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Alzheimer's disease: analysis of a mathematical model incorporating the role of prions

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Abstract We introduce a mathematical model of the in vivo progression of Alzheimer's disease with focus on the role of prions in memory impairment. Our model consists of differential equations that describe the dynamic formation of β -amyloid plaques based on the concentrations of A β oligomers, PrP^C proteins, and the A β -×-PrP^C complex, which are hypothesized to be responsible for synaptic toxicity. We prove the well-posedness of the model and provided stability results for its unique equilibrium, when the polymerization rate of β -amyloid is constant and also when it is described by a power law.

Keywords Alzheimer · Prion · Mathematical model · Well-posedness · Stability

M. Helal

E. Hingant (⊠) Cl²MA, Universidad de Concepción, Concepción, Chile e-mail: ehingant@ci2ma.udec.cl

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Département de Mathématique, Faculté des Sciences, Université Djillali Liabes, 22000 Sidi Bel Abbès, Algeria e-mail: mhelal_abbes@yahoo.fr

L. Pujo-Menjouet Institut Camille Jordan, Université de Lyon, CNRS UMR 5208, Université Lyon 1, 43 blvd. du 11 novembre 1918, 69622 Villeurbanne cedex, France e-mail: pujo@math.univ-lyon1.fr

G. F. Webb Department of Mathematics, Vanderbilt University, 1326 Stevenson Center, Nashville, TN 37240-0001, USA e-mail: glenn.f.webb@vanderbilt.edu

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1 Introduction

1.1 What is the link between Alzheimer disease and prion proteins?

Alzheimer's disease (AD) is acknowledged as one of the most widespread diseases of age-related dementia with \approx 35.6 million people infected worldwide according to Wimo and Prince (2010). By the 2050's, this same report has predicted three or four times more people living with AD. AD affects memory, cognizance, behavior, and eventually leads to death. Apart from the social dysfunction of patients, another notable societal consequence of AD is its economic cost (\approx \$422 billion in 2009, e.g. Wimo and Prince 2010). The human and social impact of AD has driven extensive research to understand its causes and to develop effective therapies. Among recent findings are the results that imply cellular prion protein (PrP^C) is connected to memory impairment (Cissé and Mucke 2009; Cissé et al. 2011; Gimbel et al. 2010; Laurén et al. 2009; Nath et al. 2012). This connection is the focus of our modeling here, which we hope will contribute to understanding the relation of AD to prions.

The pathogenesis of AD is related to a gradual build-up of β -amyloid(A β) plaques in the brain (Duyckaerts et al. 2009; Hardy and Selkoe 2002). *β*-amyloid plaques are formed from the A β peptides obtained from the amyloid protein precursor (APP) protein cleaved at a displaced position. There exist different forms of β -amyloids, from soluble monomers to insoluble fibrillar aggregates (Chen et al. 2010; Lomakin et al. 1996; Lomakin et al. 1997; Urbanc et al. 1999; Walsh et al. 1997). It has been revealed by Selkoe (2008) that the toxicity depends on the size of these structures and recent evidence suggest that oligomers (small aggregates) play a key role in memory impairment rather than β -amyloid plaques (larger aggregates) formed in the brain. More specifically, A β oligomers cause memory impairment *via* synaptic toxicity onto neurons. This phenomenon seems to be induced by a membrane receptor, and there is evidence that this rogue agent is the PrP^C protein (Nygaard and Strittmatter 2009; Zou et al. 2011; Resenberger et al. 2011; Gimbel et al. 2010; Laurén et al. 2009) We note that this protein, when misfolded in a pathological form called PrP^{Sc}, is responsible for Creutzfeldt–Jacob disease. Indeed, it is believed that there is a high affinity between PrP^C and Aβ oligomers, at least theoretically by Gallion (2012). Moreover, the prion protein has also been identified as an APP regulator, which confirms that both are highly related (Nygaard and Strittmatter 2009; Vincent et al. 2008). This discovery offers a new therapeutic target to recover memory in AD patients, or at least slow memory depletion (Freir et al. 2011; Chung et al. 2010).

1.2 What is our objective?

Our objective here is to introduce and study a new in vivo model of AD evolution mediated by PrP^C proteins. To the best of our knowledge, no model such as the one proposed here has yet been advanced. There exist a variety of models specifically designed for Alzheimer's disease and their treatment (Achdou et al. 2012; Craft et al.

2002, 2005). Nevertheless, the prion protein has never been taken into account in the way we formulate here, and our model could help in designing new experiments and treatments.

This paper is organized as follows. We present the model in Sect. 2, and provide a well-posedness result in the particular case that β -amyloids are formed at a constant rate. In Sect. 3 we provide a theoretical study of our model in a more general context with a power law rate of polymerization, i.e. the polymerization or build-up rate depends on β -amyloid plaque size.

2 The model

2.1 A model for beta-amyloid formation with prions

The model deals with four different species. First, the concentration of A β oligomers consisting of aggregates of a few A β peptides; second, the concentration of the PrP^C protein; third, the concentration of the complex formed from one A β oligomer binding onto one PrP^C protein. These quantities are soluble and their concentration will be described in terms of ordinary differential equations. Fourth, we have the insoluble β -amyloid plaques described by a density according to their size *x*. This approach is standard in modeling prion proliferation phenomena (see for instance Greer et al. 2006; Prigent et al. 2012; Calvez et al. 2009, 2010; Gabriel 2011; Greer et al. 2006, 2007; Laurençot and Walker 2007; Prüss et al. 2006; Simonett and Walker 2006 for modeling approach and analysis). Note that the size *x* is an abstract variable that could be the volume of the aggregate. Here, however, we view aggregates as fibrils that lengthen in one dimension. The size variable *x* thus belongs to the interval ($x_0, +\infty$), where $x_0 > 0$ stands for a critical size below which the plaques cannot form. To summarize we denote, for $x \in (x_0, +\infty)$ and $t \ge 0$,

- $f(t, x) \ge 0$: the density of β -amyloid plaques of size x at time t,
- $-u(t) \ge 0$: the concentration of soluble A β oligomers (unbounded oligomers) at time *t*,
- $-p(t) \ge 0$: the concentration of soluble cellular prion proteins PrP^C at time t,
- $-b(t) \ge 0$: the concentration of A β -×-PrP^C complex (bounded oligomers) at time *t*.

Note that β -amyloid plaques are formed from the clustering of $A\beta$ oligomers. The rate of agglomeration depends on the concentration of soluble oligomers and the structure of the amyloid which is linked to its size. It occurs in a mass action between plaques and oligomers at a nonnegative rate given by $\rho(x)$, where x is the size of the plaque. This is the reason why the intentionally misused word "size" considered here (and described above) accounts for the mass of $A\beta$ oligomers that form the polymer. We assume indeed, that the mass of one oligomer is given by a "sufficiently small" parameter $\varepsilon > 0$. Thus, the number of oligomers in a plaque of mass x > 0 is x/ε which justifies our assumption that the size of plaques is a continuum. Moreover, amyloids have a critical size $x_0 = \varepsilon n > 0$, where $n \in \mathbb{N}^*$ is the number of oligomers in the critical plaque size. The amyloids are prone to be damaged at a nonnegative rate μ ,

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Parameter/variable	Definition	Unit
t	Time	Days
x	size of β-amyloidplaques	_
<i>x</i> ₀	Critical size of β-amyloidplaques	-
n	Number of oligomers in a plaque of size x_0	-
ε	Mass of one oligomer	_
λ_u	Source of A ^β oligomers	Days ⁻¹
γи	Degradation rate of AB oligomers	Days ⁻¹
λ_p	Source of PrP^C	Days ⁻¹
γ_p	Degradation rate of PrP^C	Days ⁻¹
τ	Binding rate of A β oligomers onto PrP ^C	Days ⁻¹
σ	Unbinding rate of $A\beta$ -×- PrP^C	Days ⁻¹
δ	Degradation rate of A β -×-PrP ^C	Days ⁻¹
$\rho(x)$	Conversion rate of oligomers into a plaque	$(SAF/sq)^{-1} * days^{-1}$
$\mu(x)$	Degradation rate of a plaque	Days ⁻¹

 Table 1
 Parameter description of the model

* SAF/sq means Scrapie-Associated Fibrils per square unit and is explained in detail by Rubenstein et al. (1991) (we consider plaques as being fibrils here)

possibly dependent on the size *x* of the plaques. All the parameters for A β oligomers, PrP^C, and β -amyloid plaques, such as production, binding and degradation rates, are nonnegative and described in Table 1.

