

From gauge fields to direct connections on gauge groupoids

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Women at the intersection of Mathematics and Theoretical Physics

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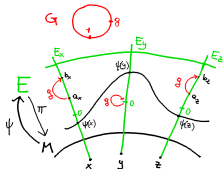
Based on a joint work with

Sara Azzali (Hamburg), **Youness Boutaïb** (Aachen) and **Sylvie Paycha** (Potsdam)

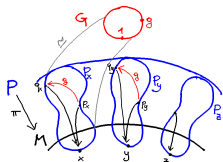
in progress also with **Alfonso Garmendia** (Potsdam)

Geometric model underlying field and gauge theories

- **space-time manifold** M
- **vector/spinor bundle** $E \rightarrow M$ with fibres V
- **(matter) field** $\psi : M \rightarrow E$ section of E
 \Rightarrow configuration space $\mathcal{E} = \Gamma(M, E)$
- **dynamics** via action $S[\psi]$ (isolated particles)
 \Rightarrow use **derivatives** $D_X \psi$ i.e. **linear connection on E**



- **symmetries** by Lie group G acting on fibres of E
- **principal G -bundle** $P \rightarrow M$ s.t. $E = P \times_G V$
 (associated bundle)
- **gauge bosons (force carriers)** $A \in \Omega_{loc}^1(M, \mathfrak{g})$
 local form of **principal connection on P**
- **gauge group** $\hat{G} = \text{Aut}_M(P) = \Gamma(M, P^{ad}) = \text{Aut}_M(E)$
 \Rightarrow acts on ψ and A
- **dynamics** via \hat{G} -inv. action $S[\psi, A]$ (interacting particles)
 \Rightarrow new **covariant derivative** $D_X^A \psi$



Idea: replace $\left\{ \begin{array}{l} \text{connection on } P \\ + \text{ gauge group} \\ (+ \text{ Lie algebra } \mathfrak{g}) \end{array} \right.$ by $\left\{ \begin{array}{l} \text{direct connection} \\ \text{on gauge groupoid of } P \\ (+ \text{ Lie algebroid of } P) \end{array} \right.$.

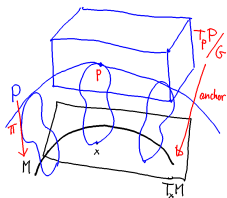
Infinitesimal structure: Lie algebroids

- **Lie algebroid:** vector bundle $A \rightarrow M$ with a Lie bracket $[\cdot, \cdot]_A$ on sections $\Gamma(A)$ and an anchor map $a : A \rightarrow TM$ inducing a derivation on sections w.r.t. vector fields on M :

$$[X, fY]_A = f[X, Y]_A + a(X)(f)Y$$

- **Atiyah Lie algebroid of a principal G -bundle $\pi : P \rightarrow M$:**

$A(P) = TP/G \rightarrow M$ with fibres $A_x(P) \cong T_{p_x}P$,
 anchor $A(P) \rightarrow TM$ induced by $d\pi : TP \rightarrow TM$
 via quotient map $j : TP \rightarrow TP/G$
 and Lie bracket of G -invariant vector fields on P .



- **Trivial G -bundle:** $P = M \times G \rightarrow M \Rightarrow A(M \times G) = TM \oplus (M \times \mathfrak{g}) \xrightarrow{id+0} TM$
 where $\mathfrak{g} = \text{Lie}(G)$.
- **Frame bundle of a vector bundle $E \rightarrow M$ of rank r :** $F(E) = \bigcup_{x \in M} \text{Iso}(\mathbb{R}^r, E_x) \rightarrow M$
 $\Rightarrow A(F(E)) = \text{Der}(E) \rightarrow TM$ bundle of derivative endomorphisms
 s.t. $\Gamma(\text{Der}(E)) = \text{derivations of } \Gamma(E)$.

