

Preliminaires sur les géodésiques

Soit  $S = \{ f(u,v) \in \mathbb{R}^3, u \in I, v \in J \}$  une surface paramétrée.

géodésique sur S = courbe  $\gamma: I \rightarrow S$  régulière t.q.

- ①  $\gamma$  minimise la distance entre deux points  $P, Q \in S$ .
- $\Leftrightarrow$  ②  $\gamma''(t) \perp T_{\gamma(t)} S \quad \forall t$ , i.e.  $\gamma''(t) \parallel N_{\gamma(t)}$
- $\Leftrightarrow$  ③  $\langle \gamma''(t), f_u(\gamma(t)) \rangle = 0 \quad \forall t \in I$ .
- $\langle \gamma''(t), f_v(\gamma(t)) \rangle = 0$

Equation des géodésiques: pour déterminer une géodésique

$\gamma(t) = f(u(t), v(t)) \in S \quad \forall t \in I$ , il suffit de déterminer les fonctions  $u(t)$  et  $v(t)$  qui donnent les paramètres.

$$\gamma'(t) = \frac{d}{dt} f(u(t), v(t)) = f_u u'(t) + f_v v'(t)$$

$$\gamma''(t) = \frac{d}{dt} (f_u(u'(t), v(t)) u'(t) + f_v(u(t), v(t)) v'(t))$$

$$= f_{uu} (u')^2 + f_{uv} u'v' + f_u u'' + f_{vu} u'v' + f_{vv} (v')^2 + f_v v'' \quad \text{donc}$$

$$\boxed{\gamma''(t) = f_u u'' + f_v v'' + f_{uu} (u')^2 + 2 f_{uv} u'v' + f_{vv} (v')^2}$$

On définit les symboles de Christoffel  $\Gamma_{jk}^i$ ,  $i, j, k = 1, 2$

en décomposant  $f_{uu}, f_{uv}, f_{vv}$  dans la base  $\{ f_u, f_v, N \}$  de  $\mathbb{R}^3$ :

$$\begin{cases} f_{uu} = \Gamma_{11}^1 f_u + \Gamma_{11}^2 f_v + l N \\ f_{uv} = \Gamma_{12}^1 f_u + \Gamma_{12}^2 f_v + m N \\ f_{vv} = \Gamma_{22}^1 f_u + \Gamma_{22}^2 f_v + n N \end{cases}$$

Les  $\Gamma_{jk}^i$  se trouvent avec les systèmes:

$$\begin{cases} \langle f_{uu}, f_u \rangle = \Gamma_{11}^1 E + \Gamma_{11}^2 F \\ \langle f_{uv}, f_u \rangle = \Gamma_{12}^1 F + \Gamma_{12}^2 G \end{cases} \quad \begin{cases} E_u = \frac{\partial}{\partial u} \langle f_u, f_u \rangle = 2 \langle f_{uu}, f_u \rangle \\ E_v = \frac{\partial}{\partial v} \langle f_u, f_u \rangle = 2 \langle f_{uv}, f_u \rangle \end{cases}$$

$$\begin{cases} \langle f_{uv}, f_u \rangle = \Gamma_{12}^1 E + \Gamma_{12}^2 F \\ \langle f_{vv}, f_u \rangle = \Gamma_{22}^1 F + \Gamma_{22}^2 G \end{cases} \quad \begin{cases} F_u = \frac{\partial}{\partial u} \langle f_u, f_v \rangle = \langle f_{uu}, f_v \rangle + \langle f_{uv}, f_u \rangle \\ F_v = \frac{\partial}{\partial v} \langle f_u, f_v \rangle = \langle f_{uv}, f_v \rangle + \langle f_{vv}, f_u \rangle \end{cases}$$

$$\begin{cases} \langle f_{uv}, f_v \rangle = \Gamma_{12}^1 E + \Gamma_{12}^2 F \\ \langle f_{vv}, f_v \rangle = \Gamma_{22}^1 F + \Gamma_{22}^2 G \end{cases} \quad \begin{cases} G_u = \frac{\partial}{\partial u} \langle f_v, f_v \rangle = 2 \langle f_{uv}, f_v \rangle \\ G_v = \frac{\partial}{\partial v} \langle f_v, f_v \rangle = 2 \langle f_{vv}, f_v \rangle \end{cases}$$

Si on combine les deux systèmes précédents, on trouve :

$$\begin{cases} \Gamma_{11}^1 E + \Gamma_{11}^2 F = \frac{1}{2} E_u \\ \Gamma_{11}^1 F + \Gamma_{11}^2 G = F_u - \frac{1}{2} E_v \end{cases}$$

$$\begin{cases} \Gamma_{12}^1 E + \Gamma_{12}^2 F = \frac{1}{2} E_v \\ \Gamma_{12}^1 F + \Gamma_{12}^2 G = \frac{1}{2} G_u \end{cases}$$

$$\begin{cases} \Gamma_{22}^1 E + \Gamma_{22}^2 F = F_v - \frac{1}{2} G_u \\ \Gamma_{22}^1 F + \Gamma_{22}^2 G = \frac{1}{2} G_v \end{cases}$$

$$\Leftrightarrow \begin{cases} \Gamma_{11}^1 = \frac{GE_u - F(2F_u - E_v)}{2(EG - F^2)} \\ \Gamma_{11}^2 = \frac{E(2F_u - E_v) - FE_u}{2(EG - F^2)} \end{cases}$$

$$\Leftrightarrow \begin{cases} \Gamma_{12}^1 = \frac{GE_v - FG_u}{2(EG - F^2)} \\ \Gamma_{12}^2 = \frac{EG_u - FE_v}{2(EG - F^2)} \end{cases}$$

$$\Leftrightarrow \begin{cases} \Gamma_{22}^1 = \frac{G(2F_v - G_u) - FG_v}{2(EG - F^2)} \\ \Gamma_{22}^2 = \frac{EG_v - F(2F_v - G_u)}{2(EG - F^2)} \end{cases}$$

La condition ③ sur le pédoncule  $\gamma(t)$  devient donc :

$$\begin{cases} -u'' = \Gamma_{11}^1 (u')^2 + 2\Gamma_{12}^1 u'v' + \Gamma_{22}^1 (v')^2 \\ -v'' = \Gamma_{11}^2 (u')^2 + 2\Gamma_{12}^2 u'v' + \Gamma_{22}^2 (v')^2 \end{cases}$$