A Lévy-Fokker-Planck equation: entropies and convergence to equilibrium

I. Gentil

CEREMADE, Université Paris-Dauphine

International Conference on stochastic Analysis and Applications Hammamet, Tunisia, November 5-10, 2007

Work written in collaboration with C. Imbert from Paris-Dauphine

<ロ><日><日><日><日><日><日><日><日><日><日><日><日><10</td>



Introduction

- Ornstein-Uhlenbeck and Fokker-Planck equations
- Tools for the asymptotic behaviour
- 2 The Lévy-Fokker-Planck equation
 - The Lévy-Fokker-Planck equation
- 3 Results
 - The equilibrium
- 4 Results
 - Entropies
 - Convergence towards the equilibrium



Introduction

- Ornstein-Uhlenbeck and Fokker-Planck equations
- Tools for the asymptotic behaviour
- 2 The Lévy-Fokker-Planck equation
 - The Lévy-Fokker-Planck equation
- 3 Results
 - The equilibrium
- 4 Results
 - Entropies
 - Convergence towards the equilibrium

Ornstein-Uhlenbeck and Fokker-Planck equations

$$\begin{cases} dX_t = 2dB_t - \nabla V(X_t)dt \\ X_0 = x \end{cases}$$

where B_t is a standard Brownian Motion in \mathbb{R}^n .

Ito formula implies : the Semigroup $P_t f(x) = E_x(f(X_t))$ satisfies the PDE

$$\frac{\partial}{\partial t} P_t f(\mathbf{x}) = L P_t f(\mathbf{x})$$
$$P_0 f = f,$$

where $Lf = \Delta f - \nabla V \cdot \nabla f$ is the IG of P_t . This is the **Ornstein-Uhlenbeck equation**.

Consider L^* or P_t^* , the dual with respect to dx,

$$\int \textit{Lfgdx} = \int \textit{fL}^*\textit{gdx}, \ \, \text{or} \ \ \, \int P_t fgdx = \int fP_t^*gdx,$$

then

$$L^*g = \Delta g + \operatorname{div}(g.\nabla V).$$

The Semigroup $P_t^* f(x)$ satisfies the PDE

$$\begin{bmatrix} \frac{\partial}{\partial t} P_t^* f(x) = L^* P_t^* f(x) \\ P_0^* f = f, \end{bmatrix}$$

This is the Fokker-Planck equation.

Let $\mu_V = e^{-V} dx$ (assume that μ_V is a probability measure), $(P_t)_{t\geq 0}$ or *L* is self adjoint in $L^2(\mu_V)$ and the by integration by parts

$$\int Lf \, g d\mu_V = - \int
abla f \cdot
abla g d\mu_V.$$

Under smooth assumptions :

$$\lim_{t\to\infty} P_t f(\mathbf{x}) = \int f d\mu_V.$$

or equivalently

$$\lim_{t\to\infty} e^{V(x)} P_t^* g(x) = \int g dx.$$

The good question is HOW FAST?



Introduction

Ornstein-Uhlenbeck and Fokker-Planck equations

- Tools for the asymptotic behaviour
- The Lévy-Fokker-Planck equation
 The Lévy-Fokker-Planck equation
- 3 Results
 - The equilibrium
- 4 Results
 - Entropies
 - Convergence towards the equilibrium

Tools for the asymptotic behaviour

▶ Poincaré inequality : a L² convergence.

$$\frac{d}{dt} var_{\mu_V}(P_t f) = 2 \int P_t f L P_t f d\mu_V - 0 = -2 \int |\nabla P_t f|^2 d\mu_V,$$

If Poincaré inequality holds

$$extsf{var}_{\mu_V}(f) \leq C \int |
abla f|^2 d\mu_V$$

 $extsf{var}_{\mu_V}(\mathcal{P}_t f) \leq e^{-2t/C} extsf{var}_{\mu_V}(f).$

Logarithmic Sobolev inequality a Llog L convergence

$$\frac{d}{dt} Ent_{\mu_V}(P_t f) := \frac{d}{dt} \int P_t f \log \frac{P_t f}{\int P_t f d\mu_V} d\mu_V = -4 \int |\nabla \sqrt{P_t f}|^2 d\mu_V,$$

If Logarithmic Sobolev inequality holds

$$Ent_{\mu_{V}}(f^{2}) \leq C \int |\nabla f|^{2} d\mu_{V}$$
$$Ent_{\mu_{V}}(P_{t}f) \leq e^{-4t/C} Ent_{\mu_{V}}(f).$$

When do we have a Poincaré or a logarithmic Sobolev inequality?

