_{\sigma \in \widehat{\mathcal{E}}_{B_\ell}} \widehat{\nu}_{\rho, N}^\ell(\sigma) > 0.$$

STEP 3 – LIMIT AS $N \rightarrow \infty$. This step is exactly identical to [11, Chapter 5, Section 4, Step 4], and mainly relies on the lower semicontinuity of the Dirichlet form and the continuity of $\widehat{\pi}_\rho^\ell(f)$ with respect to f . Letting

$$\mathcal{A}_\ell^* = \left\{ f : \widehat{\mathcal{E}}_{B_\ell} \rightarrow \mathbb{R}_+ : \widehat{\pi}_\rho^\ell(f) = 1, \quad \mathcal{D}^\ell(f) = 0 \right\},$$

thanks to Lemma 6.8, this reduces the proof of (7.12) to the proof of

$$\limsup_{\ell \rightarrow \infty} \sup_{f \in \mathcal{A}_\ell^*} \widehat{\pi}_\rho^\ell(f | \widetilde{V}_\ell(h, \cdot)) = 0. \quad (7.13)$$

STEP 4 – DECOMPOSITION ALONG HYPERPLANES WITH A FIXED NUMBER OF PARTICLES. Note that a density which has a vanishing Dirichlet form is constant on each hyperplane with a fixed total number of particles (recall Lemma 2.1). Therefore, the proof of (7.13) reduces to the proof of

$$\limsup_{\ell \rightarrow \infty} \sup_{j \in \{\ell, \dots, 2\ell+1\}} \widehat{\pi}_\rho^{\ell, j} \left(|\widetilde{V}_\ell(h, \cdot)| \right) = 0. \quad (7.14)$$

STEP 5 – AN APPLICATION OF THE EQUIVALENCE OF ENSEMBLES. Recall that we defined $\ell' = \ell - 1$, and note that

$$\widehat{\pi}_\rho^{\ell, j} \left(|\widetilde{V}_\ell(h, \cdot)| \right) = \widehat{\pi}_\rho^{\ell, j} \left(\left| \frac{1}{2\ell'+1} \sum_{x \in B_{\ell'}} \tau_x h - \pi_{j/(2\ell'+1)}(h) \right| \right). \quad (7.15)$$

Let us fix $k \in \mathbb{N}$ and decompose $B_{\ell'}$ into boxes of length $2k+1$, as follows: first, fix some $\delta > 0$, and define $q = \lfloor \ell'(1-\delta)/(2k+1) \rfloor$. We partition

$$B_{\ell'} = \mathcal{A}_\ell \sqcup \bigsqcup_{i=-q}^q \mathcal{B}_k(i), \quad \text{where } \mathcal{B}_k(i) := \{(2k+1)i - k, \dots, (2k+1)i + k\},$$

so that $\mathcal{B}_k(i)$ is the box of size $2k+1$ centered in the site $(2k+1)i$ and \mathcal{A}_ℓ its complementary set in $B_{\ell'}$, namely

$$\mathcal{A}_\ell = B_{\ell'} \setminus \bigsqcup_{i=-q}^q \mathcal{B}_k(i) = \{-\ell', \dots, -(2k+1)q - k - 1\} \cup \{(2k+1)q + k + 1, \dots, \ell'\}.$$

By construction, the size of \mathcal{A}_ℓ is

$$|\mathcal{A}_\ell| = 2(\ell' - (2k+1)q - k) \leq 2(\delta\ell' + k + 1) \leq 3\delta\ell'$$

for ℓ large enough. Then, in (7.15), the total contribution of the sites $x \in \mathcal{A}_\ell$ is less than $3\|h\|_\infty\delta$, so that (7.15) is bounded from above by

$$3\delta\|h\|_\infty + \frac{(2k+1)}{(2\ell'+1)} \sum_{i=-q}^q \widehat{\pi}_\rho^{\ell, j} \left(\left| \frac{1}{2k+1} \sum_{x \in \mathcal{B}_k(i)} \tau_x h - \pi_{j/(2\ell'+1)}(h) \right| \right). \quad (7.16)$$

