

ON A NON-SOLENOIDAL APPROXIMATION TO THE INCOMPRESSIBLE NAVIER-STOKES EQUATIONS

LORENZO BRANDOLESE

ABSTRACT. We establish an asymptotic profile that sharply describes the behavior as $t \rightarrow \infty$ for solutions to a non-solenoidal approximation of the incompressible Navier–Stokes equations introduced by Temam. The solutions of Temam’s model are known to converge to the corresponding solutions of the classical Navier–Stokes, e.g., in $L^3_{\text{loc}}(\mathbb{R}^+ \times \mathbb{R}^3)$, provided $\epsilon \rightarrow 0$, where $\epsilon > 0$ is the physical parameter related to the artificial compressibility term. However, we show that such model is no longer a good approximation of Navier–Stokes for large times: indeed, its solutions can decay much slower as $t \rightarrow \infty$ than the corresponding solutions of Navier–Stokes.

1. INTRODUCTION

Motivated by recent studies by C. Niche, M. E. Schonbek [13] and W. Rusin [14], we consider the following system, proposed by R. Temam [16] as an useful model for the effective approximation of solutions to the Navier–Stokes equations:

$$(1.1) \quad \begin{cases} \partial_t u^\epsilon - \Delta u^\epsilon + u^\epsilon \cdot \nabla u^\epsilon + \frac{1}{2} u^\epsilon \operatorname{div}(u^\epsilon) - \frac{1}{\epsilon} \nabla \operatorname{div}(u^\epsilon) = 0, & x \in \mathbb{R}^n, t > 0 \\ u(x, 0) = u_0(x), & \operatorname{div}(u_0) = 0. \end{cases}$$

Here $\epsilon > 0$ is a parameter measuring how much the vector field u^ϵ is far from being incompressible. Notice that $\operatorname{div}(u^\epsilon)$, in general, will not be equal to zero for $t > 0$. This way of approaching a Navier–Stokes flow with vector fields that are not necessarily divergence-free has several advantages in numerical simulations, as pointed out in [16]. Indeed, using (1.1) simplifies the discretization procedures, as one does not need to put the divergence-free constraint at each step for the discretized solutions. Another nice feature of this approximation is that it allows to disregard the nonlocal features (the pressure) of the original system.

The construction of global weak solutions to the system (1.1) and the convergence problem as $\epsilon \rightarrow 0$ of these solutions to the corresponding Leray solutions of the classical Navier–Stokes system are successfully addressed in [16] and more recently in [14].

However, the above mentioned papers provide little information on the long time behavior of solutions of the approximated model. The asymptotic behavior as $t \rightarrow +\infty$ for the problem (1.1) has been addressed only recently in [13]. Therein, the authors prove, among other things, that if $u_0 \in L^2(\mathbb{R}^n)$ and if the solution of the linearized equation decays to zero in the L^2 -norm at *slow rates* as $t \rightarrow \infty$, then the solutions to system (1.1) will behave like the solutions of the heat equation: this behavior was somehow the expected one because it is in agreement with the asymptotic behavior for the standard Navier–Stokes equations.

Date: May 9, 2017.

2010 Mathematics Subject Classification. 35Q30, 76D05, 76E19.

Key words and phrases. Asymptotic profiles, long-time asymptotics, artificial compressibility.

Supported by the ANR project DYFICOLTI ANR-13-BS01-0003-01.

On the other hand, when the solution of the linearized equation decays at *faster rates* (namely, faster than $\mathcal{O}(t^{-n/4})$ in the L^2 -norm), then the analysis of [13] does not apply anymore: the reason is explained in Remark 3.2 below. Typically, we encounter such faster decay rates as soon as $u_0 \in L^1(\mathbb{R}^n)$.

For this reason, the problem of the large time behavior of solutions to (1.1) arising from $u_0 \in L^1(\mathbb{R}^n) \cap L^2(\mathbb{R}^n)$, is essentially open.

One might wonder that when $u_0 \in L^1(\mathbb{R}^n) \cap L^2(\mathbb{R}^n)$ the solutions to (1.1) could behave quite differently than the solutions of the Navier–Stokes equations as $t \rightarrow +\infty$. Our analysis will prove that this is indeed the case. One reason for this is that, contrary to the case of the classical Navier–Stokes, the mean $\int u^\epsilon(\cdot, t) dx$ is no longer constant-in-time. The long time behavior of $\int u^\epsilon(\cdot, t) dx$ is itself an interesting problem to address, as one expects its limit to play a predominant role in the description of the long-time asymptotics. These issues are the motivations of the present paper. The main result in this direction, stated in Section 3, is the asymptotic profile obtained in Theorem 3.1 and its applications to sharp two-sided decay estimates for the L^q -norms.

The main tool for describing the long time behavior of solutions to (1.1) will be the construction of asymptotic profiles for solutions of linear integral equations of the form

$$(1.2) \quad u(t) = M(t)u_0 + \int_0^t M(t-s)f(s) ds.$$

Here $M(t)$ is a convolution operator with a kernel satisfying the same scaling properties as the heat kernel, and f is a given forcing term. As such, equation (1.2) is a natural generalization of the classical heat equation in the whole space. The key result for this linear problem is Theorem 2.1. This theorem is widely applicable to other nonlinear equations, besides (1.1).

2. ON A GENERALIZED HEAT EQUATION IN \mathbb{R}^n

Let $M(x, t) \in C^1(\mathbb{R}^n \times (0, +\infty))$ with the same scaling properties as the heat kernel $E(x, t)$, namely,

$$(2.1) \quad M(x, t) = t^{-n/2} M\left(\frac{x}{\sqrt{t}}, 1\right), \quad x \in \mathbb{R}^n, t > 0.$$

We also assume that M satisfies the uniform-in-time spatial decay estimates

$$(2.2) \quad |\partial_x^m M(x, t)| \leq C_m |x|^{-n-m}, \quad m = 0, 1.$$

Estimates (2.2) together with the scaling properties (2.1) imply the uniform-in-space estimates

$$(2.3) \quad |\partial_t^p M(x, t)| \leq C_p t^{-(n+2p)/2}, \quad p = 0, 1.$$

Both decay estimates hold true for the usual heat kernel. For all $1 < q \leq \infty$, and $m, p = 0, 1$ we also deduce

$$(2.4) \quad \|\partial_x^m M(\cdot, t)\|_q = C_{m,q} t^{-\frac{n}{2}(1-\frac{1}{q})-m/2}, \quad \|\partial_t^p M(\cdot, t)\|_q = C_{p,q} t^{-\frac{n}{2}(1-\frac{1}{q})-p}.$$

Next theorem describes the long time behavior of solutions to (1.2) starting from vanishing initial data.

Theorem 2.1.

- (1) Let $n \geq 1$, $f \in L^1(\mathbb{R}^n \times \mathbb{R}^+)$, with $\|f(t)\|_1 = \mathcal{O}(\frac{1}{t})$ as $t \rightarrow +\infty$. Let us introduce the constant $\lambda = \int_0^\infty \int f(y, s) dy ds$ and let also

$$\Phi(x, t) = \int_0^t \int M(x - y, t - s) f(y, s) dy ds.$$

Then, as $t \rightarrow +\infty$,

$$(2.5) \quad \left\| \Phi(t) - \lambda M(\cdot, t) \right\|_q = o(t^{-\frac{n}{2}(1-\frac{1}{q})}), \quad \text{with} \quad \begin{cases} 1 < q \leq \infty & \text{if } n = 1, \\ 1 < q < \frac{n}{n-2} & \text{if } n \geq 2. \end{cases}$$

In particular, when $\lambda \neq 0$, there exist two constants $c_q, c'_q > 0$, independent on f , such that, for $t \gg 1$,

$$(2.6) \quad \lambda c_q t^{-\frac{n}{2}(1-\frac{1}{q})} \leq \|\Phi(t)\|_q \leq \lambda c'_q t^{-\frac{n}{2}(1-\frac{1}{q})}.$$

- (2) The above results (2.5)-(2.6) extend to $q = 1$ if we have in addition $M(\cdot, 1) \in L^1(\mathbb{R}^n)$, and to $\frac{n}{n-2} \leq q \leq \infty$, provided $\|f(t)\|_\beta = \mathcal{O}(t^{-(1+\frac{n}{2}(1-\frac{1}{\beta}))})$ as $t \rightarrow +\infty$, for some β such that $\frac{1}{q} \leq \frac{1}{\beta} < \frac{1}{q} + \frac{2}{n}$.

The above asymptotic expansion, with E instead of M , provides an asymptotic profile for the solution of the heat equation with forcing $\partial_t u = \Delta u + f$ in \mathbb{R}^n , and zero initial data. In this case, the result is valid also for $q = 1$.

Proof of Theorem 2.1. We consider the following decomposition, for any $\eta > 0$ and a suitable $1/2 < a_\eta < 1$ to be chosen later.

$$\begin{aligned} & \Phi(t) - \lambda M(\cdot, t) \\ &= \int_0^t \int M(x - y, t - s) f(y, s) dy ds - M(\cdot, t) \int_0^\infty \int f(y, s) dy ds \\ &= I_1 + I_2 + I_3 + I_4, \end{aligned}$$

where

$$\begin{aligned} I_1 &= -M(x, t) \int_{a_\eta t}^\infty \int f(y, s) dy ds, \\ I_2 &= \int_0^{a_\eta t} \int [M(x, t - s) - M(x, t)] f(y, s) dy ds, \\ I_3 &= \int_0^{a_\eta t} \int [M(x - y, t - s) - M(x, t - s)] f(y, s) dy ds, \quad \text{and} \\ I_4 &= \int_{a_\eta t}^t M(x - y, t - s) f(y, s) dy ds. \end{aligned}$$

Now,

$$\|I_1(t)\|_q \leq C_q t^{-n/2(1-1/q)} \int_{a_\eta t}^\infty \int |f(y, s)| dy ds = o(t^{-\frac{n}{2}(1-\frac{1}{q})}), \quad \text{as } t \rightarrow +\infty.$$

by the assumption $f \in L^1(\mathbb{R}^+ \times \mathbb{R}^n)$ and the dominated convergence theorem.