* The Gaussian measure, $V(x) = x^2/2$ (Inequality proved by Gross).

* The Bakry-Emery Γ_2 -criterion implies that if

 $\operatorname{Hess}(V) \geq \lambda \operatorname{Id},$

with $\lambda > 0$ then logarithmic Sobolev inequality holds with $C = 2/\lambda$ and Poincaré inequality holds with $C = 1/\lambda$.

* There are also many technical methods to prove Poincaré or Log-Sobolev : Hardy, transportation...



Ornstein-Uhlenbeck and Fokker-Planck equations

- Tools for the asymptotic behaviour
- The Lévy-Fokker-Planck equation
 The Lévy-Fokker-Planck equation
 - 3 Results
 - The equilibrium
 - 4 Results
 - Entropies
 - Convergence towards the equilibrium

Definition of Lévy process

Lévy process L_t = process with stationary & indep increments

Fourier transform $(L_t) = e^{t\psi(\xi)}$ where ψ is characterized by the Lévy-Khinchine formula.

$$\psi(\xi) = -\sigma\xi \cdot \xi + ib \cdot \xi + \int (e^{iz \cdot \xi} - 1 - iz \cdot \xi \mathbf{1}_B(\xi))\nu(dz)$$

where ν is a singular measure satisfying

$$\int_{B} |z|^{2} \nu(dz) < +\infty \qquad \qquad \int_{\mathbb{R}^{d} \setminus B} \nu(dz) < +\infty,$$

 σ is a positive definite matrix and *b* is a vector.

Parameters (σ , b, ν) characterize the Lévy process (or a inifinite divisible law).

For all t > 0 the law of L_t is an infinite divisible law :

$$\mu = \underbrace{\mu_n \star \cdots \star \mu_n}_{n \text{ times}}.$$

► Associated infinitesimal generators as for the Brownian Motion.

$$I(u) = \operatorname{div} (\sigma \nabla u) + b \cdot \nabla u + \int (u(x+z) - u(x) - \nabla u(x) \cdot z \mathbf{1}_{B}(z)) \nu(dz)$$

These operators appear everywhere (mathematical finance, mechanics, fluids *etc.*)

- Laws with heavy tails (decrease as power laws)
- Example : $(\sigma, b, \nu) = (0, 0, \frac{1}{|z|^{\alpha+d}}dz)$, the $\alpha \in (0, 2)$ stable process. In that case $\psi(\xi) = |\xi|^{\alpha}$.
- The case $\alpha = 2$ is the Brownian motion, $I = \Delta$.

The Lévy-Fokker-Planck equation

Replace Δ by *I* a IG of a Lévy process in the Fokker-Planck equation :

$$\begin{cases} \frac{\partial}{\partial t}u = I(u) + \operatorname{div}(ux)\\ u(0, x) = f(x) \end{cases}$$
(LFP)

The goal is to understand the asymptotic behaviour of the semigroup.

<u>Remark :</u> We assume that $V = x^2/2$.

Questions :

- Find a steady state as e^{-V} as for the classical case Δ .
- Find the asymptotic behaviour of the Lévy-Fokker-Planck equation (LFP).
- Find conditions to get an asymptotic behaviour using inequalities as Poincaré or logarithmic Sobolev.



Ornstein-Uhlenbeck and Fokker-Planck equations

- Tools for the asymptotic behaviour
- 2 The Lévy-Fokker-Planck equation
 - The Lévy-Fokker-Planck equation
- 3 Results
 - The equilibrium
- 4 Results
 - Entropies
 - Convergence towards the equilibrium

An equilibrium $u_{\infty} \stackrel{\text{def}}{=}$ a stationary solution of the LFP u_{∞} can be seen as an invariant measure μ_V in the case of the Laplacian.

Proposition (Existence of an equilibrium)

Assume that

$$\int_{\mathbb{R}^d\setminus B}\ln|z|\nu(dz)<+\infty.$$

There then exists an positive equilibrium u_{∞} :

$$I(u_{\infty})+div(u_{\infty}x)=0.$$

Moreover, $u_{\infty}dx$ is an infinite divisible law whose characteristic exponent A is

$$A(\xi) = \int_0^1 \psi(s\xi) \frac{ds}{s}.$$

□ > < @ > < ⊇ > < ⊇ > < ⊇ > 14/22 The condition

$$\int_{\mathbb{R}^d\setminus B} \ln |z|\nu(dz) < +\infty.$$
 (Con1)

is satisfied for the α -stable Lévy process. In that case u_{∞} is the infinite divisible law of the Lévy process, $A = \psi/\lambda$.