Recall that δ and k are fixed, and that for $i \in \{-q, \dots, q\}$, we have $(2k+1)i \in B_{(1-\delta)\ell}$. We can therefore apply Corollary 6.10 to the translations $\tau_{(2k+1)i}f$, of the local function

$$f = \left| \frac{1}{2k+1} \sum_{x \in B_k} \tau_x h - \pi_{j/(2\ell'+1)}(h) \right|,$$

for any fixed δ, k : uniformly in $i \in \{-q, \dots, q\}$, the expectation converges as $\ell \rightarrow \infty$ towards

$$\pi_{j/(2\ell'+1)} \left(\left| \frac{1}{2k+1} \sum_{x \in \mathcal{B}_k(i)} \tau_x h - \pi_{j/(2\ell'+1)}(h) \right| \right).$$

Therefore, (7.16) can be rewritten

$$\begin{aligned} o_\ell(1) + 3\delta\|h\|_\infty + \frac{(2k+1)}{(2\ell'+1)} \sum_{i=-q}^q \pi_{j/(2\ell'+1)} \left(\left| \frac{1}{2k+1} \sum_{x \in \mathcal{B}_k(i)} \tau_x h(\sigma) - \pi_{j/(2\ell'+1)}(h) \right| \right) \\ \leq_{\ell \rightarrow \infty} 3\delta\|h\|_\infty + \sup_{\alpha \in [\frac{1}{2}, 1]} \pi_\alpha \left(\left| \frac{1}{2k+1} \sum_{x \in B_k} \tau_x h - \pi_\alpha(h) \right| \right), \end{aligned} \quad (7.17)$$

because π_α is translation invariant and $(2k+1)(2q+1)/(2\ell'+1) \leq 1$ for ℓ large enough. We now need to eliminate the low densities for which we do not have exponential decay of the correlations for π_α .

Fix $\alpha_0 > \frac{1}{2}$, in the limit $\ell \rightarrow \infty$ then $\delta \rightarrow 0$ we can therefore bound (7.16) by

$$\sup_{\alpha \in [\frac{1}{2}, \alpha_0]} \pi_\alpha \left(\left| \frac{1}{2k+1} \sum_{x \in B_k} \tau_x h - \pi_\alpha(h) \right| \right) + \sup_{\alpha \in (\alpha_0, 1]} \pi_\alpha \left(\left| \frac{1}{2k+1} \sum_{x \in B_k} \tau_x h - \pi_\alpha(h) \right| \right), \quad (7.18)$$

Since h is a non-negative function, we can crudely bound the first member of (7.18) from the triangular inequality by

$$\sup_{\alpha \in [\frac{1}{2}, \alpha_0]} 2\pi_\alpha(h) = \sup_{\alpha \in [\frac{1}{2}, \alpha_0]} 2 \frac{2\alpha - 1}{\alpha} \leq 4(2\alpha_0 - 1),$$

which vanishes as $\alpha_0 \rightarrow \frac{1}{2}$ uniformly in k . To estimate the second member of (7.18), we apply the law of large numbers for the grand canonical measure π_α . The variables $\{\sigma(x) : x \in \mathbb{Z}\}$ are not independent. Therefore, we need here to use the exponential decay of correlations proved in Corollary 6.6, which is uniform on $(\alpha_0, 1]$ for any fixed $\alpha_0 > \frac{1}{2}$. The strong law of large numbers has been proven in this case for instance in [13]. As a result, for any fixed $\frac{1}{2} < \alpha_0 \leq 1$, the second line of (7.18) vanishes as $k \rightarrow \infty$. We then let $\alpha_0 \rightarrow \frac{1}{2}$ to end the proof of Lemma 7.1. \square

7.3. Two-blocks estimate.

Lemma 7.3. Fix $\rho \in (\frac{1}{2}, 1]$. For every finite constant $C > 0$,

$$\limsup_{\ell \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \sup_{f : \mathcal{D}_N(f) \leq \frac{C}{N}} \sup_{|y| \leq \varepsilon N} v_{\rho, N} \left(f \times \frac{1}{N} \sum_{x \in \mathbb{T}_N} \left| \rho^{\varepsilon N}(x) - \rho^\ell(x+y) \right| \right) = 0. \quad (7.19)$$

Proof of Lemma 7.3. Most of the difficulties to prove the two-blocks estimate have been treated in the one-block estimate, so that we will merely sketch the proof. As in the one-block estimate, define $\tilde{f} = N^{-1} \sum \tau_x f$. To prove Lemma 7.3, it is enough to prove that

$$\limsup_{\ell \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \sup_{f : \mathcal{D}_N(f) \leq \frac{C}{N}} \sup_{|y| \leq \varepsilon N} v_{\rho, N} \left(\tilde{f} \left| \rho^{\varepsilon N}(0) - \rho^\ell(y) \right| \right) = 0.$$