For the estimate of I_2 we make use of the Taylor formula

$$M(x, t - s) - M(x, s) = - \int_0^1 \partial_t M(x, t - \theta s) \cdot s d\theta.$$

For all $0 \leq s \leq a_\eta t$, using estimate (2.4) we get

$$\begin{aligned} \|M(\cdot, t-s) - M(\cdot, s)\|_q &\leq C_{1,q}(t-\theta s)^{-\frac{n}{2}(1-\frac{1}{q})-1} s \\ &\leq c_{\eta,q} C_{1,q} t^{-\frac{n}{2}(1-\frac{1}{q})-1} s. \end{aligned}$$

Then we get

$$\|I_2(t)\|_q \leq c_{\eta,q} C_{1,q} t^{-\frac{n}{2}(1-\frac{1}{q})-1} \int_0^{a_\eta t} s \|f(\cdot, s)\|_1 ds.$$

But $\int_0^{a_\eta t} s \|f(\cdot, s)\|_1 ds = o(t)$ as $t \rightarrow \infty$, as it can be checked by splitting the last integral, *e.g.*, into $\int_0^{\sqrt{a_\eta t}}$ and $\int_{\sqrt{a_\eta t}}^{a_\eta t}$ and then applying the dominated convergence theorem (we only need $f \in L^1(\mathbb{R}^n \times \mathbb{R}^+)$ here). We thus get

$$\|I_2\|_q = o(t^{-\frac{n}{2}(1-\frac{1}{q})}), \quad \text{as } t \rightarrow +\infty.$$

In order to estimate the third integral we make use of the scaling properties (2.1) of M . For all $1 < q < \infty$ we have

$$\begin{aligned} \|I_3(t)\|_q &\leq \int_0^{a_\eta t} \int (t-s)^{-n/2} \|M(\frac{\cdot-y}{\sqrt{t-s}}, 1) - M(\frac{\cdot}{\sqrt{t-s}}, 1)\|_q |f(y, s)| dy ds \\ &\leq c_{\eta,q} t^{-\frac{n}{2}(1-\frac{1}{q})} \int_0^\infty \int \mathbf{1}_{[0, a_\eta t]}(s) \|M(\cdot - \frac{y}{\sqrt{t-s}}, 1) - M(\cdot, 1)\|_q |f(y, s)| dy ds, \end{aligned}$$

where $\mathbf{1}_S$ denotes the indicator function of the set S . The integrand is dominated by the integrable function $2\|M(\cdot, 1)\|_q |f(y, s)|$. Moreover, for a.e. $(y, s) \in \mathbb{R}^n \times (0, a_\eta t)$, by the continuity under translations of the L^q -norm (or, when $q = \infty$, by the fact that $M(\cdot, t)$ is uniformly continuous), we have

$$\mathbf{1}_{[0, a_\eta t]}(s) \|M(\cdot - \frac{y}{\sqrt{t-s}}, 1) - M(\cdot, 1)\|_q |f(y, s)| \rightarrow 0, \quad \text{as } t \rightarrow \infty.$$

The dominated convergence theorem then yields

$$\|I_3\|_q = o(t^{-\frac{n}{2}(1-\frac{1}{q})}), \quad \text{as } t \rightarrow +\infty.$$

Let us now consider I_4 . Applying the Young inequality we have, for $1 \leq \alpha, \beta, q \leq \infty$,

$$(2.7) \quad \|I_4(t)\|_q \leq \int_{a_\eta t}^t \|M(t-s)\|_\alpha \|f(\cdot, s)\|_\beta ds, \quad 1 + \frac{1}{q} = \frac{1}{\alpha} + \frac{1}{\beta}.$$

Let us first consider the cases $n = 1$, or $n \geq 2$ and $1 \leq q < \frac{n}{n-2}$. This ensures $(n/2)(1-1/q) < 1$. For any $\eta > 0$ small enough, take $a_\eta = 1 - \eta^{1/(1-(n/2)(1-1/q))}$, in a such way that $1/2 < a_\eta < 1$. In these cases we apply the above estimate with $\alpha = q$ and $\beta = 1$. We now make use of the assumption $\|f(t)\|_1 = \mathcal{O}(\frac{1}{t})$ to deduce the estimate, for large enough t ,

$$\begin{aligned} \|I_4\|_q &\leq C_q \int_{a_\eta t}^t (t-s)^{-\frac{n}{2}(1-\frac{1}{q})} s^{-1} ds \\ &\leq C_q t^{-1} \int_{a_\eta t}^t (t-s)^{-\frac{n}{2}(1-\frac{1}{q})} ds \\ &= C_q \eta t^{-\frac{n}{2}(1-\frac{1}{q})}. \end{aligned}$$

As $\eta > 0$ is arbitrarily small, the conclusion follows in this case.

When $n \geq 2$ and $\frac{n}{n-2} \leq q \leq \infty$, we need to take in estimate (2.7) $1 - \frac{2}{n} < \frac{1}{\alpha} \leq 1$, and so $\frac{1}{q} \leq \frac{1}{\beta} < \frac{1}{q} + \frac{2}{n}$. We consider again an arbitrary small $\eta > 0$, but we take now $a_\eta = 1 - \eta^{1/(1-(n/2)(1-1/\alpha))}$. The additional assumption $\|f(t)\|_\beta = \mathcal{O}(t^{-(1+\frac{n}{2}(1-\frac{1}{\beta}))})$ yields the estimate, for large enough t ,

$$\begin{aligned} \|I_4\|_q &\leq C_q \int_{a_\eta t}^t (t-s)^{-\frac{n}{2}(1-\frac{1}{\alpha})} s^{-1-\frac{n}{2}(1-\frac{1}{\beta})} ds \\ &\leq C_q t^{-1-\frac{n}{2}(1-\frac{1}{\beta})} \int_{a_\eta t}^t (t-s)^{-\frac{n}{2}(1-\frac{1}{\alpha})} ds \\ &= C_q \eta t^{-\frac{n}{2}(1-\frac{1}{q})}. \end{aligned}$$

Hence, the conclusion follows also in this case. \square

Calculations in the same spirit as the above were done in [9], in the particular case of the heat kernel. In fact, in [9] the computations were carried for the Navier–Stokes equations.

Theorem 2.1 is the crucial step for establishing Theorem 3.1 below. But the former theorem should be of independent interest, as the scaling methods used for obtaining asymptotic profiles can be applied to other nonlinear problems.

3. STATEMENT OF THE MAIN RESULT

A weak solution to problem (1.1) is a vector field $u^\epsilon \in L^\infty(0, \infty), L^2(\mathbb{R}^n) \cap L^2((0, \infty), \dot{H}^1(\mathbb{R}^n))$ such that for every $\phi \in \mathcal{D}(\mathbb{R}^n, [0, \infty))$, $\operatorname{div} \phi = 0$, we have

$$\int_0^\infty \int \left(\nabla u^\epsilon \cdot \nabla \phi + (u^\epsilon \cdot \nabla u^\epsilon) \cdot \phi + \frac{1}{2} u^\epsilon \cdot \phi \operatorname{div} u^\epsilon - u^\epsilon \cdot \partial_t \phi \right) dx dt = \int u_0 \cdot \phi(\cdot, 0) dx$$

If $u_0 \in L^2(\mathbb{R}^n)$, and $\operatorname{div} u_0 = 0$, then the existence (and the uniqueness for $n = 2$) of a weak solution to problem (1.1) was established in [14, 16]. As for the classical Navier–Stokes equations such solution satisfies the energy inequality

$$(3.1) \quad \int |u^\epsilon(x, t)|^2 dx + 2 \int_s^t \int \left(|\nabla u^\epsilon(x, r)|^2 + \frac{1}{\epsilon} |\operatorname{div} u^\epsilon(x, r)|^2 \right) dx dr \leq \int |u^\epsilon(x, s)|^2 dx,$$

for $s = 0$ and all $t \geq 0$. At least for $2 \leq n \leq 4$, one can also prove the validity of the so-called *strong energy inequality*, that is the validity of (3.1) for almost $s > 0$ and all $t \geq s$.

In the following, we denote the heat kernel by

$$E(x, t) = e^{-|x|^2/(4t)}/(4\pi t)^{n/2}.$$

The usual L^q -estimates for E are

$$\|E(\cdot, t)\|_q = c_q t^{-(n/2)(1-1/q)}, \quad 1 \leq q \leq \infty.$$

Our main result reads as follows:

Theorem 3.1. *Let $2 \leq n \leq 4$, $\epsilon > 0$ and $u_0 \in L^1 \cap L^2(\mathbb{R}^n)$ be a divergence-free vector field. Let u^ϵ be a weak solution to (1.1), satisfying the strong energy inequality (3.1). Then, u^ϵ becomes eventually a strong solution and the limit*

$$(3.2) \quad \vec{\lambda}_\epsilon \equiv \lim_{t \rightarrow +\infty} \int u^\epsilon(x, t) dx$$

does exist and is finite. The vector $\vec{\lambda}_\epsilon$ describes the long time behavior of u^ϵ in the following sense:

$$(3.3) \quad \left\| u^\epsilon(\cdot, t) + E(\cdot, t)\vec{\lambda}_\epsilon \right\|_q = o(t^{-\frac{n}{2}(1-\frac{1}{q})}), \quad \text{with } 1 \leq q \leq \infty.$$

Moreover, if $\vec{\lambda}_\epsilon \neq 0$, then there exist $c_\epsilon, c'_\epsilon > 0$ such that, for $1 \leq q \leq \infty$,

$$(3.4) \quad c_\epsilon |\vec{\lambda}_\epsilon| t^{-\frac{n}{2}(1-\frac{1}{q})} \leq \|u^\epsilon(t)\|_q \leq c'_\epsilon |\vec{\lambda}_\epsilon| t^{-\frac{n}{2}(1-\frac{1}{q})}, \quad \text{for } t \gg 1.$$

The above conclusions holds in any dimension $n \geq 2$ if u_0 belongs to $L^1 \cap L^n(\mathbb{R}^n)$ with $\|u_0\|_n$ small enough and u^ϵ the unique global strong solution to (1.1). The smallness condition on $\|u_0\|_n$ can be replaced by the weaker assumption that $\|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}$ is small enough, for some $n < p < 2n$.