Then, writing evolution equations for these four quantities, we obtain

$$\frac{\partial}{\partial t}f(x,t) + u(t)\frac{\partial}{\partial x}\left[\rho(x)f(x,t)\right] = -\mu(x)f(x,t) \quad \text{on } (x_0,+\infty) \times (0,+\infty),$$
(1)

$$\dot{u} = \lambda_u - \gamma_u u - \tau u p + \sigma b - nN(u) - \frac{1}{\varepsilon} u \int_{x_0}^{+\infty} \rho(x) f(x, t) dx \quad \text{on } (0, +\infty), \quad (2)$$

$$\dot{p} = \lambda_p - \gamma_p p - \tau u p + \sigma b \quad \text{on } (0, +\infty),$$
(3)

$$\dot{b} = \tau u p - (\sigma + \delta) b \quad \text{on } (0, +\infty).$$
 (4)

The term *N* accounts for the formation rate of a new β -amyloidplaque with size x_0 from the A β oligomers. In order to balance this term, we add the boundary condition

$$u(t)\rho(x_0)f(x_0,t) = N(u(t)), \quad t \ge 0.$$
(5)

The integral in the right-hand side of equation (2) is the total polymerization with parameters $1/\varepsilon$, since dx/ε counts the number of oligomers into a unit of length dx. Finally, the problem is completed with nonnegative initial data, a function $f^{in} \ge 0$ and u^{in} , p^{in} , $b^{in} \ge 0$, such that at time t = 0

$$f(\cdot, t = 0) = f^{in} \text{ on } (x_0, +\infty),$$
 (6)

and

$$u(t=0) = u^{in}, \quad p(t=0) = p^{in} \text{ and } b(t=0) = b^{in}.$$
 (7)

The above system (1-5) involves two formal balance laws: the first one for prion proteins

$$\frac{d}{dt}\left(b+p\right) = \lambda_p - \gamma_p p - \delta b$$

and the second for $A\beta$ oligomers

$$\frac{d}{dt}\left(b+u+\frac{1}{\varepsilon}\int\limits_{x_0}^{+\infty}xf\,dx\right)=\lambda_u-\gamma_uu-\delta b-\frac{1}{\varepsilon}\int\limits_{x_0}^{+\infty}x\mu f\,dx.$$

The total concentrations of both evolve in time according to the production and degradation rates. In Fig. 1 we give a schematic representation of these processes.

Before going further, we emphasis some modeling points:

- Modeling fibril formation involves many more complex features. Indeed, in vivo but also in vitro, their dynamics include phenomena of depolymerization, fragmentation and possible coagulation. A fully developed model would take into account all these processes. In this work we focus on the dynamics of oligomers and their interactions with fibrils and PrP^C. Therefore, we neglect the internal dynamics of polymers and their depolymerization to give priority to an apparent extension rate of fibrils.
- There exist various sizes of Aβ oligomers, from dimer up to ten or so peptides. Nevertheless, these are unstable until they reach a stable structure, and that is why



Fig. 1 Schematic diagram of the evolution processes of β -amyloidplaques, A β oligomers (bounded and unbounded), and PrP^C in the model

we assume here only one stable oligomer size, which is the one that interacts with PrP^C and that is able to form protofibrils (also called critical plaques here). We refer to the papers by Serpell (2000) and Fawzi et al. (2007) for their discussions about intermediate oligomers, fibril structure, and fibril nucleation.

2.2 An associated ODE system

In this section we investigate constant polymerization and degradation rates, i.e, rates independent of the size of the plaque involved in the process. This first approach is biologically less realistic, but technically more tractable, yet still quite challenging for an analytical study of the problem. In Sect. 3, the polymerization rate ρ will be taken more realistically as a power of x. Here we assume that $\rho(x) := \rho$ and $\mu(x) := \mu$ are positive constants. Moreover, without loss of generality, we let $\varepsilon = 1$, which only requires a rescaling of the units in the equations.

Then, we assume a pre-equilibrium hypothesis for the formation of β -amyloidplaques, as formulated by Portet and Arino (2009) for filaments, by setting $N(u) = \alpha u^n$, with $\alpha > 0$ the formation rate. It is obtained assuming n - 1 reactions lead to a fibril of size *n* from oligomers:

$$u + u \Leftrightarrow F_2,$$

$$F_2 + u \Leftrightarrow F_3$$

$$\dots$$

$$F_{n-1} + u \Leftrightarrow F_n$$

where F_i are pre-fibrils or aggregates of *i*-oligomers for i = 2, ..., n and the coefficient rate of each equation is given by K_i . So, taking all the equations at the equilibrium, we get $F_2 = K_2 u^2$, $F_3 = K_3 F_2 u = K_3 K_2 u^3$, etc. until $F_n = \alpha u^n$, where the formation rate α of a critical plaque, composed of $n \ge 1$ oligomers, is given by $\alpha = K_n \times \cdots \times K_2 > 0$. Once the length *n* is achieved, we assume the fibrils reach a stable structure. Therefore, we only take into account their polymerization and not their reverse reactions (see, for instance, a discussion about prion fibrils in Serpell (2000), Fawzi et al. (2007).

With these assumptions we are able to close the system (1–4) with respect to (5) into a system of four differential equations. Indeed, integrating (1) over $(x_0, +\infty)$ we get formally an equation over the quantity of amyloids at time $t \ge 0$

$$A(t) = \int_{x_0}^{+\infty} f(x, t) dx,$$

which is given by

$$\frac{d}{dt}A(t) - u(t)\rho f(x_0, t) = -\mu A(t).$$

We close the system using expression of the boundary (5), recalling that ρ is constant, and the fact that $N(u) = \alpha u^n$. This method has already been used on the prion model by Greer et al. (2006). Now the problem reads, for $t \ge 0$,

$$\dot{A} = \alpha u^n - \mu A,\tag{8}$$

$$\dot{u} = \lambda_u - \gamma_u u - \tau u p + \sigma b - \alpha n u^n - \rho u A, \tag{9}$$

$$\dot{p} = \lambda_p - \gamma_p p - \tau u p + \sigma b, \tag{10}$$

$$\dot{b} = \tau u p - (\sigma + \delta) b. \tag{11}$$

The mass of β -amyloidplaques is given by $M(t) = \int_{x_0}^{+\infty} x f(x, t) dx$ which satisfies an equation (formal integration of 1) that can be solved independently, since

$$\frac{d}{dt}M(t) - x_0u(t)\rho f(x_0, t) - \int_{x_0}^{+\infty} \rho u(t)f(x, t)dx = -\mu M(t).$$

Indeed, we use once again the boundary condition (5), the expression of the formation rate *N* and that $x_0 = n$ since $\varepsilon = 1$, in order to get

$$\dot{M} = n\alpha u^n + \rho u A - \mu M. \tag{12}$$

Notice that initial conditions for *A* and *M* are given by $A^{in} = \int_{x_0}^{+\infty} f^{in}(x)dx$ and $M^{in} = \int_{x_0}^{+\infty} x f^{in}(x)dx$, while the initial conditions for *u*, *p* and *b* are unchanged. The next subsection is devoted to the analysis of the system (8–11).

2.3 Well-posedness and stability of the ODE system

We prove in the following proposition the nonnegativity, existence, and uniqueness of a global solution to the system (8-11) with classical techniques from the theory of ordinary differential equations.

Proposition 1 (Well-posedness) Assume λ_u , λ_p , γ_u , γ_p , τ , σ , δ , ρ and μ are positive, and let $n \geq 1$ be an integer. For any $(A^{in}, u^{in}, p^{in}, b^{in}) \in \mathbb{R}^4_+$ there exists a unique nonnegative bounded solution (A, u, p, b) to the system (8–11) defined for all time t > 0, i.e, the solution A, u, p and b belong to $C^1_b(\mathbb{R}_+)$ and remains in the stable subset

$$S = \left\{ (A, u, p, b) \in \mathbb{R}_{+}^{4} : nA + u + p + 2b \le nA^{in} + u^{in} + p^{in} + 2b^{in} + \frac{\lambda}{m} \right\}$$
(13)

with $\lambda = \lambda_u + \lambda_p$ and $m = \min\{\mu, \gamma_u, \gamma_p, \delta\}$. Furthermore, let $M(t = 0) = M^{in} \ge 0$, and then there exists a unique nonnegative solution M to (12), defined for all time t > 0.

