

# A Lévy-Fokker-Planck equation: entropies and convergence to equilibrium

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# Ornstein-Uhlenbeck and Fokker-Planck equations

$$\begin{cases} dX_t = \sqrt{2}dB_t - \nabla V(X_t)dt \\ X_0 = x \end{cases}$$

where  $B_t$  is a standard Brownian Motion in  $\mathbb{R}^n$ .

Ito formula implies : the Semigroup  $P_t f(x) = E_x(f(X_t))$  satisfies the PDE

$$\begin{cases} \frac{\partial}{\partial t} P_t f(x) = L P_t f(x) \\ P_0 f = f, \end{cases}$$

where  $Lf = \Delta f - \nabla V \cdot \nabla f$  is the IG of  $P_t$ . This is the **Ornstein-Uhlenbeck equation**.

Consider  $L^*$  or  $P_t^*$ , the dual with respect to  $dx$ ,

$$\int Lfgdx = \int fL^*gdx, \quad \text{or} \quad \int P_t fgdx = \int fP_t^*gdx,$$

then

$$L^*g = \Delta g + \text{div}(g \cdot \nabla V).$$

The Semigroup  $P_t^* f(x)$  satisfies the PDE

$$\begin{cases} \frac{\partial}{\partial t} P_t^* f(x) = L^* P_t^* f(x) \\ P_0^* f = f, \end{cases}$$

This is the **Fokker-Planck equation**.

Let  $\mu_V = e^{-V} dx$  (**assume that  $\mu_V$  is a probability measure**),  $(P_t)_{t \geq 0}$  or  $L$  is self adjoint in  $L^2(\mu_V)$  and the by integration by parts

$$\int Lf g d\mu_V = - \int \nabla f \cdot \nabla g d\mu_V.$$

Under smooth assumptions :

$$\lim_{t \rightarrow \infty} P_t f(x) = \int f d\mu_V.$$

or equivalently

$$\lim_{t \rightarrow \infty} e^{V(x)} P_t^* g(x) = \int g dx.$$

**The good question is HOW FAST ?**

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# Tools for the asymptotic behaviour

- ▶ **Poincaré inequality** : a  $L^2$  convergence.

$$\frac{d}{dt} \text{var}_{\mu_V}(P_t f) = 2 \int P_t f L P_t f d\mu_V - 0 = -2 \int |\nabla P_t f|^2 d\mu_V,$$

If Poincaré inequality holds

$$\text{var}_{\mu_V}(f) \leq C \int |\nabla f|^2 d\mu_V$$

$$\text{var}_{\mu_V}(P_t f) \leq e^{-2t/C} \text{var}_{\mu_V}(f).$$

- ▶ **Logarithmic Sobolev inequality** a  $L \log L$  convergence

$$\frac{d}{dt} \text{Ent}_{\mu_V}(P_t f) := \frac{d}{dt} \int P_t f \log \frac{P_t f}{\int P_t f d\mu_V} d\mu_V = -4 \int |\nabla \sqrt{P_t f}|^2 d\mu_V,$$

If Logarithmic Sobolev inequality holds

$$\text{Ent}_{\mu_V}(f^2) \leq C \int |\nabla f|^2 d\mu_V$$

$$\text{Ent}_{\mu_V}(P_t f) \leq e^{-4t/C} \text{Ent}_{\mu_V}(f).$$

When do we have a Poincaré or a logarithmic Sobolev inequality ?

- \* The Gaussian measure,  $V(x) = x^2/2$  (Inequality proved by Gross).
- \* The Bakry-Emery  $\Gamma_2$ -criterion implies that if

$$\text{Hess}(V) \geq \lambda \text{Id},$$

with  $\lambda > 0$  then logarithmic Sobolev inequality holds with  $C = 2/\lambda$  and Poincaré inequality holds with  $C = 1/\lambda$ .

- \* There are also many technical methods to prove Poincaré or Log-Sobolev : Hardy, transportation...



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# Definition of Lévy process

Lévy process  $L_t$  = process with stationary & indep increments

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Fourier transform  $(L_t) = e^{t\psi(\xi)}$  where  $\psi$  is characterized by the Lévy-Khinchine formula.

$$\psi(\xi) = -\sigma\xi \cdot \xi + ib \cdot \xi + \int (e^{iz \cdot \xi} - 1 - iz \cdot \xi \mathbf{1}_B(\xi)) \nu(dz)$$

where  $\nu$  is a singular measure satisfying

$$\int_B |z|^2 \nu(dz) < +\infty \qquad \int_{\mathbb{R}^d \setminus B} \nu(dz) < +\infty,$$

$\sigma$  is a positive definite matrix and  $b$  is a vector.