Let us stress the fact that, in general, $\vec{\lambda}_\epsilon \neq 0$ for $\epsilon > 0$. We shall construct in Section 5.2 an explicit example of initial data in the Schwartz class $\mathcal{S}(\mathbb{R}^3)$ such that $\vec{\lambda}_\epsilon \neq (0, 0, 0)$. In particular, this means that these solutions of (1.1) will satisfy $\|u^\epsilon(t)\|_q \sim t^{-\frac{n}{2}(1-\frac{1}{q})}$ for $\epsilon > 0$, when u_0 is integrable. This contrasts with the case $\epsilon = 0$ of the Navier–Stokes equations: indeed, solutions of the Navier–Stokes equations are known to decay as $\|u(t)\|_2 \sim t^{-(n+2)/4}$ as soon as u_0 is well localized, see [17], and sometimes even at faster rates (e.g., under appropriate symmetries). See contribution [3] for an up-to-date review of decay issues for the Navier–Stokes flows.

Remark 3.2. Theorem 3.1 corrects one of the results of the paper [13]. Therein, the authors develop the theory of “decay characters” and give several applications of this notion. One application concerns equation (1.1). The authors asserted that, when $n = 3$, $\|u^\epsilon(t)\|_2 \leq C(1+t)^{-5/4}$, under the assumptions that $u_0 \in L^2(\mathbb{R}^3)$ is divergence-free and satisfies the moment condition $\int(1+|x|)|u_0(x)|dx < \infty$ (in [13] this moment condition is in fact replaced by the closely related condition $r^*(u_0) \geq 1$, where $r^*(u_0)$ is the so-called “decay character” of u_0). This assertion is in contradiction with our lower bounds in (3.4) that shows, for $n = 3$ and $q = 2$ that in general $\|u(t)\|_2 \sim t^{-3/4}$ for large t . It is possible to fix the proof of [13, Theorem 3.9] (the pointwise inequality therein for $G(\xi, t)$ and the subsequent calculations could be easily corrected), but at the price of obtaining a weaker result, namely a lower decay rate in the upper-bound estimate. This could be done using the same ideas as in [13], based on the Fourier-splitting technique. In the present paper we follow however a different approach (somehow inspired by [12]) that has the advantage providing an exact asymptotic profile.

4. GLOBAL STRONG SOLUTIONS UNIFORMLY INTEGRABLE IN TIME

4.1. The linearized equation. The associated linear problem to Eq. (1.1) is

$$(4.1) \quad \begin{cases} \partial_t u^\epsilon - \Delta u^\epsilon - \frac{1}{\epsilon} \nabla \operatorname{div}(u^\epsilon) = g \\ u(x, 0) = u_0(x). \end{cases}$$

The integral formulation associated with this linear problem reads

$$(4.2) \quad u^\epsilon(x, t) = M_\epsilon(t)u_0(x) + \int_0^t M_\epsilon(t-s)g(s) ds,$$

where $M_\epsilon(t)u_0(x)$ is given by the convolution integral

$$M_\epsilon(t)u_0(x) = \int M_\epsilon(x-y, t)u_0(y) dy.$$

The properties of the kernel $M_\epsilon(x, t)$ have been studied in detail by W. Rusin in [14]. Its symbol is

$$(4.3) \quad (\widehat{M}_\epsilon(\xi, t))_{k,l} = e^{-t|\xi|^2} \left(\delta_{k,l} - \frac{\xi_k \xi_l}{|\xi|^2} (1 - e^{-t|\xi|^2/\epsilon}) \right).$$

We complement here the analysis in [14] of the kernel $M_\epsilon(t)$ by observing that, for $t, \epsilon > 0$, the kernel $M_\epsilon(\cdot, t)$ belongs to the Schwartz class. Indeed, we can write

$$(\widehat{M}_\epsilon(\xi, t))_{k,l} = e^{-t|\xi|^2} (\delta_{k,l} + \frac{t}{\epsilon} \xi_k \xi_l H(-t|\xi|^2/\epsilon)),$$

where $H(r) = (e^{-r} - 1)/r = \sum_{k=0}^{\infty} (-r)^k / (k+1)!$. This expressions shows that the map $\xi \mapsto (\widehat{M}_\epsilon(\xi, t))_{k,l}$ can be extended smoothly in a neighborhood of $\xi = 0$, and it is simple to check that its derivatives of any order decay exponentially to zero as $|\xi| \rightarrow \infty$.

We also observe that

$$(4.4) \quad \operatorname{div} u = 0 \quad \Rightarrow \quad M_\epsilon(t)u = e^{t\Delta}u.$$

Another useful identity is the following, valid for any vector field f (not necessarily divergence-free):

$$(4.5) \quad \operatorname{div}(M_\epsilon(t)f) = e^{t(1+1/\epsilon)\Delta}(\operatorname{div} f).$$

It readily follows from the above expressions of the symbol that, for all $\epsilon > 0$,

$$M_\epsilon(x, t) = t^{-n/2} M_\epsilon(x/\sqrt{t}, 1).$$

Moreover, from these scaling properties and the fact that $M(\cdot, t)$ belongs to the Schwartz class, it follows that, for $t > 0$,

$$|\partial_x^m M_\epsilon(x, t)| \leq C_m |x|^{-n-m}, \quad m \in \mathbb{N},$$

where $C_m > 0$ is independent on x and t . Therefore, M_ϵ satisfies all the estimates (2.1)–(2.4), including for $q = 1$.

4.2. The iterative scheme. The problem (1.1) can now be conveniently reformulated in the following integral form:

$$(4.6) \quad u^\epsilon(t) = u_1(t) + B(u^\epsilon, u^\epsilon), \quad u_1(t) = M_\epsilon(t)u_0, \quad \operatorname{div}(u_0) = 0.$$

Here B is the bilinear operator

$$(4.7) \quad B(u, v) = - \int_0^t M_\epsilon(t-s) [u \cdot \nabla v + \frac{1}{2} u \operatorname{div} v](s) ds.$$

The considerations below apply to any abstract equation of the form $u = u_1 + B(u, u)$: let \mathcal{F} be a Banach space, $u_1 \in \mathcal{F}$ and let $B: \mathcal{F} \times \mathcal{F} \rightarrow \mathcal{F}$ be a continuous bilinear operator, with

operator norm $\|B\|$. Let us introduce the nonlinear operators $T_k: \mathcal{F} \rightarrow \mathcal{F}$, $k = 1, 2, \dots$, defined by induction through the formulae

$$\begin{aligned} T_1 &= \text{Id}_{\mathcal{F}} \\ T_k(v) &\equiv \sum_{l=1}^{k-1} B(T_l(v), T_{k-l}(v)), \quad k \geq 2. \end{aligned}$$

Notice that the operator T_k is the restriction to the diagonal of a k -multilinear operator defined from $\mathcal{F}^k = \mathcal{F} \times \dots \times \mathcal{F}$ to \mathcal{F} . The following estimate holds true, see [1], [11]:

$$(4.8) \quad \|T_k(u_1)\|_{\mathcal{F}} \leq \frac{C}{\|B\|} k^{-3/2} (4 \|B\| \|u_1\|_{\mathcal{F}})^k.$$

Thus, under the smallness assumption

$$(4.9) \quad \|u_1\|_{\mathcal{F}} \leq 1/(4\|B\|),$$

the series

$$(4.10) \quad \Psi(u_1) \equiv \sum_{k=1}^{\infty} T_k(u_1),$$

is absolutely convergent in \mathcal{F} and its sum $\Psi(u_1)$ is a solution of the equation $u = u_1 + B(u, u)$. Furthermore, $\Psi(u_1)$ is the only solution in the closed ball $\overline{B_{\mathcal{F}}}(0, \frac{1}{2\|B\|})$ (see [1], [11]).

Coming back to our model (1.1), recalling $u_1(t) = M_{\epsilon}(t)u_0$, we would like to establish the existence and the uniqueness of a solution in a suitable functional setting and to write it as $u = \Phi(u_0)$, where $\Phi(u_0)(t) = \Psi(u_1)(t) = \sum_{k=1}^{\infty} T_k(u_1)$. For later use, we will need the series being absolutely convergent in $L^{\infty}([0, \infty), L^1(\mathbb{R}^d))$. There are several ways to achieve this: the quickest way is to choose \mathcal{F} to be an appropriate subspace of $L^{\infty}([0, \infty), L^1(\mathbb{R}^d))$. A good choice is the following: we define \mathcal{F} to be the Banach space of all $L^{\infty}([0, \infty), L^1(\mathbb{R}^d))$ functions such that $\|f\|_{\mathcal{F}} < \infty$, where

$$(4.11) \quad \begin{aligned} \|f\|_{\mathcal{F}} &= \sup_{t>0} \|f(t)\|_1 + \sup_{t>0} \sqrt{t} \|\nabla_x f(t)\|_1 \\ &+ \sup_{t>0} (1+t)^{n/2} \|f(t)\|_{\infty} + \sup_{t>0} (1+t)^{(n+1)/2} \|\nabla_x f(t)\|_{\infty} < \infty. \end{aligned}$$

We are now in the position of establishing the following proposition:

Proposition 4.1. *There are two constants $\eta > 0$ and $c > 0$ such that if*

$$(4.12) \quad \|u_0\|_1 + \|u_0\|_{\infty} < \eta,$$

then there is a solution $u^{\epsilon} \in \mathcal{F}$ of equation (1.1), such that

$$(4.13) \quad u^{\epsilon} = \Phi(u_0)(t) \equiv \sum_{k=1}^{\infty} T_k(M_{\epsilon}(t)u_0),$$

belonging to and unique in the ball $\{u \in \mathcal{F} : \|u\|_{\mathcal{F}} \leq c\eta\}$. The series (4.13) is absolutely convergent in the \mathcal{F} -norm. In particular, u^{ϵ} is uniformly integral in time and, for all $t \geq 0$,

$$(4.14) \quad \int u^{\epsilon}(x, t) \, dx = \sum_{k=3}^{\infty} \int T_k(M_{\epsilon}(t)u_0) \, dx.$$

Moreover, the limit

$$\lim_{t \rightarrow +\infty} \int u^\epsilon(x, t) \, dx$$

does exist and is finite.