Proof Let $F : \mathbb{R}^4 \mapsto \mathbb{R}^4$ be given by

$$F(A, u, p, b) = \begin{pmatrix} F_1 := \alpha u^n - \mu A \\ F_2 := \lambda_u - \gamma_u u - \tau up + \sigma b - \alpha n u^n - \rho u A \\ F_3 := \lambda_p - \gamma_p p - \tau up + \sigma b \\ F_4 := \tau up - (\sigma + \delta)b \end{pmatrix}.$$

F is obviously C^1 and locally Lipschitz continuous on \mathbb{R}^4 . Moreover, if $(A, u, p, b) \in \mathbb{R}^4_+$, $F_1 \ge 0$ when A = 0, $F_2 \ge 0$ when u = 0, $F_3 \ge 0$ when p = 0, and $F_4 \ge 0$ when b = 0. Thus, the system is quasi-positive and the solution remains in \mathbb{R}^4_+ . Finally, we remark that

$$\frac{d}{dt}\left(nA+u+p+2b\right) \le \lambda - m\left(nA+u+p+2b\right),$$

with $\lambda = \lambda_u + \lambda_p$ and $m = \min \{\mu, \gamma_u, \gamma_p, \delta\} > 0$, and Gronwall's lemma ensures that

$$nA(t) + u(t) + p(t) + 2b(t) \le nA^{in} + u^{in} + p^{in} + 2b^{in} + \frac{\lambda}{m}$$

This proves the global existence of a unique nonnegative bounded solution (A, u, p, b). The claim for the mass M is straightforward.

We next consider the existence of a steady state A_{∞} , u_{∞} , p_{∞} , b_{∞} and the asymptotic behavior of solutions to (8–11). It is easy to compute the steady state by solving the problem

$$\mu A_{\infty} - \alpha u_{\infty}^n = 0 \tag{14}$$

$$\lambda_u - \gamma_u u_\infty - \tau u_\infty p_\infty + \sigma b_\infty - \alpha n u_\infty^n - \rho u_\infty A_\infty = 0 \tag{15}$$

$$\lambda_p - \gamma_p p_\infty - \tau u_\infty p_\infty + \sigma b_\infty = 0 \tag{16}$$

$$\tau u_{\infty} p_{\infty} - (\delta + \sigma) b_{\infty} = 0 \tag{17}$$

From the structure of the second equation, we cannot give an explicit formula for this problem. To obtain u_{∞} we have to solve an algebraic equation, which involves a polynomial of degree *n*. However, we can prove that the solution exists, and then u_{∞} is given implicitly. The next proposition establishes the local stability of the steady state.

Theorem 1 (Linear Stability) Under hypothesis of the Proposition 1, there exists a unique positive steady state A_{∞} , u_{∞} , p_{∞} and b_{∞} to (8–11) with

$$A_{\infty} = \frac{\alpha}{\mu} u_{\infty}^n, \quad p_{\infty} = \frac{\lambda_p}{\tau^* u_{\infty} + \gamma_p}, \quad b_{\infty} = \frac{1}{\sigma} \frac{\lambda_p (\tau - \tau^*)}{\tau^* u_{\infty} + \gamma_p} u_{\infty},$$

where $\tau^* = \tau (1 - \sigma/(\delta + \sigma))$ and u_{∞} is the unique positive root of Q, defined by

$$Q(x) = \gamma_p \lambda_u + ax - P(x)$$
, for every $x \ge 0$

with $a = \tau^* (\lambda_u - \lambda_p) - \gamma_u \gamma_p$ and

$$P(x) = \tau^* \gamma_u x^2 + \alpha \gamma_p n x^n + \left(\alpha \tau^* n + \rho \gamma_p \frac{\alpha}{\mu}\right) x^{n+1} + \rho \tau^* \frac{\alpha}{\mu} x^{n+2}$$

Moreover, this equilibrium is locally linearly asymptotically stable.

Proof First, Eq. (14) gives A_{∞} with respect to u_{∞} . Then, combining (16) and (17) we get p_{∞} and b_{∞} as functions of u_{∞} . Now replacing p_{∞} and b_{∞} in (15) we get u_{∞} as the root of Q. It is straightforward that Q has a unique positive root. Indeed, it is the intersection between a line and a monotonic polynomial on the half plane. Now, we linearize the system in A_{∞} , u_{∞} , p_{∞} and b_{∞} . Let $X = (A, u, p, b)^T$ and the linearized system reads

$$\frac{d}{dt}X = DX$$

where

$$D = \begin{pmatrix} -\mu & \alpha n u_{\infty}^{n-1} & 0 & 0 \\ -\rho u_{\infty} & \gamma_u - \tau p_{\infty} - \alpha n^2 u_{\infty}^{n-1} - \rho A_{\infty} & -\tau u_{\infty} & \sigma \\ 0 & -\tau p_{\infty} & -(\gamma_p + \tau u_{\infty}) & \sigma \\ 0 & \tau p_{\infty} & \tau u_{\infty} & -(\sigma + \delta) \end{pmatrix}$$

The characteristic polynomial is of the form

$$P(\lambda) = \lambda^4 + a_1 \lambda^3 + a_2 \lambda^2 + a_3 \lambda + a_4,$$

with the $a_i > 0$, $i = 1 \dots 4$ given in the Appendix. Moreover it satisfies

$$a_1 a_2 a_3 > a_3^2 + a_1^2 a_4.$$

Then, according to the Routh–Hurwitz criterion (see Allen 2007*Th. 4.4, page 150), all the roots of the characteristic polynomial P are negative or have negative real part, thus the equilibrium is locally asymptotically stable.

To go further, we give a conditional global stability result when no nucleation is considered, i.e., $\alpha = 0$.

Proposition 2 (Global stability) *Assume that* $\alpha = 0$. *Under the condition*

$$\left(1+2\frac{\delta+\gamma_u}{\sigma}\right)>\frac{\delta}{2\gamma_p}>\frac{\gamma_p}{\sigma},$$

the unique equilibrium is given by

$$A_{\infty} = 0, \quad p_{\infty} = \frac{\lambda_p}{\tau^* u_{\infty} + \gamma_p}, \quad b_{\infty} = \frac{1}{\sigma} \frac{\lambda_p (\tau - \tau^*)}{\tau^* u_{\infty} + \gamma_p} u_{\infty},$$

where u_{∞} is the unique positive root of $Q(x) = \gamma_p \lambda_u + ax - \tau^* \lambda_u x^2$, with $a = \tau^*(\lambda_u - \lambda_p) - \gamma_u \gamma_p$. Further, this equilibrium is globally asymptotically stable in the stable subset S defined in (13).

Proof The proof is given by a Lyapunov function Φ stated in the Appendix. It is positive when the condition above is fulfilled and its derivative along the solution to the system (8–11) is negative definite. Thus, from the LaSalle's invariance principle, we get that under these hypotheses the equilibrium of (8–11) is globally asymptotically stable.

In the next section we will study from a mathematical point of view a more realistic model. Nevertheless, our model emphasizes a major dilemma in AD. Indeed, consider the steady state given in Theorem 1. If the rate of polymerization increases, it increases the growth rate of the polynomial P, so the intersection occurs faster (the positive root of Q). This means that u_{∞} decreases, and likewise b_{∞} . The balance law of oligomers suggests that when b_{∞} decreases in such a way, the mass of fibrils M will increase. So a question remains, what is the less toxic quantity, and is there any criteria under which we could optimize ρ .

3 A power law polymerization rate

The assumption that the polymerization rate ρ and the degradation rate μ are constant is not always biologically realistic, as recognized by Calvez et al. (2010) and Gabriel (2011). Consequently, we study here the more realistic case $\rho(x) \sim x^{\theta}$, and in the following we restrict our analysis to $\theta \in (0, 1)$. We will see that we are able to obtain a result of existence and uniqueness of solutions for this more general case.