Similarly to [11, Section 5.5], by triangular inequality, to prove Lemma 7.3, it is enough to prove that

$$\limsup_{\ell \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \sup_{f : \mathcal{D}_N(f) \leq \frac{C}{N}} \sup_{\sqrt{N} \leq |y| \leq 2\varepsilon N} v_{\rho, N} \left(\tilde{f} \left| \rho^\ell(0) - \rho^\ell(y) \right| \right) = 0. \quad (7.20)$$

Fix $|y| \geq \sqrt{N}$ and define

$$\mathcal{D}_{N, y, 1}^\ell(h) := \sum_{x : \{x, x+1\} \subset B_\ell \cup \tau_y B_\ell} v_{\rho, N}(I_x(h, \cdot))$$

$$\mathcal{D}_{N, y, 2}(h) := v_{\rho, N}(J_{0, y}(h, \cdot))$$

where I_x was defined in (7.10) and

$$J_{0,y}(g, \eta) := (\eta(-1)\eta(0)\eta(1) + \eta(y-1)\eta(y)\eta(y+1)) (\sqrt{g}(\eta^{0,y}) - \sqrt{g}(\eta))^2.$$

This last element allows jumps from 0 to y and vice-versa, while ensuring that one never leaves the ergodic component. We now condition the density f to $B_\ell \cup \tau_y B_\ell$, by letting, for any pair of configurations $(\sigma_1, \sigma_2) \in \widehat{\mathcal{E}}_{B_\ell} \times \widehat{\mathcal{E}}_{\tau_y B_\ell}$

$$\widetilde{f}_\ell^y(\sigma_1, \sigma_2) := \frac{1}{\nu_{\rho,N}(\sigma_1, \sigma_2)} \sum_{\substack{\eta \in \{0,1\}^{\mathbb{Z}^N} \\ \eta|_{B_\ell} = \sigma_1 \text{ and } \eta|_{\tau_y B_\ell} = \sigma_2}} \widetilde{f}(\eta) \nu_{\rho,N}(\eta).$$

Following the same steps as in [11, Section 5.5], we obtain that, since $\mathcal{D}_N(f) \leq C/N$,

$$\max_{\sqrt{N} \leq |y| \leq 2\epsilon N} \mathcal{D}_{N,y,1}^\ell(\widetilde{f}_\ell^y) + \mathcal{D}_{N,y,2}^\ell(\widetilde{f}_\ell^y) \leq C(\ell) o_\epsilon(1). \quad (7.21)$$

For any pair of configurations $(\sigma_1, \sigma_2) \in \widehat{\mathcal{E}}_{B_\ell} \times \widehat{\mathcal{E}}_{\tau_y B_\ell}$, denote $j(\sigma_1, \sigma_2)$ their total number of particles

$$j(\sigma_1, \sigma_2) = \sum_{x \in B_\ell} \sigma_1(x) + \sigma_2(y+x).$$

We claim that for any \widetilde{f}_ℓ^y satisfying (7.21), we can write for $(\sigma_1, \sigma_1', \sigma_2, \sigma_2') \in \widehat{\mathcal{E}}_{B_\ell}^2 \times \widehat{\mathcal{E}}_{\tau_y B_\ell}^2$ such that $j(\sigma_1, \sigma_2) = j(\sigma_1', \sigma_2')$,

$$|\widetilde{f}_\ell^y(\sigma_1, \sigma_2) - \widetilde{f}_\ell^y(\sigma_1', \sigma_2')| \leq C'(\ell) o_\epsilon(1), \quad (7.22)$$

for some large constant $C'(\ell)$ which does not depend neither on $\sqrt{N} \leq |y| \leq 2\epsilon N$ nor on the chosen configurations $(\sigma_1, \sigma_1', \sigma_2, \sigma_2')$. This bound is proved straightforwardly by the following three steps:

- Choose a sequence of licit jumps either in B_ℓ , in $\tau_y B_\ell$, or between 0 and y , allowing to reach (σ_1', σ_2') from (σ_1, σ_2) . Uniformly in (σ_1, σ_2) and (σ_1', σ_2') , the number $n := n(\sigma_1, \sigma_2, \sigma_1', \sigma_2')$ of jumps necessary for the construction of this path can be crudely bounded from above by some constant K_ℓ independent of the configurations and of y . Denote $(\sigma_1, \sigma_2) = (\sigma_1^0, \sigma_2^0), (\sigma_1^1, \sigma_2^1), \dots, (\sigma_1^n, \sigma_2^n) = (\sigma_1', \sigma_2')$ such a path.
- Then,

$$|\widetilde{f}_\ell^y(\sigma_1, \sigma_2) - \widetilde{f}_\ell^y(\sigma_1', \sigma_2')| \leq \sum_{k=0}^{n-1} |\widetilde{f}_\ell^y(\sigma_1^{k+1}, \sigma_2^{k+1}) - \widetilde{f}_\ell^y(\sigma_1^k, \sigma_2^k)|. \quad (7.23)$$

- Since all the jumps from (σ_1^k, σ_2^k) to $(\sigma_1^{k+1}, \sigma_2^{k+1})$ are by assumption allowed for our non-ergodic dynamics, each of the terms in the sum in the right hand side above can be crudely bounded from above by

$$\begin{aligned} & 2 \left[\|\widetilde{f}_\ell^y\|_\infty \left(\sqrt{\widetilde{f}_\ell^y(\sigma_1^{k+1}, \sigma_2^{k+1})} - \sqrt{\widetilde{f}_\ell^y(\sigma_1^k, \sigma_2^k)} \right)^2 \right]^{1/2} \\ & \leq 2 \left[\frac{\|\widetilde{f}_\ell^y\|_\infty \left(\mathcal{D}_{N,y,1}^\ell(\widetilde{f}_\ell^y) + \mathcal{D}_{N,y,2}^\ell(\widetilde{f}_\ell^y) \right)}{\nu_{\rho,N}(\eta|_{B_\ell} = \sigma_1^k, \eta|_{\tau_y B_\ell} = \sigma_2^k)} \right]^{1/2} \leq \widetilde{C}(\ell) o_\epsilon(1), \quad (7.24) \end{aligned}$$

where the last bound comes from a straightforward adaptation of Lemma 7.2.

Together, the two bounds (7.23) and (7.24) prove (7.22), by letting $C'_\ell = K_\ell \tilde{C}_\ell$.

We can now project $\nu_{\rho,N}$ as in the previous section on $B_\ell \cup \tau_y B_\ell$, and obtain that in the limit $N \rightarrow \infty$ and then $\varepsilon \rightarrow 0$, f is ultimately constant on the sets with fixed number of particles in $B_\ell \cup \tau_y B_\ell$. Note that the number of particles is not fixed in each of those two boxes, because we allowed jumps from 0 to y . Specifically, denote

$$\widehat{\mathcal{E}}_{\ell,y}^k = \left\{ (\sigma_1, \sigma_2) \in \widehat{\mathcal{E}}_{B_\ell} \times \widehat{\mathcal{E}}_{\tau_y B_\ell}, \quad j(\sigma_1, \sigma_2) = k \right\}$$

the set of ergodic configurations on $B_\ell \cup \tau_y B_\ell$ with k particles, and denote $\nu_{\rho,N}^{y,\ell,k}$ the measure $\nu_{\rho,N}$ conditioned to $\widehat{\mathcal{E}}_{\ell,y}^k$, defined for any $(\sigma_1, \sigma_2) \in \widehat{\mathcal{E}}_{\ell,y}^k$ by

$$\nu_{\rho,N}^{y,\ell,k}(\sigma_1, \sigma_2) := \frac{\nu_{\rho,N}(\eta_{B_\ell} = \sigma_1, \eta_{\tau_y B_\ell} = \sigma_2)}{\nu_{\rho,N}(j(\eta_{B_\ell}, \eta_{\tau_y B_\ell}) = k)}.$$

To prove (7.20), thanks to (7.22), it is therefore enough to prove that

$$\limsup_{\ell \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} \sup_{\sqrt{N} \leq |y| \leq 2\varepsilon N} \sup_k \sum_{(\sigma_1, \sigma_2) \in \widehat{\mathcal{E}}_{\ell,y}^k} \left| \rho^\ell(0)(\sigma_1) - \rho^\ell(y)(\sigma_2) \right| \nu_{\rho,N}^{y,\ell,k}(\sigma_1, \sigma_2) = 0. \quad (7.25)$$