The smallness condition (4.12) is somewhat unpleasant —one usually express smallness conditions in scaling invariant norm, like the L^n -norm or even weaker norms— and can indeed be relaxed. We will indicate in Section 6 how to obtain global solutions belonging to \mathcal{F} replacing condition (4.12) by a much weaker condition on a scaling invariant Besov space, at the price of needing more involved bilinear estimates.

Proof. The bicontinuity of the bilinear operator B is easily proved in this space \mathcal{F} . Indeed, the two L^1 -estimates for $B(u, v)$ and $\nabla B(u, v)$ follow easily applying to M_ϵ the first of (2.4) with $q = 1$ and $m = 0, 1$: this gives

$$\|B(u, v)(t)\|_1 \leq \int_0^t \|u(s)\|_1 \|\nabla v(s)\|_\infty \, ds \lesssim \|u\|_{\mathcal{F}} \|v\|_{\mathcal{F}}$$

and

$$\sqrt{t} \|\nabla B(u, v)(t)\|_1 \leq \sqrt{t} \int_0^t (t-s)^{-1/2} \|u(s)\|_1 \|\nabla v(s)\|_\infty \, ds \lesssim \|u\|_{\mathcal{F}} \|v\|_{\mathcal{F}}.$$

We obtain the two L^∞ estimates for $B(u, v)$ and $\nabla B(u, v)$ by splitting the integral at $t/2$: for the integral $\int_0^{t/2} \dots$ we apply to M_ϵ the first of (2.4) with $q = \infty$ and $m = 0, 1$; for the integral $\int_{t/2}^\infty \dots$ we apply to M_ϵ the first of (2.4) with $q = 1$ and $m = 0, 1$. This implies the required estimate

$$\|B(u, v)\|_{\mathcal{F}} \lesssim C \|u\|_{\mathcal{F}} \|v\|_{\mathcal{F}}.$$

On the other hand, if $u_0 \in L^1 \cap L^\infty(\mathbb{R}^n)$, and if u_0 is divergence-free, then $M_\epsilon(t)u_0 \in \mathcal{F}$, because $M_\epsilon(t)u_0$ agrees with $e^{t\Delta}u_0$. In fact, by the standard heat kernel estimates:

$$\|M_\epsilon(t)u_0\|_{\mathcal{F}} \lesssim \|u_0\|_1 + \|u_0\|_\infty.$$

As discussed before, the above estimates imply that, if $\eta > 0$ is small enough, then the series (4.13) converges in the \mathcal{F} -norm. Its sum is the unique solution of equation (1.1) in a ball of \mathcal{F} centered at zero and with small radius. This series converges in particular in the $L^\infty([0, \infty), L^1(\mathbb{R}^n))$ -norm, so we can integrate (4.13) and exchange the integral and summations symbols. Hence, $\int u^\epsilon(x, t) \, dx = \sum_{k=1}^\infty \int T_k(M_\epsilon(t)u_0) \, dx$.

To deduce (4.14) it only remains to check that $\int u_1(t) \, dx = 0$ and $\int T_2(u_1) \, dx = 0$. But

$$(4.15) \quad \int u_1(t) \, dx = \int M_\epsilon(t)u_0(x) \, dx = \int u_0(x) \, dx = 0,$$

because u_0 is integrable and divergence-free (the last inequality is well known and easy to check using the Fourier transform). Moreover,

$$\begin{aligned} \int T_2(u_1) dx &= \int B(u_1, u_1)(t) dx \\ &= - \int_0^t \int M_\epsilon(t-s) [u_1 \cdot \nabla u_1 + \frac{1}{2} \operatorname{div}(u_1)u_1] dx ds \\ &= - \int_0^t \int [u_1 \cdot \nabla u_1 + \frac{1}{2} \operatorname{div}(u_1)u_1] dx ds \\ &= \frac{1}{2} \int_0^t \int [\operatorname{div}(u_1)u_1](s) dx ds = 0, \end{aligned}$$

because, by (4.5),

$$\operatorname{div}(u_1(s)) = M_\epsilon(s) \operatorname{div}(u_0) = 0.$$

Let us now discuss the existence of the limit $\lim_{t \rightarrow +\infty} \int u^\epsilon(t) dx$. First of all, observe that

$$\int T_k(u_1)(t) dx = - \sum_{l=1}^{k-1} \int_0^t \int [T_l(u_1) \cdot \nabla T_{k-l}(u_1) + \frac{1}{2} \operatorname{div}(T_l(u_1))T_{k-l}(u_1)] dx ds.$$

Moreover, using the definition of the \mathcal{F} -norm we get

$$\|T_l(u_1)(s)\|_2 \|\nabla T_{k-l}(u_1)(s)\|_2 \leq c_{l,k} (1+s)^{-(n+1)/2},$$

so, by Schwarz inequality, the integrand in the last equality is in $L^1(\mathbb{R}^+ \times \mathbb{R}^n)$. Hence the limit

$$\ell_k = \lim_{t \rightarrow +\infty} \int T_k(u_1)(t) dx$$

does exist and equals $-\sum_{l=1}^{k-1} \int_0^\infty \int [T_l(u_1) \cdot \nabla T_{k-l}(u_1) + \frac{1}{2} \operatorname{div}(T_l(u_1))T_{k-l}(u_1)] dx ds$. Moreover, $|\ell_k| \leq \|T_k(u_1)\|_{\mathcal{F}}$ and (4.8) ensures that the series $\sum_{k=3}^\infty \ell_k$ converges under our smallness assumption on u_0 . On the other hand, the convergence of series (4.13) in the $L^\infty(\mathbb{R}^+, L^1(\mathbb{R}^n))$ -norm allows us to exchange in (4.14) the limit as $t \rightarrow +\infty$ with the summation, leading to $\lim_{t \rightarrow +\infty} \int u^\epsilon(t) dx = \sum_{k=3}^\infty \ell_k$. \square

5. LONG TIME BEHAVIOR OF $\int u^\epsilon(x, t) dx$: AN EXPLICIT EXAMPLE WITH $\vec{\lambda}_\epsilon \neq 0$

This section can be skipped on a first reading. Its goal is to prove that in general $\vec{\lambda}_\epsilon \neq 0$. This is interesting to appreciate the relevance of the statement of Theorem 3.1, in particular the interest of the two-sided bound (3.4). But the computations of this section are not needed for proving Theorem 3.1 itself.

5.1. The first term in the expansion of $\int u^\epsilon(t) dx$. We will need an explicit formula for the first term in the right-hand side of (4.14).

$$\begin{aligned} \int T_3(u_1) dx &= \int B(u_1, T_2(u_1)) dx + \int B(T_2(u_1), u_1) dx \\ &= - \int_0^t \int M_\epsilon(t-s) \left[u_1 \cdot \nabla T_2(u_1) + \frac{1}{2} u_1 \operatorname{div}(T_2(u_1)) + \frac{1}{2} T_2(u_1) \operatorname{div}(u_1) \right] + T_2(u_1) \cdot \nabla u_1 \Big] ds dx \\ &= - \int_0^t \int \left[u_1 \cdot \nabla T_2(u_1) + \frac{1}{2} u_1 \operatorname{div}(T_2(u_1)) + T_2(u_1) \cdot \nabla u_1 \right] ds dx, \end{aligned}$$

where we have dropped the term $\frac{1}{2}T_2(u_1)\operatorname{div}(u_1)$ that is identically zero as $\operatorname{div}(u_1) = 0$. But integrating by parts and using again $\operatorname{div}(u_1) = 0$ shows that $\int u_1 \cdot \nabla T_2(u_1) \, dx = 0$, so the first term inside the integral can also be dropped. Another integration by parts finally yields,

$$(5.1) \quad \int T_3(u_1) \, dx = \frac{1}{2} \int_0^t \int [u_1 \operatorname{div}(T_2(u_1))](s) \, ds \, dx.$$

On the other hand,

$$T_2(u_1) = B(u_1, u_1) = - \int_0^t M_\epsilon(t-s)(u_1 \cdot \nabla u_1(s)) \, ds,$$

and so

$$\operatorname{div}(T_2(u_1)) = - \int_0^t e^{(t-s)(1+1/\epsilon)\Delta} \operatorname{div}(u_1 \cdot \nabla u_1)(s) \, ds,$$

where we applied identity (4.5). Replacing this expression in the formula (5.1), next using $u_1(t) = M_\epsilon(t)u_0$ and (4.4) leads to (we omit the summation on the repeated subscripts):

$$(5.2) \quad \int T_3(u_1)(t) \, dx = -\frac{1}{2} \int_0^t \int_0^s e^{s\Delta} u_0 \left[\int_0^s e^{(s-\tau)(1+1/\epsilon)\Delta} \partial_k \left((e^{\tau\Delta} u_{0,h}) (\partial_h e^{\tau\Delta} u_{0,k}) \right) \, d\tau \right] \, ds \, dx.$$