3.1 Hypotheses and main result

We are interested in nonnegative solutions to the system (1-4) with the boundary condition (5), completed by initial data (6) and (7), but with the new assumption $\rho(x) \sim x^{\theta}$. Moreover, we require that our solution preserves the total mass of β -amyloidin order to be biologically relevant. Hence, the solution f will be sought in the natural space $L^1(x_0, +\infty; xdx)$, since xdx measures the mass at any time. Our hypotheses for the system (1-4) are

(H1)
$$f^{in} \in L^1(x_0, +\infty; xdx), f^{in} \ge 0, a.e. x > x_0.$$

I

(H2)
$$\rho \ge 0, \ \rho \in W^{2,\infty}([x_0,\infty)), \ \mu \ge 0, \ \mu \in W^{1,\infty}([x_0,\infty)).$$

(H3) $N \ge 0, \ N \in W^{1,\infty}_{loc}(\mathbb{R}_+), \ N(0) = 0.$
(H4) $\lambda_u, \ \gamma_u, \ \lambda_p, \ \gamma_p, \ \tau, \ \sigma, \ \delta > 0.$

We note that (H2) implies the existence of a constant C > 0 such that $\rho(x) \le Cx$, with for example, $C = 2 \|\rho'\|_{L^{\infty}} + \rho(x_0)/x_0$. For any $x \ge x_0$, we have

$$\rho(x) \le \|\rho'\|_{L^{\infty}}(x+x_0) + \rho(x_0) \le \left(2\|\rho'\|_{L^{\infty}} + \frac{\rho(x_0)}{x_0}\right)x.$$

We remark that this kind of regularity of the rate ρ covers the case that $\rho(x) \sim x^{\theta}$ with $\theta \in (0; 1)$. Also, (H3) implies the existence of a constant $K_M > 0$ such that $N(w) \leq K_M w$, for any $w \in [0, M]$. Further, The nonnegativity of the parameters of Table 1 (hypothesis (H4)) is a natural assumption with regard to their biological meaning.

We introduce the definition of a solution to system (1-4).

Definition 1 Consider a function f^{in} satisfying (H1) and let u^{in} , p^{in} , b^{in} be three nonnegative real data. Assume that ρ , μ , N and all the parameters of Table 1 verify assumptions (H2)–(H4), and let T > 0. Then a quadruplet (f, u, p, b) of nonnegative functions is said to be a *solution* on the interval (0, T) to the system (1-4) with the boundary condition (5) and the initial data (6) and (7), if it satisfies, for any $\varphi \in C_c^{\infty}([0, T] \times [x_0, +\infty))$ and $t \in (0, T)$

$$\int_{x_0}^{+\infty} f(x,t)\varphi(x,t)dx = \int_{x_0}^{+\infty} f^{in}(x)\varphi(x,0)dx + \int_{0}^{t} N(u(s))\varphi(x_0,s)ds$$
$$+ \int_{0}^{t} \int_{x_0}^{+\infty} f(x,s) \left[\frac{\partial}{\partial t}\varphi(x,s) + u(s)\rho(x) \frac{\partial}{\partial x}\varphi(x,s) - \mu(x)\varphi(x,s) \right] dxds,$$

and

$$u(t) = u^{in} + \int_0^t \left[\lambda_u - \gamma_u u - \tau u p + \sigma b - x_0 N(u) - u \int_{x_0}^{+\infty} \rho(x) f(x, s) dx \right] ds,$$

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$$p(t) = p^{in} + \int_{0}^{t} \left[\lambda_p - \gamma_p p - \tau u p + \sigma b\right] ds,$$
$$b(t) = b^{in} + \int_{0}^{t} \left[\tau u p - (\sigma + \delta)b\right] ds,$$

with the regularity $f \in L^{\infty}(0, T; L^1(x_0, +\infty; xdx))$ and $u, p, b \in C^0(0, T)$.

Theorem 2 (Well-posedness) Let f^{in} be a nonnegative function satisfying (H1), let u^{in} , p^{in} and b^{in} be nonnegative real numbers, and assume hypothesis (H2) to (H4). Let T > 0. There exists a unique nonnegative solution (f, u, p, b) to (1-4) with (5) and initial conditions given by (6) and (7), in the sense of Definition 1, such that $f \in C^0([0, T], L^1(x_0, +\infty; x^r dx))$ for every $r \in [0, 1]$, and $u, p, b \in C_b^1(0, T)$.

The proof of the Theorem 2 is decomposed into two parts. First, we study the initial boundary value problem

$$\frac{\partial}{\partial t}f(x,t) + u(t)\frac{\partial}{\partial x}\left[\rho(x)f(x,t)\right] = -\mu(x)f(x,t) \quad \text{on } (x_0,+\infty) \times (0,+\infty),$$
(18)

$$u(t)\rho(x_0)f(x_0,t) = N(u(t)), \text{ on } (0,+\infty),$$
(19)

$$f(\cdot, t = 0) = f^{in}, \text{ on } (x_0, +\infty).$$
 (20)

We prove in the Sect. 3.2 the following proposition:

Proposition 3 Let $u \in C_b^0(\mathbb{R}_+)$, let f^{in} satisfy (H1), and assume hypothesis (H2) to (H3). For any T > 0, there exists a unique nonnegative solution f to (18–20) in the sense of distributions, such that $f \in C^0([0, T], L^1(x_0, +\infty; x^r dx))$ for every $r \in [0, 1]$.

The proof is in the spirit of the proof proposed by Collet and Goudon (2000) for the Lifshitz–Slyozov equation. It consists of a proof based on the concept of a mild solution in the sense of distributions, with the additional requirement of continuity from time into $L^1(xdx)$ space.

The second step of the proof of Theorem 2 is performed in Sect. 3.3. Precisely, once we have the existence of a unique density f, when u is given, we are able to construct the operator

$$S: C^{0}([0, T])^{3} \mapsto C^{0}([0, T])^{3}$$

$$(u, p, b) \mapsto (S_{u}, S_{p}, S_{b}) = S(u, p, b),$$

$$S_{u} = u^{in} + \int_{0}^{t} \left[\lambda_{u} - \gamma_{u}u - \tau up + \sigma b - x_{0}N(u) - u \int_{x_{0}}^{+\infty} \rho(x)f(x, s)dx \right] ds,$$
(21)

$$S_p = p^{in} + \int_0^t \left[\lambda_p - \gamma_p p - \tau u p + \sigma b \right] ds,$$

$$S_b = b^{in} + \int_0^t \left[\tau u p - (\sigma + \delta) b \right] ds,$$

where f is the unique solution associated to u given by Proposition 3. Then, Theorem 2 is finally proven in Sect. 3.3 applying the Banach fixed point theorem to the operator S.

3.2 Existence of a solution to the autonomous problem

In the following we let $u \in C_b^0(\mathbb{R}_+)$ and we use the notations $a(x, t) = u(t)\rho(x)$ and $c(x, t) = -u(t)\rho'(x)$ for every $(x, t) \in [x_0, +\infty) \times \mathbb{R}_+$. From (H2) and noting that $\rho(x) \leq Cx$, we have for any t > 0

$$a(t, x) \le Ax, \text{ for } x > x_0, \tag{22}$$

 $|a(t, x) - a(t, y)| \le A|x - y|, \text{ for } x, y > x_0,$ (23)

$$|c(t,x)| \le B,\tag{24}$$

where $A = \max (C ||u||_{L^{\infty}}, ||u||_{L^{\infty}} ||\rho'||_{L^{\infty}})$ and $B = ||u||_{L^{\infty}} ||\rho'||_{L^{\infty}(x_0, +\infty)}$. In order to establish the mild formulation of the problem, we define the characteristic reaching $x \ge x_0$ at time $t \ge 0$, that is, the solution to

$$\frac{d}{ds}X(s;x,t) = a(t,X(s;x,t)),$$

$$X(t;x,t) = x.$$
(25)

From property (23), their exists a unique characteristic that reaches (x, t). We note that it makes sense as long as $X(s; x, t) \ge x_0$. Thus, we define the starting time of the characteristic as

$$s_0(x, t) := \inf \{ s \in [0, t] : X(s; x, t) \ge x_0 \}.$$

The characteristic will be defined for any time $s \ge s_0$ and takes its origin from the initial or the boundary condition, respectively, if $s_0 = 0$ or $s_0 > 0$. We recall the classical properties of these characteristics

$$X(s; X(\sigma; x, t), \sigma) = X(s; x, t)$$
$$J(s; x, t) := \frac{\partial}{\partial x} X(s; x, t) = \exp\left(\int_{s}^{t} c(\sigma, X(\sigma; x, t)) d\sigma\right)$$
$$\frac{\partial}{\partial t} X(s; x, t) = -a(t, x) J(s; x, t).$$

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Also, remarking that $s_0(X(t; x_0, 0), t) = 0$, then by monotonicity and continuity of X for any t > 0, we get $x \in (x_0, X(t; x_0, 0)) \iff s_0(x, t) \in (0, t)$, and for any $x \in (x_0, X(t; x_0, 0))$ we have $X(s_0(x, t); x, t) = x_0$. It follows that for every x belongs to $(x_0, X(t; x_0, 0))$

$$I(x,t) := -\frac{\partial}{\partial x} s_0(x,t) = J(s_0(x,t); x,t) / a(s_0(x,t), x_0).$$

Considering the derivative of f(s, X(s; x, t)) in *s*, and integrating over (s_0, t) we obtain the mild formulation of the problem. The mild solution is defined for *a.e.* $(x, t) \in (x_0, +\infty) \times \mathbb{R}_+$ by

$$f(x,t) = \begin{cases} f^{in}(X(0;x,t))J(0;x,t)e^{-\int_0^t \mu(X(\sigma;x,t))d\sigma} & x \ge X(t;x_0,0), \\ N(u(s_0(x,t)))I(x,t)e^{-\int_{s_0(x,t)}^t \mu(X(\sigma;x,t))d\sigma} & x \in (x_0,X(t;x_0,0)). \end{cases}$$
(26)

We infer from the formulation (26) that for a.e $(x, t) \in [x_0, +\infty) \times \mathbb{R}_+$, f is non-negative, since J and I are nonnegative, and f^{in} satisfies (H1). We recall some useful properties that are derived in Lemma 1 from the paper by Collet and Goudon (2000).