Principal connections, gauge fields and covariant derivative

- **Principal connection on P :** five equivalent presentations
 - 1) G -equivariant **horizontal subbundle** $HP \subset TP \rightarrow P$ s.t. $TP = HP \oplus VP$, where the **vertical bundle** VP (spaces tangent to the fibres) is canonical.
 - 2) G -equivariant **connection 1-form** $\omega \in \Omega^1(P, \mathfrak{g})$ s.t. $\omega_p(\hat{A}_p) = A$ if \hat{A} is the (vertical) fundamental vector field on P det. by $A \in \mathfrak{g}$, and $\omega_p(B_p) = 0$ if $B_p \in H_p P$ is horizontal.
 - 3) **infinitesimal connection** $\delta: TM \rightarrow A(P)$ s.t. $j^{-1}(\delta(X)) \in HP$ is the horizontal lift of $X \in TM$.
 - 4) **parallel transport** $\tau_\gamma(y, x): P_x \xrightarrow{\cong} P_y$ along a curve γ of M from x to y given by the horizontal lift of γ .
 - 5) (local) **gauge fields** $\{A_\alpha \in \Omega^1(U_\alpha, \mathfrak{g})\}$ = pull back of ω along local sections of P .
- **Covariant derivative on sections of E :** bundle map $D^A: TM \rightarrow \text{Der}(E)$ equivalent to $C^\infty(M)$ -derivation $D^A_X: \Gamma(E) \rightarrow \Gamma(E)$ for $X \in \Gamma(TM)$ given only locally by

$$D^A_X(\psi)|_{U_\alpha} = \sum_{\mu, i, j} (X^\mu \partial_\mu \psi^i + X^\mu A_{\mu j}^i \psi^j) e_i$$

if $X = X^\mu \partial_\mu$ in coordinates x^μ on $U_\alpha \subset M$

$\psi = \psi^i e_i$ on a local basis (e^i) of E_{U_α}

and $A_j^i = A_{\mu j}^i dx^\mu$ are the components of the **gauge field A** in terms of generators of \mathfrak{g} .

Lie groupoids

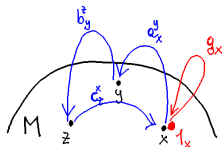
- **Lie groupoid** $\boxed{\mathcal{G} \rightrightarrows M}$: bi-fibred manifold $\mathcal{G} = \bigcup_{(y,x) \in M \times M} \mathcal{G}_x^y$

with elements $a_{yx} \in \mathcal{G}_x^y$ called **arrows**

projections $s, t : \mathcal{G} \rightarrow M$ $\left\{ \begin{array}{l} \text{source } s(a_{yx}) = x \\ \text{target } t(a_{yx}) = y \end{array} \right.$

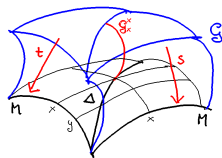
such that

- arrows can be **composed**: $b_{zy} a_{yx} \in \mathcal{G}_x^z$ if $s(b_{zy}) = t(a_{yx})$ (composition is associative),
- there are **units** $u(x) = 1_x \in \mathcal{G}_x^x$ and $M \equiv u(M) \subset \mathcal{G}$,
- each arrow $a_{yx} \in \mathcal{G}_x^y$ has an **inverse** $a_{yx}^{-1} \in \mathcal{G}_y^x$.



The induced map $(t, s) : \mathcal{G} \rightarrow M \times M$ is called the **anchor**.

Each \mathcal{G}_x^x is a Lie group, called the **vertex group** or **isotropy**.



- **Infinitesimal structure of Lie groupoid = Lie algebroid:**

$$\boxed{\mathcal{L}\mathcal{G} = \bigcup_{x \in M} T_{1_x} \mathcal{G}_x \rightarrow TM}$$

Gauge groupoids

- **Gauge groupoid of principal G -bdl $P \rightarrow M$:**

$$\mathcal{G}(P) = P \times_G P \rightrightarrows M$$

contains equivalence classes $[p, q]$ under $(p, q) \sim (pg, qg)$ for $g \in G$.

$$\Rightarrow 1) \quad \mathcal{L}(\mathcal{G}(P)) = A(P) \rightarrow TM$$

- 2) $\mathcal{G}(P)$ acts on $E = P \times_G V$:

$$\mathcal{G}(P) \times_M E = \{([p, q], [r, v]), \pi(q) = \pi(r)\} \rightarrow E$$

with action $\rho_{[p, q]}([r, v]) = [p, gv]$ where $g \in G$ s.t. $r = qg$,

- 3) $\hat{G} = \text{Aut}_M(P) \subset \mathcal{G}(P)$ given by $\Phi \mapsto [\Phi(p), p]$ for any $p \in P$.

