

## **Exponential convergence to the stationary measure and hyperbolicity of the minimisers for random Lagrangian systems.**

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What happens in the deterministic case?

Random Hamilton-Jacobi equation

Minimisers and hyperbolicity

## Variational description

Consider  $\phi$  satisfying the Hamilton-Jacobi equation for a mechanical Hamiltonian on a compact manifold :

$$\phi_t + \frac{1}{2}\phi_x^2 = -V, \quad x \in M.$$

The Legendre transform of  $H(t, x, \phi_x) = \phi_x^2/2 + V$  gives us :

$$L(t, x, \dot{x}) = \frac{1}{2}\dot{x}^2 - V$$

For an initial condition  $g(x)$  at time  $t_1$ , we get the following variational description of the solutions :

$$\phi(t_2, x) = \min_{\gamma(t_2)=x} \left( g(\gamma(t_1)) + \int_{t_1}^{t_2} L(t, \gamma, \dot{\gamma}) \right).$$

Generalising to the case of random forcing is not a problem.

## Minimisers (I)

For an initial condition  $g(x)$  at time  $t_1$ , we have the following variational description :

$$\phi(t_2, x) = \min_{\gamma \in \Gamma} \left( g(\gamma(t_1)) + \int_{t_1}^{t_2} L(t, \gamma, \dot{\gamma}) \right).$$

Here  $\Gamma$  is the set of curves such that  $\gamma(t_2) = x$ . Such curves are called  **$g$ -minimisers**.

When we minimise the action  $\int L$  on a time interval with fixed endpoints and without fixing an initial condition, the corresponding curves are called **minimisers**.

## Minimisers (II)

In the same way, one can also define **one-sided minimisers** :

$$\gamma^x : [t_0, +\infty) \text{ (or } (-\infty, t_0]) \rightarrow \mathcal{S}^1$$

(one minimises the action for compactly supported in time perturbations of  $\gamma^x$  such that  $\tilde{\gamma}^x(t) = x$ ) and the **global minimisers** :

$$\gamma : (-\infty, +\infty) \rightarrow \mathcal{S}^1$$

(without  $x$  or  $t$  : again, one minimises the action for compactly supported in time perturbations of  $\gamma$ ).

A  $g$ -minimiser is a minimiser.

The expected restriction properties hold.

## Deterministic generic setting

We consider the equation with a deterministic **generic** forcing. In other words :

$$\phi_t + \frac{1}{2}\phi_x^2 = -V(x), \quad (1)$$

where  $V$  is smooth and has a unique non-degenerate maximum. We assume that this maximum is reached for  $x = 0$  and equals 0 and denote  $-V''(0) =: \lambda^2 > 0$ .

Now we linearise in a neighbourhood of  $(x, v) = (0, 0)$  and we consider the Euler-Lagrange equation :

$$\frac{dx}{dt} = v, \quad \frac{dv}{dt} = -V_x.$$

Therefore in the  $(x, v)$  coordinates, the point  $(0, 0)$  is "exponentially attractive", i.e. a minimiser on  $[0, T]$  is in a  $\exp(-\lambda T/2)$ -neighbourhood of  $(0, 0)$  at time  $T/2$ .

# Exponential convergence of the minimisers

There exists a unique global minimiser : the line  $\tilde{\gamma} \equiv 0$ . Moreover :

## Theorem 1

*There exist a constant  $C > 0$  s.th. for a minimiser*

*$\gamma : [0, +\infty) \mapsto S^1$  one has :*

$$|\gamma(t) - \tilde{\gamma}(t)| = |\gamma(t)| \leq C \exp(-\lambda t/2), \quad t \geq 0.$$

# Exponential convergence of the solutions

## Corollary 1

*There exist a constant  $\tilde{C} > 0$  s.th. if we consider a pair of initial conditions  $\phi_0, \psi_0$ , then the corresponding solutions  $\phi, \psi : [0, +\infty) \mapsto S^1$  of (1) converge to each other exponentially up to an additive constant, i.e. :*

$$\max_{x \in S^1} |\phi(t, x) - \psi(t, x) - A(t)| \leq \tilde{C} \exp(-\lambda t), \quad t \geq 0,$$

*where  $A = \phi(t/2, 0) - \psi(t/2, 0)$ .*

The results above were obtained (in a more general setting) by Iturriaga-Sanchez Morgado in 2009.



# The stochastic periodic 1D Hamilton-Jacobi equation

$$\phi_t + \phi_x^2/2 = (\nu\phi_{xx}) + \eta, \quad t \geq 0, \quad x \in S^1 = \mathbb{R}/\mathbb{Z}. \quad (1DB)$$

$\eta(t, x) = \eta^\omega(t, x)$  : smooth in space **random** force, irregular in time (white noise or "kicked").