Parameters  $(\sigma, b, \nu)$  characterize the Lévy process (or a infinite divisible law).

► For all  $t > 0$  the law of  $L_t$  is an infinite divisible law :

$$\mu = \underbrace{\mu_n \star \cdots \star \mu_n}_{n \text{ times}}$$

► **Associated infinitesimal generators** as for the Brownian Motion.

$$I(u) = \operatorname{div}(\sigma \nabla u) + b \cdot \nabla u + \int (u(x+z) - u(x) - \nabla u(x) \cdot z \mathbf{1}_B(z)) \nu(dz)$$

These operators appear **everywhere**

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- Laws with **heavy tails** (decrease as power laws)
- Example :  $(\sigma, b, \nu) = (0, 0, \frac{1}{|z|^{\alpha+d}} dz)$ , the  $\alpha \in (0, 2)$  stable process. In that case  $\psi(\xi) = |\xi|^\alpha$ .
- The case  $\alpha = 2$  is the Brownian motion,  $I = \Delta$ .

# The Lévy-Fokker-Planck equation

Replace  $\Delta$  by  $I$  a IG of a Lévy process in the Fokker-Planck equation :

$$\begin{cases} \frac{\partial}{\partial t} u = I(u) + \operatorname{div}(ux) \\ u(0, x) = f(x) \end{cases} \quad (\text{LFP})$$

The goal is to understand the asymptotic behaviour of the semigroup.

Remark : We fix now  $(\sigma, b, \nu)$  and assume that  $V = x^2/2$ .

Starting point of this work : [Biler and Karch \(2001\)](#)

Questions :

- Find a steady state as  $e^{-V}$  as for the classical case  $\Delta$ .
- Find the asymptotic behaviour of the Lévy-Fokker-Planck equation (LFP).
- Find conditions to get an asymptotic behaviour using inequalities as Poincaré or logarithmic Sobolev.

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An equilibrium  $u_\infty \stackrel{\text{def}}{=} a$  stationary solution of the LFP  
 $u_\infty$  can be seen as an **invariant measure**  $\mu_V$  in the case of the  
Laplacian.

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## Proposition (Existence of an equilibrium)

Assume that

$$\int_{\mathbb{R}^d \setminus B} \ln |z| \nu(dz) < +\infty. \quad (\text{Con 1})$$

There then exists an **positive** equilibrium  $u_\infty$  :

$$l(u_\infty) + \text{div}(u_\infty x) = 0.$$

Moreover,  $u_\infty dx$  is an **infinite divisible law** whose characteristic  
exponent  $A$  is

$$A(\xi) = \int_0^1 \psi(s\xi) \frac{ds}{s}.$$

The condition

$$\int_{\mathbb{R}^d \setminus B} \ln |z| \nu(dz) < +\infty. \quad (\text{Con1})$$

is satisfied for the  $\alpha$ -stable Lévy process. In that case  $u_\infty$  is the infinite divisible law of the Lévy process,  $A = \psi/\lambda$ .

Proof : The Fourier transform  $\hat{u}_\infty$  satisfies

$$\psi(\xi) \hat{u}_\infty + \xi \cdot \nabla \hat{u}_\infty = 0$$

so that  $\hat{u}_\infty = \exp(-A)$  with  $A$  such that :

$$\nabla A(\xi) \cdot \xi = \psi(\xi),$$

then

$$A(\xi) = \int_0^1 \psi(s\xi) \frac{ds}{s}.$$

Con1 proves that  $A$  is well defined and is the characteristic exponent of a Lévy process.

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For  $\phi : \mathbb{R}^+ \rightarrow \mathbb{R}$  **convex** and smooth and  $\mu$  a probability measure, consider the  $\phi$ -entropy

$$E_{\mu}^{\phi}(f) = \int \phi(f) d\mu - \phi\left(\int f d\mu\right)$$

### ► Examples

For  $\phi(x) = \frac{1}{2}x^2$  ( $E_{\mu}^{\phi}$ =the variance),  $D_{\phi}(a, b) = \frac{1}{2}(a - b)^2$

$$F_{\mu}^{\phi}(v) = \frac{1}{2} \iint (v(x+z) - v(x))^2 \nu(dz) \mu(dx)$$

For  $\phi(x) = x \ln x - x - 1$  ( $E_{\mu}^{\phi}$ =entropy),  $D_{\phi}(a, b) = a \ln \frac{a}{b} + b - a$

This is a natural interpolation between the variance and the Entropy.