Lemma 1 Let $u \in C_b^0(\mathbb{R}_+)$ be a given data and assume that (H2) holds. Then for any $x \ge x_0$ and t > 0, as long as the characteristic curve $s \mapsto X(s; x, t)$ defined in (25) exists, i.e., $s \ge s_0(x, t)$, we have

for
$$s_1 \leq s_2$$
, $X(s_1; x, t) \leq X(s_2; x, t) \leq X(s_1; x, t)e^{A(s_2-s_1)}$
if $x_n \to +\infty$, then for all $t \geq s \geq 0$, $X(s; x, t) \to +\infty$
for $s \geq t$, $X(s; x, t) \leq xe^{A(s-t)}$.

Proof We refer to the proof given by Collet and Goudon (2000), where the result follows from the fact that for any $x \ge x_0$, t > 0 and $s_0(x, t) \le s_1 \le s_2$, we have

$$\begin{aligned} x_0 &\leq X(s_2; x, t) = X(s_1; x, t) + \int_{s_1}^{s_2} a(s, X(s; x, t)) ds &\leq X(s_1; x, t) \\ &+ A \int_{s_1}^{s_2} X(s; x, t) ds, \end{aligned}$$

where A is given by (22).

In the sequel we will repeatedly refer to the changes of variables

$$y = X(0; x, t)$$
 over $x \in (X(t, x_0, 0), +\infty)$, with Jacobian $J(0; x, t)$,
 $s = s_0(x, t)$ over $x \in (x_0, X(t; x_0, 0))$, with Jacobian $-I(x, t)$.

The first is a C^1 -diffeomorphism from $(X(t, x_0, 0), +\infty)$ into $(x_0, +\infty)$, and the second from $(x_0, X(t; x_0, 0))$ into (0, t). Integrating f defined by (26) over (0, R) with $R > X(t; x_0, 0)$, using the change of variables above, using Lemma 1, and taking the limit $R \to +\infty$, we get

$$\int_{x_0}^{+\infty} x |f(t,x)| dx \le \int_{x_0}^{+\infty} X(t;y,0) |f^{in}(y)| dy + \int_{0}^{t} X(t;s,x_0) |N(u(s))| ds$$
$$\le e^{At} \left(\int_{x_0}^{+\infty} y |f^{in}(y)| dy + \int_{0}^{t} x_0 |N(u(s))| ds \right),$$
(27)

where we have split the integral into two parts and uses both the previous changes of variables. Thus, for any T > 0, $f \in L^{\infty}(0, T; L^{1}(x_{0}, +\infty; xdx))$, and therefore in $L^{\infty}(0, T; L^{1}(x_{0}, +\infty; x^{r}dx))$, for any $r \in [0, 1]$. In the next lemma we claim that f defined by (26) is a weak solution.

Lemma 2 Let f be the mild solution defined by (26). Then for any t > 0

$$\int_{x_0}^{+\infty} f(x,t)\varphi(x,t)dx = \int_{x_0}^{+\infty} f^{in}(x)\varphi(x,0)dx + \int_{0}^{t} N(u(s))\varphi(x_0,s)ds$$
$$+ \int_{0}^{t} \int_{x_0}^{+\infty} f(x,s) \left[\frac{\partial}{\partial t}\varphi(x,s)u(s)\rho(x) \frac{\partial}{\partial x}\varphi(x,s) - \mu(x)\varphi(x,s) \right] dxds,$$

for all $\varphi \in C_c^{\infty}([0, T] \times [x_0, +\infty))$.

Proof Since *f* belongs to $L^{\infty}(0, T; L^1(x_0, +\infty; xdx))$, it is possible to multiply the mild solution *f* against a test function $\varphi \in C_c^{\infty}([0, T] \times [x_0, +\infty))$ and integrate over $(x_0, +\infty)$ to obtain

$$\int_{x_0}^{+\infty} f(x,t)\varphi(x,t)dx = \int_{x_0}^{+\infty} f^{in}(y)\varphi(X(t;y,0))e^{-\int_0^t \mu(X(\sigma;y,0))d\sigma}dy$$
$$-\int_0^t N(u(s))\varphi(X(t;x_0,s),t)e^{-\int_s^t \mu(X(\sigma;x_0,s))d\sigma}ds,$$
(28)

by the same change of variable made above for (27). Furthermore, we have

$$\int_{0}^{t} \int_{x_{0}}^{X(s;x_{0},0)} f(x,s) \left[\partial_{t}\varphi(x,s) + a(s,x)\partial_{x}\varphi(x,s) - \mu(x)\varphi(x,s)\right] dxds$$

$$= \int_{0}^{t} \int_{x_{0}}^{+\infty} f^{in}(x) \frac{d}{ds} \left(\varphi(X(s;x,0),s)e^{-\int_{0}^{s}\mu(X(\sigma;x,0))d\sigma}\right) dyds$$

$$= \int_{x_{0}}^{+\infty} f^{in}(x)\varphi(X(t;x,0),t)e^{-\int_{0}^{t}\mu(X(\sigma;y,0))d\sigma} dx - \int_{x_{0}}^{+\infty} f^{in}(x)\varphi(x,0)dx,$$
(29)

still using the change of variable mentioned above and

$$\int_{0}^{t} \int_{X(s;x_{0},0)}^{\infty} f(x,s) \left[\partial_{t}\varphi(x,s) + a(s,x)\partial_{x}\varphi(x,s) - \mu(x)\varphi(x,s)\right] dxds$$

$$= -\int_{0}^{t} \int_{0}^{s} N(u(z)) \frac{d}{ds} \left(\varphi(X(s;x_{0},z),s)e^{-\int_{z}^{s} \mu(X(\sigma;x_{0},z))d\sigma}\right) dzds$$

$$= -\int_{0}^{t} N(u(s))\varphi(X(t;x_{0},s),t)e^{-\int_{s}^{t} \mu(X(\sigma;x_{0},s))d\sigma} dzds$$

$$-\int_{0}^{t} N(u(s))\varphi(x_{0},s)ds.$$
(30)

Finally, combining (28), (29) and (30) we obtain that f is a weak solution.

The aim of the following lemma is to prove that the moments of f less than 1 are continuous in time.

Lemma 3 Let hypothesis (H1) to (H3) hold. Let f be the mild solution given by (26). Then for any T > 0,

$$f \in C^0([0, T], L^1(x_0, +\infty; x^r dx)), \text{ for every } r \in [0, 1].$$

Proof Let T > 0 and $r \in [0, 1]$, since $f \in L^{\infty}_{loc}(\mathbb{R}_+, L^1(x_0, +\infty; x^r dx))$, we have for any t > 0 and $\delta t > 0$ such that $t + \delta t \leq T$

$$\int_{x_0}^{+\infty} x^r \left| f(x, t + \delta t) - f(x, t) \right| dx = I_1 + I_2 + I_3,$$

where

$$I_{1} = \int_{x_{0}}^{X(t;x_{0},0)} x^{r} |f(x, t + \delta t) - f(x, t)| dx,$$

$$I_{2} = \int_{X(t;x_{0},0)}^{X(t+\delta t;x_{0},0)} x^{r} |f(x, t + \delta t) - f(x, t)| dx,$$

$$I_{3} = \int_{X(t+\delta t;x_{0},0)}^{+\infty} x^{r} |f(x, t + \delta t) - f(x, t)| dx.$$

Our goal is to prove that each term goes to zero when δt goes to zero. We first bound I_3 , which results from the initial condition, since for $x \ge X(t + \delta t; x_0, 0) \ge X(t; x_0, 0)$, it follows that

$$I_{3} = \int_{X(t+\delta t;x_{0},0)}^{+\infty} x^{r} \left| f^{in}(X(0;x,t+\delta t))J(0;x,t+\delta t)e^{-\int_{0}^{t+\delta t} \mu(X(\sigma;x,t+\delta t))d\sigma} - f^{in}(X(0;x,t))|J(0;x,t)e^{-\int_{0}^{t} \mu(X(\sigma;x,t))d\sigma} \right| dx.$$