We can consider a more general nonlinear term  $f(\phi_x)$  (under Tonelli-type convexity-growth assumptions on  $f$ ).

## Different types of forcing

- "Kicked" force :

$$\eta^\omega(x) = \sum_{i=1}^{+\infty} \delta_{t=i} \zeta_i^\omega(x),$$

where  $\zeta_i^\omega$  are non-trivial smooth i.i.d. random variables in  $L_2(S^1)$  with support in  $C^\infty(S^1)$  and finite moments in all Sobolev spaces  $H^m(S^1)$ .

**Example :** At each integer time, a kick equals 0,  $\cos(2\pi x)$ ,  $\sin(2\pi x)$  or  $\sin(2\pi x) + \cos(2\pi x)$ , with probability 1/4 each.

- White noise-type force :  $\eta^\omega(x) = w_t^\omega(x)$ , where  $w^\omega$  is an  $L_2(S^1)$ -valued Wiener process in time with all moments of  $w(t)$  bounded in  $H^m(S^1)$ ,  $m \geq 1$ .

## Stationary measure

The arguments below in this part of the talk hold uniformly with respect to the viscosity coefficient  $\nu \geq 0$  (i.e. for the stochastic HJ equation with or without the viscous term  $\nu\phi_{xx}$ ) and on the torus  $\mathbb{T}^d$ ,  $d \geq 1$ .

Solutions  $\phi$  of the stochastic HJ equation define a Markov process. We have  $\nu$ -uniform upper estimates [Bor12, Bor13] : the Bogolyubov-Krylov argument implies the existence of a stationary measure.

## Speed of convergence

The semigroup  $G_t^\omega$  corresponding to the Markov process is  $L_\infty$ -nonexpanding :

$$|G_t^\omega \phi_0 - G_t^\omega \tilde{\phi}_0|_\infty \leq |\phi_0 - \tilde{\phi}_0|_\infty.$$

Then a coupling argument (cf. [Kuksin-Shirikyan '12]) gives us algebraic convergence to the unique stationary measure if 0 is in the support of the forcing.

**Idea :** The distance between two solutions corresponding to different initial values and the same forcing becomes small since the solutions themselves become small after a long "small-force" period, and then this distance is nonincreasing.

## Speed of convergence : comments

The assumption  $0 \in \text{Supp}$  is (more or less the only) necessary assumption.

Analogous arguments are used in [Gomes-Iturriaga-Khanin-Padilla '05] (and also by Dirr-Souganidis, Debussche-Vovelle...) but there is no explicit estimate of the speed of convergence.

## Random dynamics

Two types of questions :

- Existence-uniqueness-properties of the stationary measure
- Properties of the minimisers (in particular : existence, uniqueness and hyperbolicity of the global minimiser).

Sinai '91 (using the Cole-Hopf transform) ;

E-Khanin-Mazel-Sinai '00 (**hyperbolicity**).

**Multi-d case** : [Iturriaga-Khanin '03],

[Gomes-Iturriaga-Khanin-Padilla '05].

In these papers there are additional assumptions on the forcing.

## Minimisers and the sets $\Omega_{s,t}$

### Definition 1

For an initial condition  $\psi_0(0, \cdot) : S^1 \rightarrow \mathbb{R}$  and  $0 < s < t$ , let  $\Omega_{s,t}$  be the set of points reached, at time  $s$ , by the  $\psi_0$ -minimisers on  $[0, t]$ .

One notices that  $\Omega_{s,t}$  is closed.

For a closed subset  $Z$  of  $S^1$ , the diameter of  $Z$  is defined as :

$$d(Z) = 1 - m(Z),$$

where  $m(Z)$  is the maximal length among connected components of  $S^1 - Z$ .

Other possible definition : the minimal length of a closed interval in  $S^1$  containing  $Z$ .

## Assumptions on the potentials

The goal is to prove exponential convergence of the minimisers to the global minimiser.

**Idea :** In the kicked case, if we denote by  $\mu$  the probability measure on  $L_2(S^1)$  for the kicks, we want the following :

- The forcing can be arbitrarily small :  $0 \in \text{Supp } \mu$ .
- The forcing is not too structured : there exist potentials  $f_i$ ,  $i = 1, 2, 3$  reaching their nondegenerate maxima  $x_i$  at 3 different points, s.th.  $f_i \in \text{Supp } \mu$ .

For instance, a forcing generated by  $(\cos 4\pi x, \sin 4\pi x)$  does not work (periodicity);  $(\cos 2\pi x, \sin 2\pi x)$  works.



## Assumptions on the potentials (II)

### Assumptions 1

"Kicked case" :

(i) The kicks at integer time moments  $j$  have the form

$$F^\omega(j) = \sum_{k=1}^K c_k^\omega(j) F^k,$$

where the  $F^k$  are smooth functions defined on  $S^1 = \mathbb{R}/\mathbb{Z}$ .