Define also a **Bregman distance**

$$D_\phi(a, b) = \phi(a) - \phi(b) - \phi'(b)(a - b) \geq 0$$

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## Theorem

Let  $\mu(dx) = u_\infty(x)dx$ ,  $\nu$  the Lévy measure associated to  $I$  and consider  $v(t, x) = \frac{u(t, x)}{u_\infty(x)}$ , then

$$\frac{d}{dt} E_\mu^\phi(v(t, \cdot)) = - \iint D_\phi(v(x+z), v(x)) \nu(dz) \mu(dx).$$

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### ► Fisher information

$$F_\mu^\phi(v) = \iint D_\phi(v(x+z), v(x)) \nu(dz) \mu(dx).$$

Can be seen as a **Dirichlet form** with respect to the measure  $u_\infty(x)dx$

The proof is related to :

► **A related equation : the Lévy-Ornstein-Uhlenbeck equation (LOU)**

The function  $v = u/u_\infty$  satisfies

$$\partial_t v = \frac{1}{u_\infty} \left( I(u_\infty v) - I(u_\infty) v \right) + x \cdot \nabla v \stackrel{\text{def}}{=} Lv.$$

Dual operator of  $L$  wrt  $\mu = u_\infty(x) dx$

$$\int w_1 \left( Lw_2 \right) d\mu = \int \left( \check{I}(w_1) - x \cdot \nabla w_1 \right) w_2 d\mu,$$

where  $\check{I}$  is  $I$  with  $\check{\nu}(dx) = \nu(-dx)$ .

*Recall that in the classical case  $L$  is a self-adjoint operator with respect to  $\mu$ .*

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# Convergence towards the equilibrium

## Theorem

We assume that  $\nu_t$  has a density  $N$  with respect to  $dx$  and satisfies

$$\int_{\mathbb{R}^d \setminus B} \ln |z| N(z) dz < +\infty. \quad (\text{Con 1})$$

If  $N$  is even and satisfies,

$$\forall z, \int_1^{+\infty} N(sz) s^{d-1} ds \leq CN(z) \quad (\text{Con 2})$$

then for any smooth convex function  $\Phi$  one gets :

$$\forall t \geq 0, \quad \text{Ent}_{u_\infty}^\Phi \left( \frac{u(t)}{u_\infty} \right) \leq e^{-\frac{t}{c}} \text{Ent}_{u_\infty}^\Phi \left( \frac{u_0}{u_\infty} \right).$$

$$\frac{d}{dt} E_{\mu}^{\phi}(v(t)) = -F_{\mu}^{\phi}(t)$$

it is enough to compare  $F_{\mu}^{\phi}$  with  $E_{\mu}^{\phi}$ .

► **A functional inequality** [Ane-Ledoux'00, Wu'00, Chafaï'04]

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If  $\mu$  is an infinite divisible law (without gaussian part)  
 $\phi$  satisfies  $\phi'' > 0$  and ...

Then 
$$E_{\mu}^{\phi}(f) \leq 1 \int D_{\phi}(v(x+z), v(x)) \nu_{\mu}(dz) \mu(dx)$$

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$\nu_{\mu}$  is the derivation associated to the probability measure  $\mu$ .

Example : If  $\phi(x) = \ln x - x - 1$  and for the Gaussian measure this is exactly the Log-Sobolev inequality

→ Generalization of Log-Sobolev inequality to the infinite divisible law.

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---

If

$$\nu_{\mu} \leq C\nu_I,$$

This is true for fractional Laplacians and for Lévy process near  $\alpha$ -stable process.

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Then  $E_{\mu}^{\phi}(f) \leq CF_{\mu}^{\phi}(f)$

---

If

$$\nu_{\mu} \leq C\nu_1,$$

then

$$E_{\mu}^{\phi}(u/u_{\infty}) \leq E_{\mu}^{\phi}\left(\frac{u_0}{u_{\infty}}\right) e^{-\frac{t}{C}}$$

This is true for fractional Laplacians and for Lévy process near  $\alpha$ -stable process.

- **Conclusion**

- Family of Entropies  $\rightarrow$  Associated Fisher information
- Sufficient condition for exponential decay to the equilibrium

- **Perspectives**

- $V(x) = x^2/2 \rightarrow$  General potential ?
- What happens if the Lévy measure  $\nu$  has atoms ?