Let $f_{\varepsilon}^{in} \in C_0^{\infty}$ with compact support $supp(f_{\varepsilon}^{in}) \subset (0, R_{\varepsilon})$ and converge in the space $L^1([x_0, +\infty), xdx)$ to f^{in} . We write I_3 as follows

$$I_3 = I_3^1 + I_3^2 + I_3^3, (31)$$

where

$$\begin{split} I_{3}^{1} &= \int_{X(t+\delta t;x_{0},0)}^{+\infty} x^{r} \left| f^{in}(X(0;x,t+\delta t)) - f_{\varepsilon}^{in}(X(0;x,t+\delta t)) \right| \\ &\times J(0;x,t+\delta t) e^{-\int_{0}^{t+\delta t} \mu(X(\sigma;x,t+\delta t))d\sigma} dx, \\ I_{3}^{2} &= \int_{X(t+\delta t;x_{0},0)}^{+\infty} x^{r} \left| f_{\varepsilon}^{in}(X(0;x,t+\delta t)) J(0;x,t+\delta t) \right| \\ &\times e^{-\int_{0}^{t+\delta t} \mu(X(\sigma;x,t+\delta t))d\sigma} \\ &- f_{\varepsilon}^{in}(X(0;x,t)) J(0;x,t) e^{-\int_{0}^{t} \mu(X(\sigma;x,t))d\sigma} \left| dx, \right| \\ I_{3}^{3} &= \int_{X(t+\delta t;x_{0},0)}^{+\infty} x^{r} \left| f_{\varepsilon}^{in}(X(0;x,t)) - f^{in}(X(0;x,t)) \right| \\ &\times J(0;x,t) e^{-\int_{0}^{t} \mu(X(\sigma;x,t))d\sigma} dx. \end{split}$$

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Dropping the exponential term, which is bounded by one, and changing variables $y = X(0; x, t + \delta t)$ in I_3^1 and y = X(0; x, t) in I_3^3 , we get

$$I_{3}^{1} + I_{3}^{3} \le 2e^{AT} \int_{x_{0}}^{+\infty} y^{r} |f^{in}(y) - f_{\varepsilon}^{in}(y)| dy = C_{3}^{1}(T,\varepsilon),$$
(32)

with the help of Lemma 1. Next we bound I_3^2 by

$$\begin{split} I_{3}^{2} &\leq \int_{X(t+\delta t;x_{0},0)}^{+\infty} x^{r} |f_{\varepsilon}^{in}(X(0;x,t+\delta t)) - f_{\varepsilon}^{in}(X(0;x,t))| J(0;x,t+\delta t) dx \\ &+ \int_{X(t+\delta t;x_{0},0)}^{+\infty} x^{r} f_{\varepsilon}^{in}(X(0;x,t))| J(0;x,t+\delta t) - J(0;x,t)| dx \\ &+ \int_{X(t+\delta t;x_{0},0)}^{+\infty} x^{r} f_{\varepsilon}^{in}(X(0;x,t)) J(0;x,t) \\ &\times |e^{-\int_{0}^{t+\delta t} \mu(X(\sigma;x,t+\delta t)) d\sigma} - e^{-\int_{0}^{t} \mu(X(\sigma;x,t)) d\sigma} | dx, \end{split}$$

and we denote the integrals by J_3^1 to J_3^3 , respectively. We remark that $J(0, x, t) \le e^{BT}$ by (24) and so

$$J_{3}^{1} \leq e^{BT} \| f_{\varepsilon}^{in} \|_{L^{\infty}} \int_{X(t+\delta t;x_{0},0)}^{C_{\varepsilon}} x^{r} |X(0;x,t+\delta t) - X(0;x,t)| dx$$

$$\leq \delta t e^{BT} \| f_{\varepsilon}^{in} \|_{L^{\infty}} \int_{X(t+\delta t;x_{0},0)}^{C_{\varepsilon}} x^{r} \sup_{s \in [t,t+\delta t]} \left| \frac{\partial}{\partial t} X(0;x,s) \right| dx$$

$$\leq \delta t A e^{2BT} \| f_{\varepsilon}^{in} \|_{L^{\infty}} \int_{x_{0}}^{C_{\varepsilon}} x^{r+1} dx, \qquad (33)$$

where C_{ε} depends on T, A and R_{ε} i.e., the compact support of f_{ε}^{in} . Then

$$J_3^2 \le e^{BT} \|f_{\varepsilon}^{in}\|_{L^{\infty}} \int_{X(t+\delta t;x_0,0)}^{R_{\varepsilon}} x^r |e^{G(t,\delta t,x)} - 1| dx$$

with

$$\begin{aligned} |G(t,\delta t,x)| &= \Big| \int_{0}^{t+\delta t} c(\sigma,X(\sigma;x,t+\delta t))d\sigma - \int_{0}^{t} c(\sigma,X(\sigma;x,t))d\sigma \Big| \\ &\leq \int_{0}^{t+\delta t} \Big| \rho'(X(\sigma;x,t+\delta t)) - \rho'(X(\sigma;x,t)) \Big| u(\sigma)d\sigma \\ &+ \int_{t}^{t+\delta t} \Big| c(\sigma,X(\sigma;x,t)) \Big| d\sigma. \end{aligned}$$

Thus, with (22) and (24),

$$\begin{aligned} |G(t,\delta t,x)| &\leq K \|u\|_{L^{\infty}} \int_{0}^{T} \left| X(\sigma;x,t+\delta t) - X(\sigma;x,t) \right| d\sigma + \delta t B \\ &\leq \delta t K \|u\|_{L^{\infty}} \int_{0}^{T} \sup_{s \in [t,t+\delta t]} \left| \frac{\partial}{\partial t} X(\sigma;x,s) \right| d\sigma + \delta t B \\ &\leq \delta t \left(K \|u\|_{L^{\infty}} A T e^{BT} x + B \right), \end{aligned}$$

where *K* is the Lipschitz constant of ρ' . Since $x \leq R_{\varepsilon}$, let

$$C_G(T,\varepsilon) = K \|u\|_{L^{\infty}} AT e^{BT} R_{\varepsilon} + B,$$

and if $|x| \leq y$, then

$$|e^{x} - 1| \le |e^{y} - 1| + |e^{-y} - 1|.$$

Thus, we get

$$J_3^2 \le e^{BT} \|f_{\varepsilon}^{in}\|_{L^{\infty}} \left(\left| e^{\delta t C_G(T,\varepsilon)} - 1 \right| + \left| e^{-\delta t C_G(T,\varepsilon)} - 1 \right| \right) \int_{x_0}^{R_{\varepsilon}} x^r dx.$$
(34)

Since μ is nonnegative, $J_3^3 \leq$

$$e^{BT} \| f_{\varepsilon}^{in} \|_{L^{\infty}} \int\limits_{X(t+\delta t;x_0,0)}^{R_{\varepsilon}} x^r \left| e^{-\left(\int_0^{t+\delta t} \mu(X(\sigma;x,t+\delta t)) d\sigma - \int_0^t \mu(X(\sigma;x,t)) d\sigma \right)} - 1 \right| dx.$$

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Exactly as above,

$$\left|\int_{0}^{t+\delta t} \mu(X(\sigma; x, t+\delta t)) d\sigma - \int_{0}^{t} \mu(X(\sigma; x, t)) d\sigma\right| \leq \delta t M A T e^{BT} x + \delta t \|\mu\|_{L^{\infty}},$$

with $M = \text{Lipschitz constant of } \mu$. Denoting by $C_M(T, \varepsilon) = MATe^{BT}R_{\varepsilon} + \|\mu\|_{L^{\infty}}$, we get

$$J_3^3 \le e^{BT} \|f_{\varepsilon}^{in}\|_{L^{\infty}} \left(\left| e^{\delta t C_M(T,\varepsilon)} - 1 \right| + \left| e^{-\delta t C_M(T,\varepsilon)} - 1 \right| \right) \int_{x_0}^{R_{\varepsilon}} x^r dx.$$
(35)

From (32), (33), (34) and (35) we can conclude that for any $\varepsilon > 0$,

$$I_3(\delta t) \le C_3^1(T,\varepsilon) + C_3^2(T,\delta t,\varepsilon), \tag{36}$$

with $\lim_{\varepsilon \to 0} C_3^1(T, \varepsilon) = 0$ and $\lim_{\delta t \to 0} C_3^2(T, \delta t, \varepsilon) = 0$.