Proof : The Fourier transform \hat{u}_{∞} satisfies

 $\psi(\xi)\hat{u}_{\infty}+\xi\cdot\nabla\hat{u}_{\infty}=0$

so that $\hat{u}_{\infty} = \exp(-A)$ with A such that :

 $\nabla A(\xi) \cdot \xi = \psi(\xi),$

then

$$A(\xi) = \int_0^1 \psi(s\xi) \frac{ds}{s}$$

Con1 prove that *A* is well defined and is the characteristic exponent of a Lévy process.

(ロ) (聞) (言) (言) 言 のへで 15/22



Ornstein-Uhlenbeck and Fokker-Planck equations

- Tools for the asymptotic behaviour
- 2 The Lévy-Fokker-Planck equation
 - The Lévy-Fokker-Planck equation
- 3 Results
 - The equilibrium



- Entropies
- Convergence towards the equilibrium

For $\phi : \mathbb{R}^+ \to \mathbb{R}$ convex and smooth and μ a probability measure, consider the ϕ -entropy

$$oldsymbol{E}^{\phi}_{\mu}({\it f}) = \int \phi({\it f}) {\it d} \mu - \phi\left(\int {\it f} {\it d} \mu
ight)$$

Examples

For
$$\phi(x) = \frac{1}{2}x^2$$
 (E^{ϕ}_{μ} =the variance), $D_{\phi}(a, b) = \frac{1}{2}(a-b)^2$
 $F^{\phi}_{\mu}(v) = \frac{1}{2} \iint (v(x+z) - v(x))^2 \nu(dz) \mu(dx)$
For $\phi(x) = x \ln x - x - 1$ (E^{ϕ}_{μ} =entropy), $D_{\phi}(a, b) = a \ln \frac{a}{b} + b - a$

This is a natural interpolation between the variance and the Entropy.

Define also a Bregman distance

$$oldsymbol{D}_{\phi}(a,b)=\phi(a)-\phi(b)-\phi'(b)(a-b)\geq 0$$

Theorem

Let $\mu(d\mathbf{x}) = u_{\infty}(\mathbf{x})d\mathbf{x}$, ν the Lévy measure associated to I and consider $v(t, \mathbf{x}) = \frac{u(t, \mathbf{x})}{u_{\infty}(\mathbf{x})}$, then

$$\frac{d}{dt} \frac{E^{\phi}_{\mu}(v(t,\cdot)) = -\iint D_{\phi}\bigg(v(x+z),v(x)\bigg)\nu(dz)\mu(dx).$$

Fisher information

$$\mathcal{F}^{\phi}_{\mu}(v) = \iint \mathcal{D}_{\phi}\bigg(v(x+z),v(x)\bigg)
u(dz)\mu(dx).$$

Can be seen as a Dirichlet form with respect to the measure $u_{\infty}(x)dx$

The proof is related to :

► A related equation : the Lévy-Ornstein-Ulenbeck equation (LOU)

The function $v = u/u_{\infty}$ satisfies

$$\partial_t v = \frac{1}{u_\infty} \left(I(u_\infty v) - I(u_\infty) v \right) + x \cdot \nabla v \stackrel{\text{def}}{=} L v.$$

Dual operator of *L* wrt $\mu = u_{\infty}(x)dx$

$$\int w_1(Lw_2)d\mu = \int \left(\check{I}(w_1) - x \cdot \nabla w_1\right)w_2d\mu,$$

where \check{I} is I with $\check{\nu}(dx) = \nu(-dx)$.

Recall that in the classical case L is a self-adjoint operator with respect to μ .