Thanks to (6.12) and Corollary 6.6, (recall formula (6.4) for $\pi_\rho(\xi_{|B_\ell} = \sigma_1)$), we can substitute for any $(\sigma_1, \sigma_2) \in \widehat{\mathcal{E}}_{\ell,y}^k$

$$\begin{aligned} & \frac{\pi_\rho(\xi_{|B_\ell} = \sigma_1) \pi_\rho(\xi_{|B_\ell} = \sigma_2)}{\sum_{(\sigma'_1, \sigma'_2) \in \widehat{\mathcal{E}}_{\ell,y}^k} \pi_\rho(\xi_{|B_\ell} = \sigma'_1) \pi_\rho(\xi_{|B_\ell} = \sigma'_2)} \\ &= \frac{\gamma^{\sigma_1(1) + \sigma_1(\ell) + \sigma_2(y+1) + \sigma_2(y+\ell)}}{\sum_{(\sigma'_1, \sigma'_2) \in \widehat{\mathcal{E}}_{\ell,y}^k} \gamma^{\sigma'_1(1) + \sigma'_1(\ell) + \sigma'_2(y+1) + \sigma'_2(y+\ell)}} \leq \frac{C(\rho)}{\#\{(\sigma'_1, \sigma'_2) \in \widehat{\mathcal{E}}_{\ell,y}^k\}} \end{aligned} \quad (7.26)$$

to $\nu_{\rho,N}^{y,\ell,k}(\sigma_1, \sigma_2)$ in equation (7.25). For $k \leq 2|B_\ell|$, let us introduce the following measure on $\{0, 1, \dots, k\}$:

$$\nu_{k,\ell}(j) = \frac{\#\{(\sigma'_1, \sigma'_2) \in \widehat{\mathcal{E}}_{\ell,y}^k, \sum_{x \in B_\ell} \sigma'_1(x) = j\}}{\#\{(\sigma'_1, \sigma'_2) \in \widehat{\mathcal{E}}_{\ell,y}^k\}},$$

which does not depend on $y \geq \sqrt{N}$. All quantities are now independent from N , ε and y , therefore to prove (7.25) it is sufficient to show that

$$\limsup_{\ell \rightarrow \infty} \sup_k \sum_{j \leq k} \frac{|2j - k|}{|B_\ell|} \nu_{k,\ell}(j) = 0. \quad (7.27)$$

It is now straightforward to project on the possible values at the border of the two configurations, and use Lemma 6.15 to prove that the measure $\nu_{k,\ell}$ is concentrated as $\ell \rightarrow \infty$ on j which is of order $k/2$, in the sense that for any $\delta > 1/2$,

$$\inf_k \nu_{k,\ell}(|j - k/2| \leq \ell^\delta) \xrightarrow{\ell \rightarrow \infty} 1,$$

which proves (7.27) and concludes the proof of the two-blocks estimate. \square

7.4. Dirichlet estimates. We prove here an estimate for the Dirichlet form. Recall that we denote by \mathcal{D}_N the Dirichlet form

$$\begin{aligned}\mathcal{D}_N(g) &= -\nu_{\rho,N}(\sqrt{g}\mathcal{L}_N\sqrt{g}) \\ &= \frac{1}{2}\nu_{\rho,N}\left(\sum_{x\in\mathbb{T}_N}\eta(x+2)\eta(x-1)(\sqrt{g}(\eta^{x,x+1})-\sqrt{g}(\eta))^2\right).\end{aligned}$$

Proposition 7.4 (Dirichlet form estimate). *For any time T , there exists a constant $C = C(T)$ such that*

$$\mathcal{D}_N\left(\frac{1}{T}\int_0^T f_s^N ds\right) \leq \frac{C}{N}.$$

Proof of Proposition 7.4. Let us define the usual relative entropy with respect to $\nu_{\rho,N}$ as

$$\mathcal{H}_\rho^N(f) = \nu_{\rho,N}(f \log f),$$

we have for any $s \in [0, T]$

$$\partial_s \mathcal{H}_\rho^N(f_s^N) = N^2 \nu_{\rho,N}(f_s^N \mathcal{L}_N \log f_s^N) \leq -N^2 \mathcal{D}_N(f_s^N),$$

so that

$$0 \leq \mathcal{H}_\rho^N(f_T^N) + N^2 \int_0^T \mathcal{D}_N(f_s^N) ds \leq \mathcal{H}_\rho^N(f_0^N).$$