Next, concerning I_1 , f can be written from the boundary condition. Let $u^{\varepsilon} \in C_0^{\infty}$ such that $u^{\varepsilon} \longrightarrow u$ uniformly on [0, T]. Then we write I_1 as follows:

$$\begin{split} I_{1} &\leq \int_{x_{0}}^{X(t+\delta t;x_{0},0)} x^{r} |N(u(s_{0}(x,t+\delta t)) - N(u^{\varepsilon}(s_{0}(x,t+\delta t)))|I(x,t+\delta t)dx \\ &+ \int_{x_{0}}^{X(t;x_{0},0)} x^{r} \left| N(u^{\varepsilon}(s_{0}(x,t+\delta t))I(x,t+\delta t)e^{-\int_{s_{0}(x,t)}^{t} \mu(X(\sigma;x,t+\delta t))d\sigma} - N(u^{\varepsilon}(s_{0}(x,t))I(x,t)e^{-\int_{s_{0}(x,t)}^{t} \mu(X(\sigma;x,t))d\sigma} \right| dx \\ &+ \int_{x_{0}}^{X(t;x_{0},0)} x^{r} |N(u(s_{0}(x,t)) - N(u^{\varepsilon}(s_{0}(x,t)))|I(x,t)dx. \end{split}$$

From (H3) we obtain, similarly to I_3 , that there exist two constants $C_1^1(T, \varepsilon)$ and $C_1^2(T, \delta t, \varepsilon)$ such that

$$I_1(\delta t) \le C_1^1(T,\varepsilon) + C_1^2(T,\delta t,\varepsilon), \tag{37}$$

with $\lim_{\varepsilon \to 0} C_1^1(T, \varepsilon) = 0$ and $\lim_{\delta t \to 0} C_1^2(T, \delta t, \varepsilon) = 0$.

Finally, for I_2 , we use the two formulas of f,

$$I_{2} = \int_{X(t;x_{0},0)}^{X(t+\delta t;x_{0},0)} x^{r} \left| N(u(s_{0}(x,t+\delta t)))I(x,t+\delta t)e^{-\int_{s_{0}(x,t+\delta t)}^{t+\delta t} \mu(X(\sigma;x,t+\delta t))d\sigma} - f^{in}(X(0;x,t))J(0;x,t)e^{-\int_{s_{0}(x,t)}^{t} \mu(X(\sigma;x,t))d\sigma} \right| dx$$

Using the Lipschitz constant of N denoted by K_N , from the definition of I and with the help of Lemma 1, we get

$$I_{2} \leq x_{0}^{r} e^{(rA+B)T} K_{N} |X(t+\delta t; x_{0}, 0) - X(t; x_{0}, 0)| + x_{0}^{r} e^{rAT} \int_{X(t; x_{0}, 0)} \left| f^{in}(X(0; x, t)) J(0; x, t) \right| dx$$

Using the regularization f_{ε}^{in} of f^{in} , there exist two constants $C_2^1(T, \varepsilon)$ and $C_2^2(T, \delta t, \varepsilon)$ such that for any $\varepsilon > 0$,

$$I_2(\delta t) \le C_2^1(T,\varepsilon) + C_2^2(T,\delta t,\varepsilon), \tag{38}$$

with $\lim_{\varepsilon \to 0} C_2^1(T, \varepsilon) = 0$ and $\lim_{\delta t \to 0} C_2^2(T, \delta t, \varepsilon) = 0$.

In conclusion, combining (36), (37) and (38), we get for any $\varepsilon > 0$ and $\delta t > 0$,

$$\int_{x_0}^{+\infty} x^r |f(x,t+\delta t) - f(x,t)| dx \le C^1(T,\varepsilon) + C^2(T,\delta t,\varepsilon).$$

where $C^1(T, \varepsilon)$ and $C^2(T, \delta t, \varepsilon)$ are two constants such that $\lim_{\varepsilon \to 0} C^1(T, \varepsilon) = 0$ and $\lim_{\delta t \to 0} C^2(T, \delta t, \varepsilon) = 0$. Noticing that the proof remains the same when δt is negative, taking the lim sup in δt we get

$$0 \le \limsup_{\delta t \to 0} \int_{x_0}^{+\infty} x^r |f(x, t + \delta t) - f(x, t)| dx \le C^1(T, \varepsilon), \text{ for any } \varepsilon > 0.$$

The proof is completed by taking the limit as ε goes to zero, which yields to the required regularity, $f \in C^0([0, T], L^1([x_0, +\infty), x^r dr))$ for all $r \in [0, 1]$.

We finish this section with a useful estimate for the uniqueness investigation.

Proposition 4 Let T > 0 and $u_1, u_2 \in C_b^0(0, T)$. Let f_1 and f_2 be two mild solutions to (18)–(20), associated, respectively to u_1 and u_2 , with initial data f_1^{in}, f_2^{in} given by formula (26). Then, for any $t \in (0, T)$

$$\int_{x_0}^{+\infty} x |f_1(x,t) - f_2(x,t)| dx \le \int_{x_0}^{+\infty} x |f_1^{in}(x) - f_2^{in}(x)| dx$$
$$- \int_{0}^{t} \int_{x_0}^{+\infty} \mu(x) x |f_1^{in}(x,s) - f_2^{in}(x,s)| dx ds$$
$$+ A_1 \int_{0}^{t} \int_{x_0}^{+\infty} x |f_1(x,s) - f_2(x,s)| dx ds$$
$$+ \int_{0}^{t} (K_{1,2} + C ||f_2(\cdot,s)||_{L^1(xdx)}) |u_1(s) - u_2(s)| ds,$$

where A_1 is given by (22) for u_1 and $K_{1,2}$ is the Lipschitz constant of N on [0, R]with $R = \max(||u_1||_{L^{\infty}(0,T)}, ||u_2||_{L^{\infty}(0,T)})$. Finally C > 0 denotes a constant such that $\rho(x) < Cx$.

Proof This estimation is obtained from a classical argument of approximation. Let $h = f_1 - f_2$ and

$$\int_{x_0}^{+\infty} h(x,t)\varphi(x,t)dx = \int_{x_0}^{+\infty} h^{in}(x)\varphi(x,0)dx$$
$$+ \int_{0}^{t} \int_{x_0}^{+\infty} h(x,s) \left[\frac{\partial}{\partial t}\varphi(x,s) + a_1(s,x) \frac{\partial}{\partial x}\varphi(x,s) - \mu(x)\varphi(x,s) \right] dxds$$
$$+ \int_{0}^{t} \left(N(u_1(s)) - N(u_2(s)) \right) \varphi(x_0,s)ds$$
$$+ \int_{0}^{t} \int_{x_0}^{+\infty} (a_1(s,x) - a_2(s,x)) f_2(x,s) \frac{\partial}{\partial x}\varphi(x,s) dxds.$$

Let h_{ε} be a regularization of h and S_{δ} a regularization of the *Sign* function. Take $\varphi(x, s) = S_{\delta}(h_{\varepsilon}(s, x))g(x)$ with $g \in C_{c}^{\infty}([x_{0}, +\infty))$. Then, letting $\delta \to 0$ and then $\varepsilon \to 0$, we get

$$\int_{x_0}^{+\infty} |h(x,t)|g(x)dx = \int_{x_0}^{+\infty} |h^{in}(x)|g(x)dx$$
$$+ \int_{0}^{t} \int_{x_0}^{+\infty} |h(x,s)| \left[a_1(s,x)\frac{\partial}{\partial x}g(x) - \mu(x)g(x)\right] dxds$$

$$+ \int_{0}^{t} |N(u_{1}(s)) - N(u_{2}(s))| Sign(h_{0}(x_{0}))g(x_{0})ds + \int_{0}^{t} \int_{x_{0}}^{+\infty} (a_{1}(s, x) - a_{2}(s, x)) f_{2}(x, s)Sign(h(s, x))\frac{\partial}{\partial x}g(x)dxds$$

Finally, we approximate the identity function with a regularized function given by $\eta_R \in C_c^{\infty}([x_0, +\infty))$ such that $\eta_R(x) = x$ over (0, R), and then taking the limit $R \to +\infty$ ends the proof.

It is straightforward from Proposition 2 that f defined by (26) is a weak solution and the only one from Proposition 4. Indeed, getting $u_1 = u_2$ and $f_1^0 = f_2^0$ in Proposition 4 leads to the uniqueness. Finally, Proposition 3 provides the continuity in time of the moments with order less or equal to one. This concludes the proof of Proposition 3.

