

# Field theory approach to off critical $SLE_2$ and $SLE_4$

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Renormalization: algebraic, geometric and probabilistic aspects

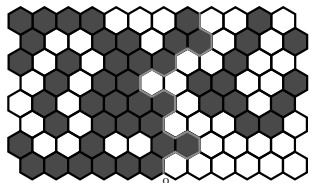
Based on work in collaboration with M. Bauer & D. Bernard

# Outline

- 1 Introduction: Interfaces in Discrete Models in 2 D
- 2 Definition of SLE
- 3 SLE/CFT correspondence
- 4 Off critical SLE
- 5 Conclusions & open problems

# Introduction

# Percolation



Each site is colored:

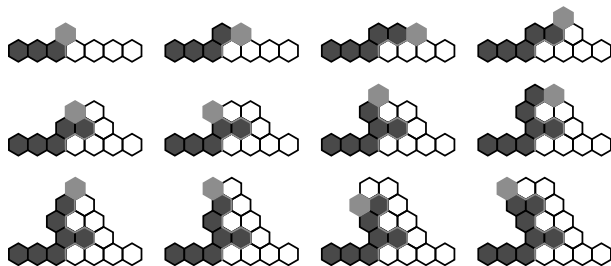
**White** with probability  $p$

**Black** with probability  $1 - p$

The model is critical on the honeycomb lattice at  $p = 1/2$

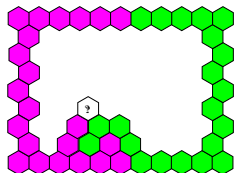
- We want to look at the properties of the interface

# Interface as a growth process



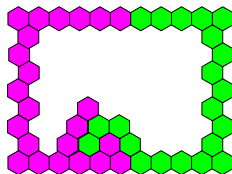
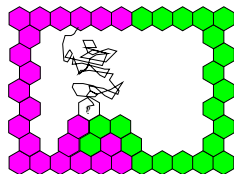
- At each step choose the color of the site with Bernoulli distribution.

# Harmonic explorer



New rule for choosing the new color.

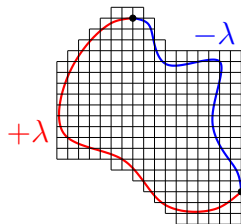
Start a random walk from “?” and stop it when it touches the boundary



# Gaussian Free Field

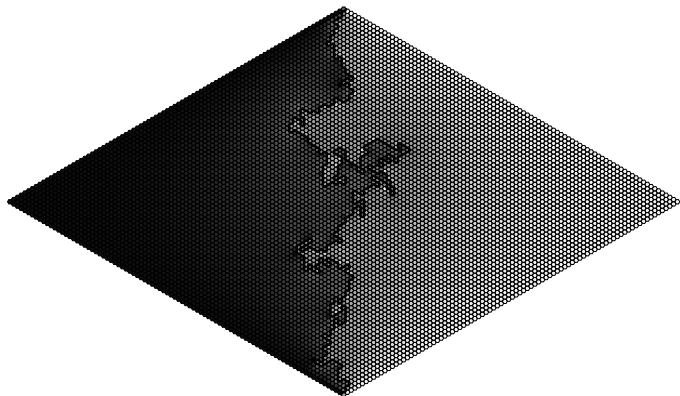
- On each site  $i$  of the exagonal lattice there is a variable  $h(i)$
- Boltzmann weight  $W(\{h(i)\}) = \exp(-\frac{1}{2} \sum_{\langle i,j \rangle} (h(i) - h(j))^2)$ .

Boundary conditions  $\pm\lambda$



- Consider an affinization and look at the Zero Level Line

# Gaussian Free Field

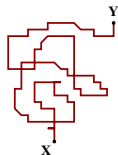


For  $\lambda = \sqrt{2}$  the continuum limit of Harmonic explorer and GFF Zero level line are statistically the same SLE<sub>4</sub> [Schramm, Sheffield '03].