- **Pair groupoid:**

$$\text{Pair}(M) = M \times M \rightrightarrows M$$

$$\text{for } P = M \times \{1\} \rightarrow M$$

$$\Rightarrow \text{Lie algebroid} = \mathcal{L}(\text{Pair}(M)) = TM \xrightarrow{id} TM.$$

- **Trivial Lie groupoid with fibre G :**

$$M \times G \times M \rightrightarrows M$$

$$\text{for } P = M \times G \rightarrow M$$

$$\Rightarrow \text{Lie algebroid} = TM \oplus (M \times \mathfrak{g}) \xrightarrow{id+0} TM.$$

- **Frame groupoid of $E \rightarrow M$:**

$$\text{Iso}(E) = \bigcup_{x, y} \text{Iso}(E_y, E_x)$$

$$\text{for } P = F(E)$$

$$\Rightarrow \text{Lie algebroid} = \text{Der}(E) \rightarrow TM$$

If the structure gp of E reduces to $G \subset GL_r(\mathbb{R})$ and $P \subset F(E)$ then

$$\mathcal{G}(P) \hookrightarrow \text{Iso}(E)$$

Direct connections on Lie groupoids

- **Local map** $\psi : \mathcal{G} \rightarrow \mathcal{G}'$ between two groupoids: map $\psi : \mathcal{U} \subset \mathcal{G} \rightarrow \mathcal{G}'$ defined on an open neighborhood \mathcal{U} of the units $u(M) \subset \mathcal{G}$, which commutes with s , t and u .
Local morphism: local map which also preserves composition (hence inversion).

- [Teleman 2004 in the linear case, Kock 2007 similar, ABFP general]

Direct connection on $\mathcal{G} \rightrightarrows M$: **local right inverse of the anchor which preserves units**, i.e. $\Gamma : \text{Pair}(M) \rightarrow \mathcal{G}$ defined on an open n . \mathcal{U}_Δ of the diagonal $\Delta \subset \text{Pair}(M)$ s.t.

$$\Gamma(y, x) \in \mathcal{G}_x^y \text{ for all } (y, x) \in \mathcal{U}_\Delta \quad \text{and} \quad \Gamma(x, x) = 1_x \in \mathcal{G}_x^x \text{ for all } x \in M.$$

- A Lie groupoid with a direct connection is a **gauge groupoid**.
- If $\mathcal{G} \times_M E \rightarrow E$ is a linear action, then a direct connection Γ on \mathcal{G} induces a **transport on fibres** $E_x \rightarrow E_y$ which is **not necessarily a parallel displacement!**
- Γ **natural** if $\Gamma(x, y) \Gamma(y, x) = 1_x$ for all $x \in M$ and suitable y .
- **Curvature of Γ at x** : $R^\Gamma(z, y, x) = \Gamma(z, x)^{-1} \Gamma(z, y) \Gamma(y, x) \in \mathcal{G}_x^x$ for suitable y, z .
 Γ is **flat** if $R^\Gamma(-, -, x) = 1_x$ for any x , i.e. Γ is a **groupoid morphism**.

Relationship to usual connections

Assume M is a manifold with affine connection ∇^M and local geodesics.

- **Parallel displacement** τ on $P \rightarrow M$ along small geodesics (equivalent to a principal connection ω on P hence to gauge fields A) defines a **direct connection** Γ^τ on $\mathcal{G}(P) \rightrightarrows M$ by

$$\Gamma^\tau(y, x) = [\tau(y, x)(p), p] \quad \text{for any choice of } p \in P_x$$

Same for $E \rightarrow M$ and $\text{Iso}(E)$ [Teleman 2004].