We now apply formula (5.2) to initial data of the following form: $u_0 = \eta v_0$, where v_0 is a fixed divergence-free vector field in $L^1 \cap L^\infty(\mathbb{R}^n)$ and the parameter $\eta > 0$ will be chosen small enough, ensuring in this way the validity of the smallness condition (4.9). For such choice of u_0 , we see that the first term in the summation appearing in the right-hand side of (4.14) satisfies

$$\int T_3(u_1)(t) \, dx = \eta^3 \int T_3(v_1)(t) \, dx, \quad v_1 = M_\epsilon(t)v_0.$$

The sum of all the other terms of (4.14) are $\mathcal{O}(\eta^4)$ as $\eta \rightarrow 0$. Hence,

$$\int u^\epsilon(x, t) \, dx = \eta^3 \int T_3(v_1)(t) \, dx + \mathcal{O}(\eta^4), \quad \text{as } \eta \rightarrow 0.$$

Let us choose a divergence-free vector field $v_0 \in L^1 \cap L^\infty$ such that, for a fixed t , $\int T_3(v_1)(t) \, dx \neq 0$: then it follows that, choosing a $\eta = \eta(t) > 0$ small enough, $\int T_3(u_1)(t) \, dx \neq 0$, and so $\int u^\epsilon(t) \, dx \neq 0$ by (4.14). In the same way, if we choose v_0 in a such way that

$$(5.3) \quad \lim_{t \rightarrow +\infty} \int T_3(v_1)(t) \, dx \neq 0,$$

then, for $u_0 = \eta v_0$ and $\eta > 0$ small enough, we get

$$\lim_{t \rightarrow +\infty} \int u^\epsilon(t) \, dx \neq 0.$$

Now, one expects $\int T_3(v_1)(t) \, dx$ not to be zero generically, but this claim should be proved rigorously. In next subsection, we will content ourselves to construct, by an explicit computation, a simple example of a divergence-free vector field v_0 in the Schwartz class such that the non-vanishing limit condition (5.3) holds.

5.2. **Example of initial data such that $\vec{\lambda}_\epsilon \neq 0$.** Let construct an initial datum satisfying (5.3). Let us consider a divergence-free three-dimensional vector field of the form

$$v_0 = \begin{pmatrix} -\partial_2 g \\ \partial_1 g \\ 0 \end{pmatrix},$$

where g is a smooth and well decaying scalar function to be chosen later on. Expanding the term $\operatorname{div}(v_1 \cdot \nabla v_1)$ for such a vector field v_0 (recalling $v_1(\tau) = M_\epsilon(\tau)v_0 = e^{t\Delta}v_0$) we get after a few simplifications

$$(5.4) \quad \operatorname{div}(v_1 \cdot \nabla v_1)(\tau) = 2\left(\partial_1 \partial_2 E(\cdot, \tau) * g\right)^2 - 2\left(\partial_1^2 E(\cdot, \tau) * g\right)\left(\partial_2^2 E(\cdot, \tau) * g\right).$$

We want to choose g in order to make the computations as explicit as possible, avoiding to take however a radial function (we cannot put too many symmetries, otherwise the integral $\int T_3(v_1) dx$ could vanish). A good choice will be to take a derivative of the gaussian function, as the constants in the subsequent computations can be easily performed in this way by a formal calculus program. For example, let

$$g(x) = \partial_2 E(x, 1).$$

We denote by κ in the following computations a non-zero constant that may change from line to line. Computing the Fourier transform in (5.4) gives (the second equality is computer-assisted, the result can be easily checked e.g. with Maple):

$$\begin{aligned} \mathcal{F}\left[\operatorname{div}(v_1 \cdot \nabla v_1)(\tau)\right](\tau) &= 2\mathcal{F}\left[\left(\partial_1 \partial_2^2 E(\cdot, \tau + 1)\right)^2 - \left(\partial_1^2 \partial_2 E(\cdot, \tau + 1)\right)\left(\partial_2^3 E(\cdot, \tau + 1)\right)\right] \\ &= \frac{\sqrt{2}}{512 \pi^{3/2}} \left((\xi_1^2 \xi_2^2 + \xi_2^4)(\tau + 1) - 3\xi_1^2 - \xi_2^2 \right) \frac{e^{-(\tau+1)|\xi|^2/2}}{(\tau + 1)^{7/2}}. \end{aligned}$$

Applying Plancherel theorem to Eq. (5.2) yields:

$$(5.5) \quad \int T_3(v_1)(x, t) dx = -\kappa \int_0^t \int_0^s \int \frac{e^{-[(s-\tau)(1+1/\epsilon)+s+(\tau+3)/2]|\xi|^2}}{(\tau + 1)^{7/2}} \begin{pmatrix} \xi_2^2 \\ -\xi_1 \xi_2 \\ 0 \end{pmatrix} \left((\xi_1^2 \xi_2^2 + \xi_2^4)(\tau + 1) - 3\xi_1^2 - \xi_2^2 \right) d\xi ds d\tau$$

where $\kappa = \frac{\sqrt{2}}{512 \pi^{3/2}}$. Let us focus on the first component of the above integral (the second and the third components vanish for symmetry reasons for all $t > 0$). Letting $\lambda = (s - \tau)(1 + 1/\epsilon) + s + (\tau + 3)/2$ we are led to calculate (with a computer assisted computation) first the integral in the ξ -variable:

$$\int e^{-\lambda|\xi|^2} \xi_2^2 \left((\xi_1^2 \xi_2^2 + \xi_2^4)(\tau + 1) - 3\xi_1^2 - \xi_2^2 \right) d\xi = -\frac{3 \pi^{3/2}}{4 \lambda^{9/2}} (2\lambda - 3\tau - 3).$$

But $s > \tau$, so $2\lambda - 3\tau - 3 > 0$ and it follows that the first component of $\int T_3(v_1)(x, t) dx$ is strictly positive for all $t > 0$. In fact, the first component of the map $t \mapsto \int T_3(v_1)(x, t) dx$ is strictly increasing and so the first component of the limit $\lim_{t \rightarrow \infty} \int T_3(v_1)(x, t) dx$ is a strictly positive real number.

Summarizing, recalling also the discussion at the beginning of this section, we established that if $\eta > 0$ is small enough and

$$u_0(x) = \eta \begin{pmatrix} -\partial_2^2 E(x, 1) \\ \partial_1 \partial_2 E(x, 1) \\ 0 \end{pmatrix}$$

then the solution of (1.1) starting from u_0 satisfies

$$\vec{\lambda}_\epsilon = \lim_{t \rightarrow \infty} \int u^\epsilon(x, t) dx \neq 0,$$

as the first component of $\vec{\lambda}$ is strictly positive.

6. RELAXING THE SMALLNESS ASSUMPTION

6.1. Global solutions with small data in Besov spaces. Inspired by T. Kato's arguments [10] for the classical Navier–Stokes equations, we introduce the following Banach space, for $n \leq p < \infty$, where we are going to construct our solutions.

(6.1a)

$$\mathcal{X}_p = \{u \in C((0, \infty), L^p(\mathbb{R}^n)) : \|u\|_{\mathcal{X}_p} = \sup_{t>0} t^{(1/2)(1-n/p)} \|u(t)\|_p + \sup_{t>0} t^{1-n/(2p)} \|\nabla u(t)\|_p < +\infty\}.$$

Notice that the \mathcal{X}_p -norm is left invariant by the natural rescaling $u \mapsto u_\lambda = \lambda u(\lambda x, \lambda^2 t)$. Eq: (1.1) is itself left invariant by the above rescaling. If $u_0 \in L^n(\mathbb{R}^n)$, then the solution of the heat equation $e^{t\Delta} u_0$ belongs to \mathcal{X}_p , for $n \leq p \leq \infty$, and

$$\|e^{t\Delta} u_0\|_{\mathcal{X}_p} \leq C_p \|u_0\|_n.$$

The solutions constructed in \mathcal{X}_p will often have faster time decay as $t \rightarrow \infty$ than predicted by the \mathcal{X}_p -norm. The following space will be useful to describe the decay properties of solutions arising from integrable initial data. To this purpose, let us introduce, for $1 \leq q \leq \infty$,

$$(6.2) \quad \mathcal{Y}_q = \{u \in C((0, \infty), L^q(\mathbb{R}^n)) : \|u\|_{\mathcal{Y}_q} = \sup_{t>0} (1+t)^{(n/2)(1-1/q)} \|u(t)\|_q + \sup_{t>0} t^{1/2} (1+t)^{(n/2)(1-1/q)} \|\nabla u\|_q < \infty\}.$$

First of all, observe that, for $n < p \leq \infty$, a tempered distribution in \mathbb{R}^n u_0 belongs to the Besov space $\dot{B}_{p,\infty}^{-1+n/p}$ if and only if $t^{(1/2)(1-n/p)} e^{t\Delta} u_0 \in L^\infty(\mathbb{R}^+, L^p(\mathbb{R}^n))$ and we have the norm equivalence

$$(6.3) \quad \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}} \simeq \sup_{t>0} t^{(1/2)(1-n/p)} \|e^{t\Delta} u_0\|_p.$$

See [4], [11]. The scaling of $\dot{B}_{p,\infty}^{-1+n/p}$ agrees with that of $L^n(\mathbb{R}^n)$ and for $n < p \leq \infty$ we have the inclusion $L^n(\mathbb{R}^n) \subset \dot{B}_{p,\infty}^{-1+n/p}$.

Proposition 6.1.