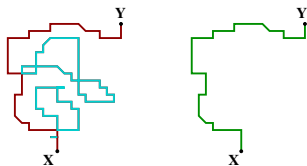


# Loop Erased Random Walks [Lawler]

Let us consider a Random Walk  $w$  from  $X$  to  $Y$



Erase the loops in an ordered way; call  $\mathcal{L}(w)$  the **Loop Erasure** of  $w$



Boltzmann weight of a simple walk

$$\omega_{\mathcal{L}}(\gamma) = \sum_{w|\mathcal{L}(w)=\gamma} \omega(w) \quad \omega(w) = \mu^{|w|}$$

Relations with **Uniform Spanning Trees** [Pemantle '91, Wilson '96]  
Continuum limit SLE<sub>2</sub>[Schramm '99]

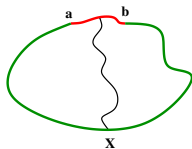
# A few questions we can ask?

- Left passage

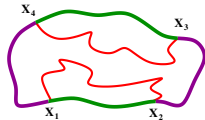
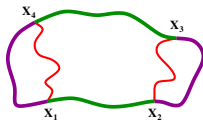


- Fractal dimension  $\sim$  passage through a ball.

- Boundary hitting probabilities



- Crossing formulae
- Topology of interfaces



# Definition of SLE

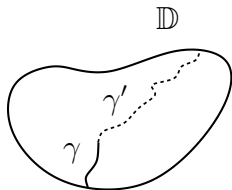
# Two key properties of these models

- Domain Markov Property (often valid already in the discrete model)
- Conformal Invariance (valid in the continuum limit)

# Domain Markov Property

Suppose we know part of our interface: **what is the law of the remaining of the curve?**

The conditioned law is the same as the one in the cut domain



$$P_{\mathbb{D}}[\gamma' | \gamma] = P_{\mathbb{D} \setminus \gamma}[\gamma]$$

# Conformal Invariance: preliminaries

*Conformal invariance is a statement about transport of probabilities under conformal mappings.*

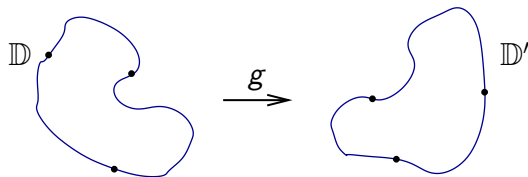
# Conformal Invariance: preliminaries

*Conformal invariance is a statement about transport of probabilities under conformal mappings.*

Consider two simply connected domains  $\mathbb{D}$  and  $\mathbb{D}'$ :

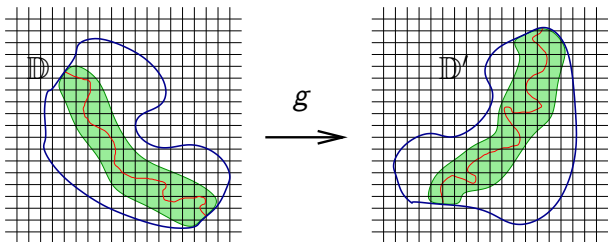
**Riemann Mapping Theorem:**

$\exists!$  a conformal map between  $g(z)$  from  $\mathbb{D}$  to  $\mathbb{D}'$  mapping the three marked points.



# Conformal Invariance

The underlying lattice is **NOT** transformed under the mapping  $g(z)$



The statement of Conformal Invariance

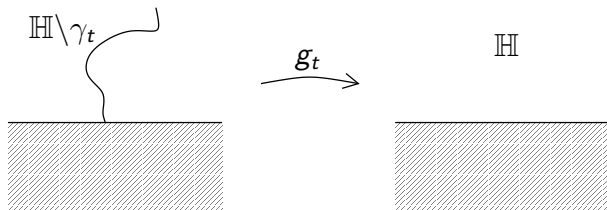
$$P_{\mathbb{D}}[\gamma] = P_{g(\mathbb{D})=\mathbb{D}'}[g(\gamma)]$$



# Löwner equation

Idea: code the shape of the curve by a uniformizing map.

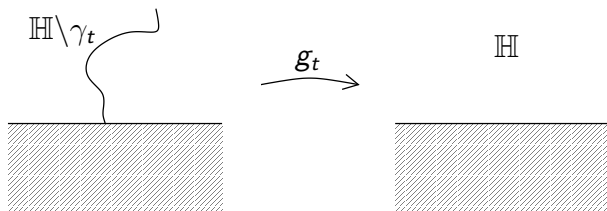
$$\text{Chordal Löwner equation } \mathbb{H} \setminus \gamma_t \rightarrow \mathbb{H}$$



# Löwner equation

Idea: code the shape of the curve by a uniformizing map.