- Viceversa, a **direct connection** Γ on $\mathcal{G}(P) \rightrightarrows M$ induces an infinitesimal connection on the Lie algebroid $A(P) \rightarrow TM$ by

$$\nabla^\Gamma(\dot{\gamma}(0)) = D\Gamma|_M(\dot{\gamma}(0)) = \frac{d}{dt}|_{t=0} \Gamma(\gamma(t), x),$$

hence a **principal connection** ω^Γ on P .

- Apply maps $\omega \mapsto \tau \mapsto \Gamma^\tau \mapsto \omega^{\Gamma^\tau}$, then $\boxed{\omega^{\Gamma^\tau} = \omega}$ on P .
- Viceversa, if apply maps $\Gamma \mapsto \omega^\Gamma \mapsto \tau^\Gamma \mapsto \Gamma^{\tau^\Gamma}$, then $\boxed{\Gamma^{\tau^\Gamma} \neq \Gamma}$ on $\mathcal{G}(P)$ in general. There are direct connections on $\mathcal{G}(P)$ which are **not parallel displacements!**

Examples

- $M = \mathbb{R}$ with flat connection $\nabla_{\partial_x}^M (h(x) \partial_x) = h'(x) \partial_x$.
- $E = M \times \mathbb{R} \rightarrow M$ with global section $e_1(x) = (x, 1) \in E_x$ and linear connection $\nabla_{\partial_x}^E : \Gamma(E) \rightarrow \Gamma(E)$ given by $f \in C^\infty(M)$ s.t. $\nabla_{\partial_x}^E e_1 = f e_1$.
- The induced **parallel transport** along a geodesic from x to y is the isomorphism $\tau(y, x) : E_x \rightarrow E_y$ defined by $\tau(y, x) \xi_0 e_1(x) = \xi(y) e_1(y)$ solution of the ODE

$$\nabla_{\partial_x}^E (\xi(x) e_1(x)) = (\xi'(x) + \xi(x)f(x))e_1(x) = 0$$

with initial value $\xi(x) e_1(x) = \xi_0 e_1(x)$. Set $F(x) = \int -f(x) dx$.
Then the direct connection on $\text{Iso}(E)$ is

$$\tau(y, x) : E_x \rightarrow E_y, \quad e_1(x) \mapsto \tau(y, x) e_1(x) = e^{F(y)-F(x)} e_1(y)$$

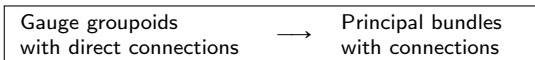
The associated direct connection is flat. For instance:

$$\nabla_{\partial_x}^E e_1(x) = -2x e_1(x) \Rightarrow \tau(y, x) e_1(x) = e^{y-x+y^2-x^2} e_1(y),$$

$$\nabla_{\partial_x}^E e_1(x) = -3x^2 e_1(x) \Rightarrow \tau(y, x) e_1(x) = e^{y-x+y^3-x^3} e_1(y).$$

- Instead, the following direct connections are **not parallel transports**:
 $\Gamma(y, x) = e^{y-x+(y-x)^2}$ non natural ($\Gamma(x, y)\Gamma(y, x) = e^{2(y-x)^2} \neq 1_x$),
 $\Gamma(y, x) = e^{y-x+(y-x)^3}$ natural but non-flat.

Conclusion: there is a surjective functor



which admits an inverse, but it is not an equivalence of categories.

Further results:

- Jet prolongation of direct connections to jet groupoids $J^n\mathcal{G} \rightrightarrows M$ (existence, examples).
- Applications to geometric regularity structures for solving stochastic PDEs (cf. M. Hairer and coll.).

Next:

- Look for more examples of direct connections which are not parallel displacements.
- Adapt to α -Hölder sections of bundles i.e. define distributional direct connections and compare to usual propagators.
- Study the whole geometry of groupoids with direct connections and compare with usual gauge theory.

Thank you for the attention!