- (1) *Let $n < p < 2n$ and $u_0 \in \dot{B}_{p,\infty}^{-1+n/p}$ be a divergence-free vector field. There exists $\eta_p > 0$ such that if $\|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}} < \eta$ then there is only one solution $u^\epsilon \in \mathcal{X}_p$ to the problem (1.1), such that $\|u^\epsilon\|_{\mathcal{X}_p} < 2\eta_p$. Such solution u^ϵ belongs also to \mathcal{X}_q , for $p \leq q \leq \infty$. If, more precisely, $u_0 \in L^n(\mathbb{R}^n)$, then u^ϵ belongs also to \mathcal{X}_q , for $n \leq q \leq \infty$ and in this case $u^\epsilon \in C([0, \infty), L^n(\mathbb{R}^n))$.*

(2) Under the additional assumption $u_0 \in L^1(\mathbb{R}^n)$, then, for all $1 \leq q \leq \infty$, $u^\epsilon \in \mathcal{Y}_q$.

Remark 6.2. When $u_0 \in L^n(\mathbb{R}^n)$ with a small enough L^n -norm, then Proposition 6.1 (1) does apply. This follows from the fact the inclusion map $L^n(\mathbb{R}^n) \subset \dot{B}_{p,\infty}^{-1+n/p}$, $n < p \leq \infty$, is continuous. Proposition 6.1, however allows to construct global solution for some initial data with large $L^n(\mathbb{R}^n)$ -norm. Indeed, the smallness condition on the Besov norm $\dot{B}_{p,\infty}^{-1+n/p}$ will be fulfilled as soon as u_0 is fast oscillating. This idea of relaxing the smallness condition using rough spaces goes back to [6].

Proof. The proof of the first part is conceptually close to that in [10] or [6], with slight changes in the choice of the exponents of some relevant estimates. For this reason we will be rather sketchy, omitting in particular to discuss the continuity in the time variable that is standard. The proof consists in applying the standard fixed point argument in \mathcal{X}_p to the integral equation (4.6).

If $u_0 \in \dot{B}_{p,\infty}^{-1+n/p}$, then $\nabla u_0 \in \dot{B}_{p,\infty}^{-2+n/p}$ and therefore

$$(6.4) \quad \sup_{t>0} t^{1-n/(2p)} \|\nabla e^{t\Delta} u_0\|_p \lesssim \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}.$$

As u_0 is divergence-free, $M_\epsilon(t)u_0 = e^{t\Delta}u_0$, as already observed in Eq. (4.4). Hence,

$$(6.5) \quad \|M_\epsilon(t)u_0\|_{\mathcal{X}_p} \leq C_p \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}, \quad n < p \leq \infty.$$

The estimates on the bilinear term rely on Young convolution inequality and the estimates, valid for $1 \leq r \leq \infty$ (see the discussion on M_ϵ at the end of Subsection 4),

$$(6.6) \quad \|M_\epsilon(t-s)\|_r \simeq (t-s)^{-\frac{n}{2}(1-1/r)}, \quad \|\nabla M_\epsilon(t-s)\|_r \simeq (t-s)^{-1/2-\frac{n}{2}(1-1/r)}.$$

More precisely, choosing $r = p'$ (the conjugate exponent of p), with $n < p < 2n$, we get

$$\|B(u, v)(t)\|_p \simeq \int_0^t (t-s)^{-n/(2p)} \|u(s)\|_p \|\nabla v(s)\|_p ds \lesssim t^{-(1/2)(1-n/p)},$$

and

$$\|\nabla B(u, v)(t)\|_p \simeq \int_0^t (t-s)^{-1/2-n/(2p)} \|u(s)\|_p \|\nabla v(s)\|_p ds \lesssim t^{-(1-n/(2p))}.$$

Combining these two estimates we see that, for some constant $C'_p > 0$ and all u, v :

$$(6.7) \quad \|B(u, v)\|_{\mathcal{X}_p} \leq C'_p \|u\|_{\mathcal{X}_p} \|v\|_{\mathcal{X}_p}, \quad n < p < 2n.$$

Applying the fixed point lemma ([4, Lemma 5.5]), using estimates (6.5) and (6.7) yields the existence and the unicity of the solution in the ball $\{u: \|u\|_{\mathcal{X}_p} < 1/2(C_p C'_p)\}$ of the space \mathcal{X}_p , provided $\|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}} < 1/(4C_p C'_p)$, and $n < p < 2n$.

Let us now prove that such solution belongs to \mathcal{X}_q , for all $n < p \leq q \leq \infty$ (or $n \leq q \leq \infty$ if we assume $u_0 \in L^n(\mathbb{R}^n)$). We write $u^\epsilon = M_\epsilon(t)u_0 + B(u^\epsilon, u^\epsilon)$ and separate the contributions of the linear and the nonlinear terms.

First of all, by the semigroup properties of the heat kernel, $\|M_\epsilon(t)u_0\|_{\mathcal{X}_q} \lesssim \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}$. Under the more stringent condition $u_0 \in L^n(\mathbb{R}^n)$ the linear term satisfies also $\|M_\epsilon(t)u_0\|_{\mathcal{X}_n} \lesssim \|u_0\|_{L^n(\mathbb{R}^n)}$. Let us prove that the nonlinear term $B(u^\epsilon, u^\epsilon)$ belongs to $\mathcal{X}_n \cap \mathcal{X}_\infty$ (whether or not $u_0 \in L^n(\mathbb{R}^n)$).

Applying twice (6.6) with r such that $1 + 1/n = 1/r + 2/p$ we get, for $n < p < 2n$:

$$\|B(u^\epsilon, u^\epsilon)(t)\|_n + t^{1/2} \|\nabla B(u^\epsilon, u^\epsilon)(t)\|_n \lesssim \|u^\epsilon\|_{\mathcal{X}_p}^2 \lesssim \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}},$$

and so

$$(6.8) \quad \|u^\epsilon\|_n + t^{1/2}\|\nabla u^\epsilon\|_n \lesssim \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}.$$

It only remains to establish the L^∞ estimates for u^ϵ and ∇u^ϵ . The former is easy: applying (6.6) with r such that $1 = 1/r + 2/p$, we get

$$t^{1/2}\|B(u^\epsilon, u^\epsilon)(t)\|_\infty \lesssim \|u^\epsilon\|_{\mathcal{X}_p}^2,$$

implying

$$(6.9) \quad t^{1/2}\|u^\epsilon(t)\|_\infty \lesssim \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}.$$

We perform the latter in two steps: first, applying (6.6) for all $n < p \leq q < \tilde{p} < \infty$, with $1 + 1/q = 1/r + 1/\tilde{p} + 1/n$,

$$\begin{aligned} \|\nabla B(u^\epsilon, u^\epsilon)(t)\|_q &\leq \int_0^t (t-s)^{-1-\frac{n}{2}(1/\tilde{p}-1/q)} \|u^\epsilon(s)\|_{\tilde{p}} \|\nabla u^\epsilon(s)\|_n \\ &\lesssim t^{-1+n/(2q)} \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}, \end{aligned}$$

where we used $\|u^\epsilon(t)\|_{\tilde{p}} \lesssim t^{-(1/2)(1-n/\tilde{p})} \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}$, that follows interpolating the two previous estimates on $\|u^\epsilon(t)\|_p$ and $\|u^\epsilon(t)\|_\infty$. Using the above with $q = 2n$ we get

$$\|\nabla u^\epsilon(t)\|_{2n} \lesssim t^{-3/4} \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}.$$

Next,

$$\begin{aligned} \|\nabla B(u^\epsilon, u^\epsilon)(t)\|_\infty &\leq \int_0^{t/2} (t-s)^{-1/2-n/p} \|u^\epsilon(s)\|_p \|\nabla u^\epsilon\|_p ds \\ &\quad + \int_{t/2}^t (t-s)^{-3/4} \|u^\epsilon(s)\|_\infty \|\nabla u^\epsilon(s)\|_{2n} ds \\ &\lesssim t^{-1} \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}. \end{aligned}$$

We finally get $\|\nabla u^\epsilon(t)\|_\infty \lesssim t^{-1} \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}$. Summarizing, u^ϵ belongs to $\mathcal{X}_p \cap \mathcal{X}_\infty$ and by interpolation for all $p \leq q \leq \infty$ and $n < p < 2n$ (with $q = n$ being allowed under the more stringent assumption $u_0 \in L^n(\mathbb{R}^n)$),

$$\|u^\epsilon\|_{\mathcal{X}_q} \lesssim \|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}},$$

for all divergence-free distribution u_0 such that $\|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}} < 1/(4C_p C'_p) = \eta_p$. This achieves the proof of Part (1) of Proposition 6.1.

Let us now prove Part (2), of Proposition 6.1. We have now the additional assumption $u_0 \in L^1(\mathbb{R}^n)$. More in general, we can prove that under the additional assumption $u_0 \in L^r(\mathbb{R}^n)$, with $1 \leq r < n$, then the solution constructed in Part (1) satisfies the estimates

$$(6.10) \quad \begin{aligned} (1+t)^{(n/2)(1/r-1/q)} u^\epsilon &\in L^\infty((0, \infty), L^q(\mathbb{R}^n)) \quad \text{and} \\ t^{1/2}(1+t)^{(n/2)(1/r-1/q)} \nabla u^\epsilon &\in L^\infty((0, \infty), L^q(\mathbb{R}^n)), \end{aligned}$$

provided $\|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}} < \eta'$ and $\eta' > 0$ is small enough. Notice that, at least for $1 < r < n$ and $r \leq q < \infty$ this type of estimates were proved in [10, Theorem 5] for the Navier–Stokes equations.