Ornstein-Uhlenbeck and Fokker-Planck equations

- Tools for the asymptotic behaviour
- 2 The Lévy-Fokker-Planck equation
 - The Lévy-Fokker-Planck equation
- 3 Results
 - The equilibrium



- Entropies
- Convergence towards the equilibrium

Theorem

We assume that ν_l has a density N with respect to dx and satisfies

$$\int_{\mathbb{R}^d\setminus B} \ln |z| \, N(z) \, dz < +\infty.$$

If N is even and satisfies,

$$\forall z, \quad \int_{1}^{+\infty} N(sz) s^{d-1} ds \leq CN(z)$$

then for any smooth convex function Φ one gets :

$$\forall t \geq 0, \quad \operatorname{Ent}_{u_{\infty}}^{\Phi} \left(\begin{array}{c} \frac{u(t)}{u_{\infty}} \end{array} \right) \leq e^{-\frac{t}{c}} \operatorname{Ent}_{u_{\infty}}^{\Phi} \left(\begin{array}{c} \frac{u_{0}}{u_{\infty}} \end{array} \right)$$

□ > < @ > < E > < E > E の Q (~ 21/22

$$rac{d}{dt} \ {\cal E}^\phi_\mu({m v}(t)) = -{\cal F}^\phi_\mu(t)$$

it is enough to compare F^{ϕ}_{μ} with E^{ϕ}_{μ} .

A functional inequality [Wu'00,Chafaï'04]

If μ is an infinite divisible law (without gaussian part) ϕ satisfies $\phi'' > 0$ and ...

Then $E^{\phi}_{\mu}(f) \leq 1 \int D_{\phi}(v(x+z),v(x)) \nu_{\mu}(dz) \mu(dx)$

 u_{μ} is the derivation associated to the probability measure μ .

Example : If $\phi(x) = \ln x - x - 1$ and for the Gaussian measure this is exactly the Log-Sobolev inequality

 \rightarrow Generalization of Log-Sobolev inequality to the infinite divisible law.

$$rac{d}{dt} \ {\cal E}^\phi_\mu({m v}(t)) = -{\cal F}^\phi_\mu(t)$$

it is enough to compare F^{ϕ}_{μ} with E^{ϕ}_{μ} .

A functional inequality [Wu'00,Chafaï'04]

If μ is an infinite divisible law (without gaussian part) ϕ satisfies $\phi'' > 0$ and ...

Then $E^{\phi}_{\mu}(f) \leq \int D_{\phi}(v(x+z),v(x)) \, u_{\mu}(dz) \mu(dx)$

lf

 $\nu_{\mu} \leq \mathbf{C}\nu_{I},$

This is true for fractional Laplacians and for Lévy process near α -stable process.

$$rac{d}{dt} \ {\cal E}^\phi_\mu({m v}(t)) = -{\cal F}^\phi_\mu(t)$$

it is enough to compare F^{ϕ}_{μ} with E^{ϕ}_{μ} .

A functional inequality [Wu'00,Chafaï'04]

If μ is an infinite divisible law (without gaussian part) ϕ satisfies $\phi'' > 0$ and ...

Then $E^{\phi}_{\mu}(f) \leq \mathbf{C} \int D_{\phi}(v(x+z),v(x)) \, \nu(dz) \mu(dx)$

lf

 $\nu_{\mu} \leq \mathbf{C}\nu_{\mathbf{I}},$

This is true for fractional Laplacians and for Lévy process near α -stable process.

$$rac{d}{dt} \ {\cal E}^\phi_\mu({m v}(t)) = -{\cal F}^\phi_\mu(t)$$

it is enough to compare F^{ϕ}_{μ} with E^{ϕ}_{μ} .

A functional inequality [Wu'00,Chafaï'04]

If μ is an infinite divisible law (without gaussian part) ϕ satisfies $\phi'' > 0$ and ...

Then $E^{\phi}_{\mu}(f) \leq \mathbf{CF}^{\phi}_{\mu}(f)$

lf

 $\nu_{\mu} \leq \mathbf{C}\nu_{\mathbf{I}},$

This is true for fractional Laplacians and for Lévy process near α -stable process.

$$rac{d}{dt} \; {\cal E}^\phi_\mu({m v}(t)) = -{\cal F}^\phi_\mu(t)$$

it is enough to compare F^{ϕ}_{μ} with E^{ϕ}_{μ} .

A functional inequality [Wu'00,Chafaï'04]

If μ is an infinite divisible law (without gaussian part) ϕ satisfies $\phi'' > 0$ and ...

Then $E^{\phi}_{\mu}(f) \leq \mathbf{CF}^{\phi}_{\mu}(f)$

lf

 $\nu_{\mu} \leq \mathbf{C}\nu_{\mathbf{I}},$

then

$$E^{\phi}_{\mu}(u/u_{\infty}) \leq E^{\phi}_{\mu}\left(rac{u_{0}}{u_{\infty}}
ight) \mathrm{e}^{-rac{t}{C}}$$

This is true for fractional Laplacians and for Lévy process near α -stable process.

か q (や 22/22