Thanks to (6.9), one obtains that for any fixed configuration η in the ergodic component \mathcal{E}_N , we have $-\log \nu_{\rho,N}(\eta) \leq C_\rho N$, so that by convexity of the entropy

$$\mathcal{H}_\rho^N(f_0^N) \leq C_\rho N,$$

and therefore

$$0 \leq \int_0^T \mathcal{D}_N(f_s^N) ds \leq \frac{C}{N}.$$

The convexity of the Dirichlet form \mathcal{D}_N concludes the proof. \square

APPENDIX A. EXIT TIMES FOR RANDOM WALKERS

In this section we recall technical estimates related to exit times for random walkers, which we used in the proof of Lemma 4.6 and Lemma 4.5.

Lemma A.1. *Let us denote by $(X_t)_{t \in \mathbb{R}_+}$ (resp. $(Y_k)_{k \in \mathbb{N}}$) the trajectory of a continuous (resp. discrete) time symmetric random walker on \mathbb{Z} , jumping left and right at the same rate $1/(1+\delta)\ell$ (resp. with the same probability $\frac{1}{2}$). We denote by \mathbb{P}^x the probability laws of both these processes starting from site x . For any integer ℓ , we denote by T_ℓ^X (resp. T_ℓ^Y) the exit time of $B_\ell = \{-\ell, \dots, \ell\}$.*

Then, there exist two constants $C := C(\delta) > 0$ and $\bar{\ell} := \bar{\ell}(\delta) > 0$ such that for any $\ell > \bar{\ell}$,

$$\mathbb{P}^0(T_\ell^X > \ell^4) \leq \exp(-C\ell). \quad (\text{A.1})$$

Furthermore, for any $\varepsilon > 0$, there exist two universal constants $C' > 0$ and $\bar{\ell}'$ such that for any $\ell \geq \bar{\ell}'$

$$\mathbb{P}^0(T_\ell^Y > \ell^{2+\varepsilon}) \leq \exp(-C'\ell^\varepsilon). \quad (\text{A.2})$$

Proof of Lemma A.1. Both of these estimates can be proved in a similar way. Let $k := k_\ell = \lfloor \ell/(1+\delta) \rfloor$, we only need to write

$$\begin{aligned}\mathbb{P}^0(T_\ell^X > \ell^4) &= \sum_{x_1, \dots, x_\ell \in B_\ell} \mathbb{P}^0\left(\{T_\ell^X > \ell^4\} \cap \{X_{(1+\delta)\ell^3} = x_1, X_{2(1+\delta)\ell^3} = x_2, \dots, X_{k(1+\delta)\ell^3} = x_k\}\right)\end{aligned}$$

$$\begin{aligned}
&= \sum_{x_1, \dots, x_k \in B_\ell} \prod_{i=1}^k \mathbb{P}^{x_{i-1}} \left(T_\ell^X > (1 + \delta) \ell^3 \right) \mathbb{P}^{x_{i-1}} \left(X_{(1+\delta)\ell^3} = x_i \mid T_\ell^X > (1 + \delta) \ell^3 \right) \\
&\leq \left(\sup_{x \in B_\ell} \mathbb{P}^x \left(T_\ell^X > (1 + \delta) \ell^3 \right) \right)^k \\
&\leq \left(\sup_{x \in B_\ell} \mathbb{P}^x \left(T_\ell^X > (1 + \delta) \ell^3 \right) \right)^{\lfloor \ell / (1 + \delta) \rfloor}.
\end{aligned}$$

Since X jumps at rate $1/((1 + \delta)\ell)$, the quantity inside parenthesis in the right hand side above converges towards a constant $0 < p(\delta) < 1$. For any $p(\delta) < q < 1$ there exists therefore $\bar{\ell}(q) := \bar{\ell}(\delta)$, such that for any $\ell > \bar{\ell}$,

$$\mathbb{P}^0(T_\ell^X > \ell^4) \leq q^{\ell/(1+\delta)},$$

which proves (A.1).