Moreover, Kato's recursive inequalities method provides the norm estimate, for $1 < r < n$ and $r \leq q < \infty$ (and $\epsilon = 0$, *i.e.* in the Navier–Stokes case),

$$(6.11) \quad \sup_{t \geq 1} (1+t)^{(n/2)(1/r-1/q)} \|u^\epsilon\|_q + t^{1/2} (1+t)^{(n/2)(1/r-1/q)} \|\nabla u^\epsilon\|_q \lesssim \|u_0\|_r + \eta'.$$

In fact, for $1 < r < n$ and $r \leq q < \infty$ Kato's proof can be reproduced in our case (*i.e.* in the case $\epsilon > 0$) in exactly the same way: indeed the linear $L^r - L^q$ linear estimates for $M_\epsilon(t)u_0$ agree with the heat kernel estimates, and the nonlinearity in Eq. (4.7) has the same structure as in the integral Navier–Stokes equations. For this reason we can skip the proof that (6.11) hold true also for solutions of (4.6). What remains to be done, is to establish (6.10) in the limit case $r = 1$ or $q = \infty$ that were excluded in Kato's approach. This is possible because, contrary to the case of Navier–Stokes, we do not have to deal with the projection operator onto the divergence-free vector field, that is unbounded in L^1 and L^∞ . In other words, we take advantage of the fact that the kernel of $M_\epsilon(t)$ belongs to $L^1 \cap L^\infty(\mathbb{R}^n)$. More precisely, using (6.6) with $r = 1$ and applying (6.10) with $r = 6/5$ and $q = 2$ we get:

$$\begin{aligned} \|u^\epsilon(t)\|_1 &\leq C \|u_0\|_1 + \int_0^t \|u^\epsilon\|_2 \|\nabla u^\epsilon\|_2 \, ds \leq C, \\ \|\nabla u^\epsilon(t)\|_1 &\leq \|u_0\|_1 t^{-1/2} + \int_0^t (t-s)^{-1/2} \|u^\epsilon\|_2 \|\nabla u^\epsilon\|_2 \, ds \leq C t^{-1/2}. \end{aligned}$$

This establishes (6.10) for $r = 1$ and $q = 1$. Using the similar arguments one can easily establish the validity of (6.10) for $r = 1$ and $q = \infty$. By interpolation, these estimates are then valid for $r = 1$ and $1 \leq q \leq \infty$, hence Part (2) of Proposition 6.1 follows. \square

6.2. Application to weak solutions. The uniqueness of weak solutions for $n = 2$ for (1.1) is due to R.Temam [16]. When $n \geq 3$, the uniqueness of weak solutions is not known. In this case, the analogue of classical weak-strong uniqueness result established for the Navier–Stokes equations, like that of Sohr and von Wahl [15] for (1.1) will be useful. We state this as a remark:

Remark 6.3. Let u^ϵ and v^ϵ be two weak solutions as defined in Section 7, satisfying the energy inequality (3.1) (at least in its weak form, *i.e.* for $s = 0$) and let $v^\epsilon \in C([0, T], L^n(\mathbb{R}^n))$, $0 < T \leq \infty$. Then $u^\epsilon = v^\epsilon$.

Indeed, writing the equation for the difference $w^\epsilon = u^\epsilon - v^\epsilon$, formally multiplying by w^ϵ using the cancellation

$$\int (u^\epsilon \cdot \nabla) w^\epsilon \cdot w^\epsilon + \frac{1}{2} \int (w^\epsilon \operatorname{div} u^\epsilon) \cdot w^\epsilon = 0,$$

one obtains from the constructions of the solutions that w^ϵ satisfies the inequality

$$\frac{1}{2} \|w^\epsilon(t)\|_2^2 + \int_0^t \int (w^\epsilon \cdot \nabla v^\epsilon) \cdot w^\epsilon + \frac{1}{2} \int_0^t \int v^\epsilon (\operatorname{div} w^\epsilon) \cdot w^\epsilon + \frac{1}{\epsilon} \int_0^t \|\operatorname{div} w^\epsilon\|_2^2 + \int_0^t \|\nabla w^\epsilon\|_2^2 \leq 0.$$

But

$$\left| \int_0^t \int (w^\epsilon \cdot \nabla v^\epsilon) \cdot w^\epsilon + \frac{1}{2} \int_0^t \int v^\epsilon (\operatorname{div} w^\epsilon) \cdot w^\epsilon \right| \leq C \int_0^t \int |v^\epsilon| |w^\epsilon| |\nabla w^\epsilon|.$$

We can now forget about the term $\frac{1}{\epsilon} \int_0^t \|\operatorname{div} w^\epsilon\|_2^2$ and proceed as for the classical Navier–Stokes: splitting $v^\epsilon = \bar{v}^\epsilon + \tilde{v}^\epsilon$ with \bar{v}^ϵ small in the $L^\infty((0, T), L^n)$ -norm and $\tilde{v}^\epsilon \in L^\infty((0, T), L^\infty)$ and

we can absorb the above term (see [11] for details of these estimates) and get $\|w^\epsilon(t)\|_2^2 + C\|\tilde{v}^\epsilon\|_{L_{t,x}^\infty} \int_0^t \|w(s)\|_2^2 \leq 0$. As usual one concludes $w^\epsilon = 0$ by Gronwall inequality.

7. PROOF OF THE MAIN RESULT

The proof of Theorem 3.1 relies on construction of weak solutions made in [14],[16], on Proposition 6.1 and Theorem 2.1.

Proof of Theorem 3.1. Let $2 \leq n \leq 4$ and consider a weak solution u^ϵ to (1.1) satisfying the strong energy inequality (3.1), arising from the divergence-free vector field $u_0 \in L^2(\mathbb{R}^n)$. Then, for almost all t_0 , we have $u^\epsilon(t_0) \in H^1(\mathbb{R}^n) \subset L^n(\mathbb{R}^n)$, and we can find $t_0 > 0$ such that the energy inequality (3.1) holds with $s = t_0$ and $\|u(t_0)\|_n$ satisfies the smallness assumption mentioned in Remark 6.2. By this remark, there is a strong solution $v^\epsilon \in C([t_0, \infty), L^n(\mathbb{R}^n))$. As observed in Remark 6.3, $u^\epsilon = v^\epsilon$ on $[t_0, \infty)$. So u^ϵ becomes a strong solution after some time, satisfying for $t \geq t_0$ the conditions of Part (1) of Proposition 6.1. But u^ϵ is also known to satisfy the integral equation (4.6), just like weak solutions of the Navier–Stokes equations satisfying the energy inequality do solve the corresponding integral equations (we refer to [7] for a proof this claim in the more general setting of L_{loc}^2 -solutions). Then the bilinear term $B(u^\epsilon, u^\epsilon)(t)$ in (4.6) belongs to $L^1(\mathbb{R}^n)$. Of course, under the more stringent condition $u_0 \in L^1 \cap L^2(\mathbb{R}^n)$ also the linear term $M_\epsilon(t)u_0$ in (4.6) belongs to $L^1(\mathbb{R}^n)$. Hence, when $2 \leq n \leq 4$, under the assumptions of Theorem 3.1, the weak solution considered above, after some time $t_0 > 0$, satisfies in fact also the conditions of Part (2) of Proposition 6.1. Namely

$$(7.1) \quad \begin{aligned} \|u^\epsilon(t)\|_q &\leq C_q t^{-(n/2)(1-1/q)}, \\ \|\nabla u^\epsilon(t)\|_q &\leq C_q t^{-1/2-(n/2)(1-1/q)}, \end{aligned} \quad t \geq t_0, \quad 1 \leq q \leq \infty.$$

It follows from Schwarz inequality that $|u^\epsilon| |\nabla u^\epsilon| \in L^1(\mathbb{R}^n \times \mathbb{R}^+)$. In particular, we can define

$$(7.2) \quad \vec{\lambda}_\epsilon \equiv \int_0^\infty \int \left(u^\epsilon \cdot \nabla u^\epsilon + \frac{1}{2} u^\epsilon \operatorname{div}(u^\epsilon) \right) dy ds.$$

Integrating with respect to the space variable equation (4.6) and recalling that $\int M_\epsilon(t)u_0 dx = 0$ (see (4.15)), we deduce

$$\int u^\epsilon(x, t) dx = \int_0^t \int \left(u^\epsilon \cdot \nabla u^\epsilon + \frac{1}{2} u^\epsilon \operatorname{div}(u^\epsilon) \right) dy ds,$$

and so the limit $\lim_{t \rightarrow +\infty} \int u^\epsilon(x, t) dx$ does exist. Then $\vec{\lambda}_\epsilon$ is also given by

$$(7.3) \quad \vec{\lambda}_\epsilon = \lim_{t \rightarrow +\infty} \int u^\epsilon(x, t) dx.$$

This is in agreement with (3.2).

Consider again the integral formulation (4.6), and let us start with the study of the asymptotics of the linear part. As already observed, $M_\epsilon(t)$ boils down to the standard heat kernel when applied to divergence-free vector fields and $\int u_0 = 0$, by the divergence-free condition and the integrability of u_0 . Hence,

$$M_\epsilon(t)u_0(x) = e^{t\Delta}u_0(x) = \int [E(x-y, t) - E(x, t)]u_0(y) dy.$$

Then, as $t \rightarrow +\infty$,

$$\|M_\epsilon(t)u_0(x)\|_q \leq t^{-\frac{n}{2}(1-\frac{1}{q})} \int \|E(\cdot - \frac{y}{\sqrt{t}}, 1) - E(\cdot, 1)\|_q |u_0(y)| dy = o(t^{-\frac{n}{2}(1-\frac{1}{q})}).$$

In the last equality we applied the continuity of the L^q -norm under translations (when $1 \leq q < \infty$) or the uniform continuity of $E(\cdot, t)$ (when $q = \infty$), and the dominated convergence theorem to prove that the last integral goes to zero as $t \rightarrow \infty$. The nonlinear part in (4.6) can be written as

$$\Phi(\cdot, t) = - \int_0^t \int M_\epsilon(x-y, t-s) f(y, s) dx ds,$$

where $f(s) = u^\epsilon \cdot \nabla u^\epsilon + \frac{1}{2} u^\epsilon \operatorname{div}(u^\epsilon)$.