*Chordal Löwner equation*  $\mathbb{H} \setminus \gamma_t \rightarrow \mathbb{H}$



- Hydrodynamical normalization + time parametrization:

$$g_t(z) = z + \frac{2t}{z} + O(1/z^2)$$

# Löwner equation

$g_t(z)$  satisfies the equation

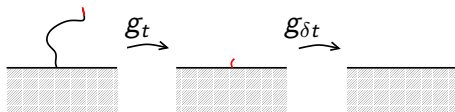
$$\dot{g}_t(z) = \frac{2}{g_t(z) - \xi_t}$$

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$$\dot{g}_t(z) = \frac{2}{g_t(z) - \xi_t}$$

Sketch of the derivation: composition of infinitesimal maps



$$g_{\delta t}(z) = \sqrt{z^2 + 4\delta t} + O(\delta t^2)$$

$$g_{t+\delta t}(z) = g_{\delta t}(g_t(z) - \xi_t) + \xi_t$$

# Schramm's idea [Schramm '99]

- If the curve  $\gamma$  is random then its measure induces a probability measure on the driving function  $\xi_t$ .

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Domain Markov Property + Conformal invariance

The driving function is proportional to a Brownian Motion

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Domain Markov Property + Conformal invariance

The driving function is proportional to a Brownian Motion

$$\xi_t = \sqrt{\kappa} B_t$$

We are left with a single parameter  $\kappa$

$\kappa = 6$ : percolation;  $\kappa = 4$ : Harmonic explorer;  $\kappa = 2$ : LERW;  
 $\kappa = 8$ : UST;  $\kappa = 3$ : Ising (spin cluster);  $\kappa = 8/3$  SAW.

# Some results

- It is always possible to define a continuous curve  $\gamma_t = \lim_{z \rightarrow 0} g_t^{-1}(z + \xi_t)$  [Rhode & Schramm]
- There are three phases [RS]
  - $\kappa \leq 4$ : the curve  $\gamma_t$  is simple.
  - $4 < \kappa < 8$ : the curve  $\gamma_t$  can touch itself.
  - $\kappa \geq 8$ : the curve  $\gamma_t$  is space filling.
- The fractal dimension of  $\gamma_t$  is  $\min\{1 + \kappa/8, 2\}$  [Beffara].
- Proof of Mandelbrot's conjecture. The boundary of 5 independent Brownian Excursions is the same as the boundary of 8 independent SLE $_{8/3}$  starting from the same point, its dimension is  $4/3$ .

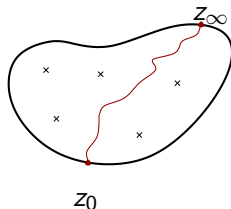


# SLE/CFT correspondence

# SLE/CFT correspondence I [Bernard, Bauer]

In Boundary Conformal Field theory

$$\langle \mathcal{O} \rangle_{\mathbb{D}} = \frac{\langle \mathcal{O} \text{ BCOs} \rangle_{\mathbb{D}}}{\langle \text{BCOs} \rangle_{\mathbb{D}}} = \frac{\langle \mathcal{O} \psi(z_0) \psi(z_\infty) \rangle_{\mathbb{D}}}{\langle \psi(z_0) \psi(z_\infty) \rangle_{\mathbb{D}}}$$



The  $\psi$ s are boundary **Primary Fields**

$$\langle \mathcal{O} \psi(z_0) \psi(z_\infty) \rangle_{\mathbb{D}} = |g'(z_0)|^h |g'(z_\infty)|^h \langle {}^g \mathcal{O} \psi(g(z_0)) \psi(g(z_\infty)) \rangle_{g(\mathbb{D})}$$