The second statement is proved in the exact same way, with Y instead of X , and $\ell^{2+\varepsilon}$ instead of ℓ^4 :

$$\begin{aligned}
\mathbb{P}^0(T_\ell^Y > \ell^{2+\varepsilon}) &= \sum_{y_1, \dots, y_{\ell^\varepsilon} \in B_\ell} \mathbb{P}^0 \left(\{T_\ell^Y > \ell^{2+\varepsilon}\} \cap \{Y_{\ell^2} = y_1, Y_{2\ell^2} = y_2, \dots, Y_{\ell^{2+\varepsilon}} = y_{\ell^\varepsilon}\} \right) \\
&= \sum_{y_1, \dots, y_{\ell^\varepsilon} \in B_\ell} \prod_{i=1}^{\ell^\varepsilon} \mathbb{P}^{y_{i-1}} \left(T_\ell^Y > \ell^2 \right) \mathbb{P}^{y_{i-1}} \left(Y_{\ell^2} = y_i \mid T_\ell^Y > \ell^2 \right) \\
&\leq \left(\sup_{y \in B_\ell} \mathbb{P}^y \left(T_\ell^Y > \ell^2 \right) \right)^{\ell^\varepsilon} \\
&\leq \left(\mathbb{P}^0 \left(T_{2\ell}^Y > \ell^2 \right) \right)^{\ell^\varepsilon}.
\end{aligned}$$

Once again, the quantity inside parenthesis converges towards a quantity which is less than one, and the proof of (A.2) is concluded in the same way as for (A.1). \square

REFERENCES

- [1] J. Baik, G. Barraquand, I. Corwin, T. Suidan: *Facilitated exclusion process*, arXiv:1707.01923
- [2] U. Basu and P. K. Mohanty: *Active-absorbing-state phase transition beyond directed percolation: A class of exactly solvable models*, Phys. Rev. E **79**, 041143 (2009).
- [3] O. Blondel, C. Cancès, M. Sasada, M. Simon: *Convergence of a degenerate microscopic dynamics to the porous medium equation*, arXiv:1802.05912.
- [4] T. Funaki: *Free boundary problem from stochastic lattice gas model*, Annales de l'I.H.P. Probabilités et statistiques, Volume 35 (1999) no. 5, p. 573–603.
- [5] T. Funaki, K. Handa, K. Uchiyama: *Hydrodynamic Limit of One-Dimensional Exclusion Processes with Speed Change*, Ann. Probab. **19** no. 1, 245–265 (1991).
- [6] T. Funaki and M. Sasada: *Hydrodynamic Limit for an Evolutional Model of Two-Dimensional Young Diagrams*, Comm. Math. Phys. **299** 335–363 (2010).
- [7] A. Gabel, P. L. Krapivsky, and S. Redner: *Facilitated Asymmetric Exclusion*, Phys. Rev. Lett. **105**, 210603 (2010).
- [8] P. Gonçalves, C. Landim, and C. Toninelli: *Hydrodynamic limit for a particle system with degenerate rates.*, Ann. IHP Probab. Stat. **45** 887-909 (2009).
- [9] F. Hernández, M. Jara, and F. Valentim: *Lattice model for fast diffusion equation*, arXiv:1706.06244.
- [10] K. Jain: *Simple sandpile model of active-absorbing state transitions*, Phys. Rev. E **72**, 017105 (2005).
- [11] C. Kipnis and C. Landim: *Scaling limits of interacting particle systems*. Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences], Volume 320. Springer-Verlag, Berlin (1999).
- [12] S. Lübeck: *Scaling behavior of the absorbing phase transition in a conserved lattice gas around the upper critical dimension*, Phys. Rev. E **64**, 016123 (2001).

- [13] R. Lyons: *Strong laws of large numbers for weakly correlated random variables*, Michigan Math. J. **35**(3), 353–359 (1988).
- [14] M. J. de Oliveira: *Conserved lattice gas model with infinitely many absorbing states in one dimension*, Phys. Rev. E **71**, 016112 (2005)
- [15] M. Rossi, R. Pastor-Satorras, and A. Vespignani: *Universality Class of Absorbing Phase Transitions with a Conserved Field*, Phys. Rev. Lett. **85**, 1803 (2000).
- [16] M. Sasada: *Hydrodynamic limit for particle systems with degenerate rates without exclusive constraints*, Alea **7**, 277–292 (2010).
- [17] J. Sondow: *Problem 11132*, American Mathematics Monthly **112**, 180, (2005).
- [18] S. R. S. Varadhan and H.T. Yau: *Diffusive limit of lattice gas with mixing conditions*, Asian J. Math., Vol.1, No. 4, 623–678 (1997).
- [19] J.L. Vazquez: *The Porous Medium Equation, Mathematical Theory*. Oxford Mathematical Monographs, Clarendon Press, Oxford, 2007.

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