Next step will consist in applying Theorem 2.1 with this choice of f and the kernel $M = M_\epsilon$. Let us check the validity of the assumptions of our proposition. We already observed in Section 4 that $M_\epsilon(\cdot, t)$ is in $L^1(\mathbb{R}^n)$. Moreover, applying the decay estimates (7.1) $f \in L^1(\mathbb{R}^n \times \mathbb{R}^+)$ and $t\|f(t)\|_1 \in L^\infty(\mathbb{R}^+)$. Furthermore combining estimates (7.1) with Hölder inequality we get, for $1 \leq \beta \leq \infty$,

$$\|f(t)\|_\beta = \|u^\epsilon(t)\|_{2\beta} \|\nabla u^\epsilon(t)\|_{2\beta} \leq C_\epsilon t^{-n+\frac{n}{2\beta}-\frac{1}{2}} = \mathcal{O}(t^{-(1+\frac{n}{2}(1-\frac{1}{\beta}))}) \quad \text{as } t \rightarrow +\infty.$$

Therefore, for all $1 \leq q \leq \infty$, both Part (1) and Part (2) of Theorem 2.1 do apply.

Hence, recalling expression (7.2),

$$(7.4) \quad \left\| \Phi(\cdot, t) + M_\epsilon(\cdot, t) \vec{\lambda}_\epsilon \right\|_q = o(t^{-\frac{n}{2}(1-\frac{1}{q})}), \quad \text{with } 1 \leq q \leq \infty.$$

Combining the integral formula (4.6) with the above results we get the asymptotics

$$\left\| u^\epsilon(t) + M_\epsilon(\cdot, t) \vec{\lambda}_\epsilon \right\|_q = o(t^{-\frac{n}{2}(1-\frac{1}{q})}), \quad 1 \leq q \leq \infty.$$

On the other hand, from (4.4).

$$M_\epsilon(\cdot, t) \vec{\lambda}_\epsilon = e^{t\Delta} \vec{\lambda}_\epsilon = E(\cdot, t) \vec{\lambda}_\epsilon.$$

and we readily obtain profile (3.3). To deduce from (3.3) the upper and lower bounds (3.4) we just write

$$\|u^\epsilon(t)\|_q = \|E(\cdot, t) \vec{\lambda}_\epsilon\|_q t^{-n/2(1-1/q)} + o(t^{-n/2(1-1/q)}) \quad \text{as } t \rightarrow +\infty.$$

and observe that there exist two constants $c_\epsilon, c'_\epsilon > 0$ such that, for all $1 \leq q \leq \infty$,

$$c'_\epsilon |\vec{\lambda}_\epsilon| t^{-n/2(1-1/q)} \leq \|E(\cdot, t) \vec{\lambda}_\epsilon\|_q \leq c_\epsilon |\vec{\lambda}_\epsilon| t^{-n/2(1-1/q)}.$$

Estimates (3.4) now follows.

The asymptotic profile (3.3) and its corollary (3.4) remain valid without restrictions on the spatial dimension, provided $u_0 \in L^1 \cap L^n(\mathbb{R}^n)$, with $\|u_0\|_{\dot{B}_{p,\infty}^{-1+n/p}}$ small enough for some $n < p < 2n$. Indeed, in this case we can directly apply Proposition (6.1) and obtain the validity of estimates (7.1) for all $t > 0$ (in fact, $u^\epsilon \cap \mathcal{Y}_1 \cap \mathcal{Y}_\infty$ that provides a better control of u^ϵ near $t = 0$ than (7.1)) and argue as above. □

Let us mention that in Theorem 3.1 it would be possible to replace the condition $u_0 \in L^1(\mathbb{R}^n)$ (that implies $\int u_0 = 0$ because u_0 is divergence-free) by a more general condition prescribing a fast L^q -decay of the linear part of the equation. For example, by the condition

$$(7.5) \quad \lim_{t \rightarrow +\infty} t^{(n/2)(1-1/q)} \|e^{t\Delta} u_0\|_q = 0.$$

Conclusions (3.3)-(3.4) would remain valid. However, formula (3.2) does require the spatial integrability of the solution and cannot be used in these more general situations. In the absence of the L^1 condition for u_0 , formula (7.2) can be used instead of (3.2) to define $\vec{\lambda}_\epsilon$.

It could be also possible to replace the L^1 -condition on u_0 by a condition involving the *decay character* of u_0 , namely, $r^*(u_0) > 0$. The decay character is very useful to get sharp algebraic decay estimates from below and above for (linear or nonlinear) dissipative systems. We do not recall here the precise definition of decay character of an L^2 function: the original definition [5, 13] has been slightly changed and improved in [2], in order to make the theory more complete and widely applicable. See also [8] for another very recent application of the decay character.

The case $r^*(u_0) < 0$ is also of interest, but this case corresponds to solutions such that the linear part decays at slow rates in L^2 : as we pointed out in the introduction, in this case, the analysis of [13] already provides a satisfactory answer to the large time decay problem for equation (1.1).

The case $r^*(u_0) = 0$ corresponds to a borderline situation: this condition ensures $ct^{-n/4} \leq \|e^{t\Delta} u_0\|_2 \leq c't^{-n/4}$ for $t \gg 1$ (with $c, c' > 0$). As u_0 is divergence-free, condition $r^*(u_0) = 0$ thus excludes $u_0 \in L^1(\mathbb{R}^n)$ and excludes also (7.5) (but condition $r^*(u_0) = 0$ is compatible, e.g., with $u_0 \in \dot{B}_{2,\infty}^{-n/2}(\mathbb{R}^n)$). In this situation, if $\vec{\lambda}_\epsilon = 0$ then the nonlinear integral term decays faster to zero than $\|e^{t\Delta} u_0\|_2$. So the linear part will govern the long time behavior of the solution. But for $\vec{\lambda}_\epsilon \neq 0$, both the linear and nonlinear terms have the same decay rates. In this borderline situation, solutions to (1.1) satisfy $\|u(t)\|_2 = \mathcal{O}(t^{-n/4})$ but lower bounds for the L^2 -decay of u are no longer available.

8. ACKNOWLEDGEMENTS

The author would like to thank the referees for their careful reading and helpful suggestions in improving the first version of the manuscript.

REFERENCES

- [1] P. Auscher and P. Tchamitcian, *Espaces critiques pour le systme des equations de Navier-Stokes incompressibles*, arXiv:0812.1158 (2008).
- [2] L. Brandolese, *Characterization of solutions to dissipative systems with sharp algebraic decay*, SIAM J. Math. Anal. **48** (2016), no. 3, 1616–1633.
- [3] L. Brandolese and M. E. Schonbek, *Large Time Behavior of the NavierStokes Flow*. In: *Handbook of Mathematical Analysis in Mechanics of Viscous Fluids. Part I. Incompressible fluids. Unsteady viscous Newtonian fluids*. Yoshikazu Giga and Antonin Novotný editors, Springer (2016), 1-67.
- [4] H. Bahouri, J.-Y. Chemin, and R. Danchin, *Fourier analysis and nonlinear partial differential equations*, Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences], vol. 343, Springer, Heidelberg, 2011.
- [5] C. Bjorland and M.E. Schonbek, *Poincaré's inequality and diffusive evolution equations*, Adv. Differential Equations **14** (2009), no. 3-4, 241–260.

- [6] M. Cannone, *Ondelettes, paraproducts et Navier-Stokes*, Diderot Editeur, Paris, 1995. With a preface by Yves Meyer.
- [7] S. Dubois, *What is a solution to the Navier-Stokes equations?*, C. R. Math. Acad. Sci. Paris **335** (2002), no. 1, 27–32 (English, with English and French summaries).
- [8] L.C.F. Ferreira, C.J. Niche, and G. Planas, *Decay of solutions to dissipative modified quasi-geostrophic equations*, Proceedings of the American Mathematical Society **145** (2017), 287–301.
- [9] Y. Fujigaki and T. Miyakawa, *Asymptotic profiles of nonstationary incompressible Navier-Stokes flows in the whole space*, SIAM J. Math. Anal. **33** (2001), no. 3, 523–544.
- [10] T. Kato, *Strong L^p -solutions of the Navier-Stokes equation in \mathbf{R}^m , with applications to weak solutions*, Math. Z. **187** (1984), no. 4, 471–480.
- [11] P. G. Lemarié-Rieusset, *Recent developments in the Navier-Stokes problem*, Chapman & Hall/CRC Research Notes in Mathematics, vol. 431, Chapman & Hall/CRC, Boca Raton, FL, 2002.
- [12] T. Miyakawa and M.E. Schonbek, *On optimal decay rates for weak solutions to the Navier-Stokes equations in \mathbb{R}^n* , Proceedings of Partial Differential Equations and Applications (Olomouc, 1999), 2001, pp. 443–455.
- [13] C. Niche and M.E. Schonbek, *Decay characterization to dissipative equations*, J. London Math. Soc. (2) **91** (2015), 573–595.
- [14] W. Rusin, *Incompressible 3D Navier-Stokes equations as a limit of a nonlinear parabolic system*, J. Math. Fluid Mech. **14** (2012), no. 2, 383–405.
- [15] H. Sohr and W. von Wahl, *On the singular set and the uniqueness of weak solutions of the Navier-Stokes equations*, Manuscripta Math. **49** (1984), no. 1, 27–59.
- [16] R. Temam, *Une méthode d’approximation de la solution des équations de Navier-Stokes*, Bull. Soc. Math. France **96** (1968), 115–152 (French).
- [17] M. Wiegner, *Decay results for weak solutions of the Navier-Stokes equations on \mathbf{R}^n* , J. London Math. Soc. (2) **35** (1987), no. 2, 303–313.

L. BRANDOLESE: UNIVERSITÉ DE LYON ; UNIVERSITÉ LYON 1 ; CNRS UMR 5208 INSTITUT CAMILLE JORDAN, 43 BD. DU 11 NOVEMBRE, VILLEURBANNE CEDEX F-69622, FRANCE.

E-mail address: Brandolese@math.univ-lyon1.fr

URL: <http://math.univ-lyon1.fr/~brandolese>