The **Martingale Condition**

$$\langle \mathcal{O} \rangle_{\mathbb{D} \setminus \gamma_t} = \frac{\langle {}^{g_t} \mathcal{O} \psi(g_t(z_0)) \psi(g_t(z_\infty)) \rangle_{\mathbb{D}}}{\langle \psi(g_t(z_0)) \psi(g_t(z_\infty)) \rangle_{\mathbb{D}}} = \text{It is a (local) martingale}$$

imposes conditions on the fields  $\psi$

# SLE/CFT correspondence II

- The conformal weight of  $\psi$  is related to  $\kappa$ :

$$h_{1,2} = \frac{6 - \kappa}{2\kappa}$$

- The correlation functions satisfy a differential equation which says that they must be **degenerate at level  $-2$**

This fixes the value of the central charge

$$c = \frac{(6 - \kappa)(3\kappa - 8)}{2\kappa}$$

consistent with the previous identification of the statistical models:  
 Percolation & SAW  $c = 0$ , GFF  $c = 1$ , Ising  $c = 1/2$ , LERW & UST  
 $c = -2$ .

# Off critical SLE

# Naive considerations I

- We want to show how the off-critical measure on random curves is related to the critical one.
- One can still use the Löwner map: how does the the law of the driving function get modified?
- At scales smaller than the correlation length (but large w.r.t. the lattice spacing) the curve should look like a critical SLE

$$" \lim_{\lambda \rightarrow 0^+} \frac{1}{\lambda} (\xi_{s+\lambda^2 t} - \xi_s) \sim \sqrt{\kappa} B_t "$$

- Does something even stronger hold?

$$d\xi_t = \sqrt{\kappa} B_t + F_t dt \quad ?$$

## Naive considerations II

The probability of a piece of interface in a discrete model out of criticality is

$$\mathbb{P}[\gamma_t] = \frac{\sum_{\rho|\gamma_t} w_{\rho}}{\sum_{\rho} w_{\rho}} = \frac{Z_{\mathbb{D}}[\gamma_t]}{Z_D} = M_t \mathbb{P}^{(c)}[\gamma_t]$$

with

$$M_t = \frac{Z_{\mathbb{D}}[\gamma_t]/Z_{\mathbb{D}}^c[\gamma_t]}{Z_D/Z_D^c}$$

By construction this is a martingale (= conserved in average).

Assume for the moment that it remains well defined in the continuum limit (need some precision).

# Girsanov's theorem

If  $B_t$  is a **Brownian Motion** w.r.t.  $\mathbb{P}^c$ , then how does it get modified by  $\mathbb{P} = M_t \mathbb{P}^c$ ?

It gets a **drift term**

$$dB_t = d\tilde{B}_t + F_t dt$$

$\tilde{B}_t$  is Brownian Motion w.r.t. the off critical measure  $\mathbb{P}$ , while

$$M_t^{-1} dM_t = F_t dB_t$$

We shall apply this to the driving function of our favourite SLE, which is proportional to a Brownian Motion.

# How to compute $M_t$ and drift: field theory I

$$M_t \sim \frac{Z_{\mathbb{D}}[\gamma_t]}{Z_{\mathbb{D}}^c[\gamma_t]} = e^{E_{\mathbb{D}}[\gamma_t] - E_{\mathbb{D}}^c[\gamma_t]} \frac{Z_{\mathbb{D} \setminus \gamma_t}}{Z_{\mathbb{D} \setminus \gamma_t}^c}$$

R.G. arguments tells us when  $E_{\mathbb{D}}[\gamma_t]$  &  $E_{\mathbb{D}}^c[\gamma_t]$  are irrelevant: this holds for LERW and GFF.

Use field theory to compute  $\frac{Z_{\mathbb{D} \setminus \gamma_t}}{Z_{\mathbb{D} \setminus \gamma_t}^c}$

- SLE<sub>4</sub> Massive free field

$$S_{FF}(\varphi) = \int \frac{dz^2}{8\pi} \partial\varphi \bar{\partial}\varphi + m^2 \varphi^2$$

- SLE<sub>2</sub> Massive symplectic fermions

$$S_{sf}(\chi^{\pm}) = \int dz^2 (4\partial\chi^+ \bar{\partial}\chi^- + m^2 \chi^+ \chi^-)$$