The vectors  $(c_k^\omega(j))_{1 \leq k \leq K}$  are i.i.d.  $\mathbb{R}^K$ -valued random variables with finite moments. Their distribution on  $\mathbb{R}^K$ , denoted  $\mu$ , is absolutely continuous with respect to  $\mu_{\text{Lebesgue}}$ .

(ii) We have  $0 \in \text{Supp } \mu$ .

(iii) The mapping  $x \mapsto (F^1(x), \dots, F^K(x))$  is an embedding (injective and homeomorphism onto its image).

## Assumptions on the potentials (III)

### Assumptions 2

*White noise case :*

(i) *The potential is of the form*

$$F^\omega(x, t) = \sum_{k=1}^K \dot{W}_k^\omega(t) F^k(x),$$

*where the  $F^k$  are smooth on  $S^1$  and the  $\dot{W}_k^\omega$  are independent white noises.*

(ii) *The mapping*

$$x \mapsto (F^1(x), \dots, F^K(x))$$

*is an embedding.*

## The separation property

### Lemma 1

*There exist  $\alpha_0 > 0$ , three pairwise disjoint open intervals  $J_i$ ,  $i = 1, 2, 3$ , and three potentials  $\tilde{F}_i$ ,  $i = 1, 2, 3$  s.t. :*

- 1) The potentials  $\tilde{F}_i$  are in the support of forcing.*
- 2) Each of the functions  $\tilde{F}_i$  reaches its minimum, denoted by  $m_i$ , at a single point  $x_i$ ; moreover these points are different.*

### Lemma 2

*Assumptions 1 or 2 imply the separation property.*

## Asymptotic behaviour of the minimisers

Here we fix  $-\infty < s \leq t < +\infty$  and we assume that the minimisers verify the assumptions mentioned above.

### Lemma 3

*Fix  $s \in \mathbb{R}$ . Then  $\omega$ -a.s., there exists a random variable with finite moments  $\tilde{C} > 0$  s.th. :*

$$d(\Omega_{s,t}) \leq \tilde{C} \exp(-\lambda(t-s)/2), \quad t \geq s.$$

### Corollary 2

*$\Omega$ -a.s., there exists a random variable with finite moments  $\tilde{C} > 0$  s.th. for any one-sided minimiser  $\gamma$  on  $[0, +\infty)$  :*

$$|\gamma(t) - \tilde{\gamma}(t)| \leq \tilde{C}(s, \omega) \exp(-\lambda t/2), \quad t \geq 0,$$

where  $\tilde{\gamma}$  is the **unique** global minimiser.

## Remarks

**Hyperbolicity** : Using the Pesin theory, Corollary 3 implies hyperbolicity of the global minimiser for the Euler-Lagrange dynamics in the phase space  $(x, u, t)$ .

**Multi-d** : A global minimiser exists and is unique ; it is hyperbolic [Khanin-Zhang]. However, we do not *a priori* have exponential contraction for the minimisers and the shock structure is not well-understood.

**Noncompact setting** :

Hoang-Khanin '03 : Forcing with a well-localised minimum.

Bakhtin-Cator-Khanin '13 : Poisson process in space-time. The minimisers coalesce exponentially.

Bakhtin '15-'16 : further extension. Positive viscosity !

## Exponential convergence

For  $\nu = 0$  et  $d = 1$  we have exponential convergence to the stationary measure [Bor '16].

In other words : two solutions corresponding to different initial conditions and the same forcing converge to each other exponentially in  $L_\infty/\mathbb{R}$  like in Iturriaga-Sanchez Morgado.

An analogous result was proved by Bec-Frisch-Khanin (2001) in  $L_\infty$  far away from the shocks.

## Exponential convergence : proof

**Proof :** For two different initial conditions, consider the position of two minimisers on  $[0, T_1 + T_2]$  at time  $T_1$ . They are localised in the  $\exp(-CT_2)$ -small sets  $\Omega_{T_1, T_1+T_2}(\psi_0^i), i = 1, 2$ .

On the other hand, the limit points  $\lim_{s \rightarrow +\infty} \Omega_{T_1, s}(\psi_0^i), i = 1, 2$  are  $\exp(-CT_1)$ -close (they are both  $\gamma(T_1)$  for some one-sided minimiser  $\gamma$ ).

Finally, we fix  $T_1 = T_2 = T/2$  and we get that at time  $T/2$  minimisers on  $[0, T]$  are  $\exp(-CT/2)$ -close, and then we conclude as in [ISM09].

# Perspectives

We want to prove that under the assumptions 1 or 2, the convergence is exponential (**uniformly in  $\nu$** ). Get rid of these assumptions?

**Main obstacle** : the viscous Lagrangian representation of Feynman-Kac type is less straightforward. We have no minimisers since now we minimise an expected value.



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