# How to compute $M_t$ and drift II: gaussian free field

$$M_t = \frac{\int [D\varphi]_{DBC} e^{\int \frac{dz^2}{8\pi} \partial\varphi \bar{\partial}\varphi + m^2 \varphi^2}}{\int [D\varphi]_{DBC} e^{\int \frac{dz^2}{8\pi} \partial\varphi \bar{\partial}\varphi}} = e^{-\int \frac{dz^2}{8\pi} m^2 \phi_t \Phi_t^{[m]}} \left[ \frac{\det[-\Delta + m^2]_{\mathbb{H}_t}}{\det[-\Delta]_{\mathbb{H}_t}} \right]^{-1/2}$$

$$\Delta\phi_t = 0, \quad \Delta\Phi_t - m^2\Phi_t = 0$$

plus discontinuous Dirichlet BCs in  $\mathbb{H}_t$ .

The rigorous definition of the Laplacian determinant is through  $\zeta$ -function regularization.

**Task:** Prove that this is a martingale.

The tricky part is to compute the derivative of the determinants ratio. Recast it in such a way one can use Hadamard formula.

# How to compute $M_t$ and drift II: gaussian free field

The result for the **drift** is

$$F_t = -2\sqrt{2} \int \frac{dz^2}{4\pi} m^2 \Phi_t(z) \theta_t(z)$$

$\theta_t(z) = -\Im m \frac{2}{g_t(z) - \xi_t}$  is related to the Poisson kernel in  $\mathbb{H}_t$

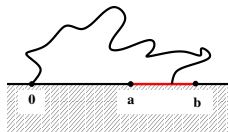
Byproduct:

- Decomposition of the field

$$\mathbb{E}[\langle e^{J_+ \star \varphi} \rangle_{\mathbb{H}_+} \langle e^{J_- \star \varphi} \rangle_{\mathbb{H}_-}] = \langle e^{J \star \varphi} \rangle_{\mathbb{H}}$$

# Drift for LERW

We consider so called **Dipolar SLE** on  $\mathbb{H}$



- Critical driving function

$$d\xi_t^c = \sqrt{2}dB_t + F_{t,[a,b]}^c dt$$

- Girsanov's martingale

$$M_t = \left[ \frac{\det[-\Delta + m^2]_{\mathbb{H}_t}}{\det[-\Delta]_{\mathbb{H}_t}} \right] \frac{\langle \psi^+(\gamma_t) \int_a^b dx \psi^-(x) \rangle_{\mathbb{H}_t}^m}{\langle \psi^+(\gamma_t) \int_a^b dx \psi^-(x) \rangle_{\mathbb{H}_t}^{m=0}}$$

$$\psi^\pm(x) = \lim_{\delta \rightarrow 0} \delta^{-1} \chi^\pm(x + i\delta)$$

- Off critical drift

$$F_{t,[a,b]} = 2\partial_{\xi_t} \log \langle \psi^+(\gamma_t) \int_a^b dx \psi^-(x) \rangle_{\mathbb{H}_t}^m$$

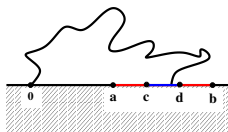
# An application: boundary hitting probability

- Consistency check: ratios of correlation functions of symplectic fermions must be martingales

$$\frac{\langle \psi^+(\gamma_t) \prod_{j=1}^{N+1} \chi^-(z_j) \prod_{j=1}^N \chi^N(z_j) \rangle}{\langle \psi^+(\gamma_t) \int_a^b dx \psi^-(x) \rangle_{\mathbb{H}_t}^m}$$

- Boundary Hitting Probability

$$P_{[a,b]}^{[c,d]} = \frac{\langle \psi^+(\gamma_t) \int_c^d dx \psi^-(x) \rangle_{\mathbb{H}_t}^m}{\langle \psi^+(\gamma_t) \int_a^b dx \psi^-(x) \rangle_{\mathbb{H}_t}^m}$$



# Conclusions & open problems

- Conclusions

- SLE techniques to describe interfaces in 2D models & relation to CFT approach.
- Girsanov's theorem to describe SLEs out of criticality.
- Application to  $SLE_2$  and  $SLE_4$ : computation of the drift.

- Perspectives

- Computation of other probabilities for LERW (left passage)
- Off criticality for other values of  $\kappa$  and other kinds of perturbations (e.g. integrable perturbations): much more ambitious.