3.3 Proof of the well-posedness

In this section we prove Theorem 2. We first study the operator S defined by (21).

Lemma 4 Consider hypothesis (H2) to (H4). Let u^{in} , p^{in} and b^{in} be nonnegative initial data, and let f^{in} satisfy (H1). Let M > 0 be large enough such that u^{in} , p^{in} , $b^{in} < M/2$ and define

$$X_M = \left\{ (u, p, b) \in \mathcal{C}^0([0, T])^3 : 0 \le u, p, b \le M \right\}$$

where $C^0([0, T])^3$ is equipped with the uniform norm. Then, there exists T > 0 (small enough) such that $S : X_M \mapsto X_M$ is a contraction.

Proof Let *M* be sufficiently large such that $\max(u^{in}, p^{in}, b^{in}) < M/2$, and let T > 0 be small enough such that

$$\begin{aligned} (\gamma_u + \tau M + \sigma + x_0 C_1(M) + C_2(M, T))MT &\leq M/2, \\ (\gamma_p + \tau M)MT &\leq M/2, \\ (\sigma + \delta)MT &\leq M/2, \\ (\lambda_u + \sigma M)T &\leq M/2, \\ (\lambda_p + \sigma M)T &\leq M/2, \\ \tau M^2T &\leq M/2, \end{aligned}$$

where $C_1(M)$ is the Lipschitz constant of N on (0, M) and

$$C_2(M,T) = Ce^{MCT} \left(\|f^{in}\|_{L^1(xdx)} + C_1(M)MT \right),$$
(39)

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where *C* is the constant such that $\rho(x) \leq Cx$, see (27). This assumption ensures that for any $(u, p, b) \in X_M$, then $S(u, p, b) \in X_M$, *i.e*, the solution is bounded by *M* and is nonnegative. It remains to prove that *S* is a contraction. Let (u_1, p_1, b_1) and (u_2, p_2, b_2) belong to X_M . Then

$$\begin{aligned} \|S_{u_{1}} - S_{u_{2}}\|_{\infty} &\leq \gamma_{u} T \|u_{1} - u_{2}\|_{\infty} + \tau T \|u_{1}p_{1} - u_{2}p_{2}\|_{\infty} + \sigma T \|b_{1} - b_{2}\|_{\infty} \\ &+ x_{0} T C_{1}(M) \|u_{1} - u_{2}\|_{\infty} \\ &+ T \sup_{t \in [0,T]} \left| u_{1} \int_{x_{0}}^{+\infty} \rho(x) f_{1}(x,s) dx - u_{2} \int_{x_{0}}^{+\infty} \rho(x) f_{2}(x,s) dx \right|. \end{aligned}$$
(40)

Then,

$$\begin{aligned} \|u_{1}p_{1} - u_{2}p_{2}\|_{\infty} &\leq M \|u_{1} - u_{2}\|_{\infty} + M \|p_{1} - p_{2}\|_{\infty}, \end{aligned}$$
(41)
$$\sup_{t \in [0,T]} \left| u_{1} \int_{x_{0}}^{+\infty} \rho(x) f_{1}(x,s) dx - u_{2} \int_{x_{0}}^{+\infty} \rho(x) f_{2}(x,s) dx \right| \\ &\leq C_{2}(M,T) \|u_{1} - u_{2}\|_{\infty} + CM \sup_{t \in [0,T]} \left| \int_{x_{0}}^{+\infty} x |f_{1}(x,t) - f_{2}(x,t)| dx \right|, \end{aligned}$$
(42)

and from Proposition 4,

$$\sup_{t \in [0,T]} \left| \int_{x_0}^{+\infty} x |f_1(x,t) - f_2(x,t)| dx \right| \le T \left(C_1(M) + CC_2(M,T) \right) \|u_1 - u_2\|_{\infty}.$$
(43)

We get similar bounds for $|S_{p_1} - S_{p_2}|_{\infty}$ and $|S_{b_1} - S_{b_2}|_{\infty}$. We infer that there exists a constant C(M, T) depending only on M and T such that

$$\|(S_{u_1}, S_{p_1}, S_{b_1}) - (S_{u_2}, S_{p_2}, S_{b_2})\|_{\infty} \le C(M, T)T\|(u_1, p_1, b_1) - (u_2, p_2, b_2)\|_{\infty},$$
(44)

with $C(M, T)T \rightarrow 0$, when T goes to 0. Hence, if T is small enough we are able to get C(M, T)T < 1, then S is a contraction.

From Lemma 4, we have a local nonnegative solution on [0, T], which is unique with the solution (u, p, b) bounded by the constant M. The solution satisfies $f \in C^0(0, T; L^1(xdx))$ and $u, p, b \in C^0(0, T)$. Furthermore from (H3), N is continuous and from (H2), $\rho(x) \leq Cx$ where C is a positive constant. Thus $\rho f \in C^0(0, T; L^1(dx))$. We conclude that u, p and b defined in Definition 1 have continuous derivatives.

Now we remark that the solutions satisfy on [0, T]

$$\frac{\mathrm{d}}{\mathrm{d}t} (u+p+2b) = \lambda_u + \lambda_p - \gamma_u u - \gamma_p p - \delta 2b - nN(u) - \frac{1}{\varepsilon} u \int_{x_0}^{+\infty} \rho(x) f(x,t) dx \le \lambda - m(u+p+2b),$$

with $m = \min(\gamma_u, \gamma_p, \delta)$ and $\lambda = \lambda_u + \lambda_p$. Using Gronwall's lemma, the solutions remain bounded at any time by

$$u + p + 2b \le u^{in} + p^{in} + 2b^{in} + \frac{\lambda}{m}.$$
 (45)

From this global bound on u, p and b, we can construct the solution on any interval of time by repetition of the local argument. The proof of the theorem is complete.

4 Perspectives and biological implications

The connection of prions and AD is not fully understood, but recent research suggests that soluble A β oligomers are possible inducers of AD neuropathology. The key element of this hypothesis is the formation of a neurotoxic complex A β -×-PrP^C, which is created by the association of A β oligomers and PrP^C proteins, and not only the progression of β -amyloid plaques by the clustering of A β oligomers.

We believe the model developed and studied here is a step forward in the understanding of the mechanisms underlying AD progression. We have introduced a mathematical model of the evolution of AD based on the hypotheses that A β oligomers exist both as bounded and unbounded to PrP^C proteins, and the agglomeration rate in the formation of β -amyloid plaques depends on the concentrations of the bound and unbound A β oligomers, the concentration of soluble PrP^C, and the size of the β amyloid plaques. Specifically, we have analyzed in detail the existence and uniqueness properties of solutions of the model, as well as the qualitative properties of solution behavior. In specific cases we have quantified the stabilization of the solutions to steady state. In future work, we will explore applications of this model to specific AD laboratory and clinical data. Nevertheless, from this approach we can deduce some suggestions for further research:

- The model suggests a stabilization to steady state for the quantities incorporated into the model. Such phenomena can be very difficult to ascertain in a progressive disease such as AD. Nevertheless, any experimental data quantifying stabilization of AD progression can be valuable in identifying the parameters of the model.
- From an experimental point of view, the investigation of the size distribution of the fibrils is an important consideration. Indeed, we have neglected some phenomena in our study (such as fragmentation-coagulation and depolymerization), and thus it remains to clarify these assumptions. Moreover, in the case of a size-dependent polymerization rate, we would also investigate the character of the polymerization rate from experimental data.

– Finally, we emphasize one further point. One of the main issues in AD is to mitigate progressive memory impairment. Both Aβ-×-PrP^C and β-amyloidplay a key role in the evolution of disease progression, but disappearance of β-amyloid plaques *via* a vaccine does not mitigate neurodegeneration (Holmes et al. 2008). One simple answer would be to increase the degradation rate of Aβ-×-PrP^C by some treatments, which are at present not available. But, as suggested by the model, the polymerization rate could be a key point in the control of disease progression. Indeed, increasing this rate would exhaust the availability of oligomers, and thus reduce the formation of complexes. An important issue remains, namely, what is the best balance between Aβ-×-PrP^C and β-amyloid plaques such that AD patients live the longest without toxicity effects. Perhaps the solution is not to suppress the β-amyloid plaques, but rather control their progression. The question is open and a deeper analysis of the model, together with biological data, would provide understanding in this direction.

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Appendix A: Characteristic polynomials of the linearized ODE system

Here we give the coefficient a_i , i = 1, ..., 4 for the characteristic polynomial of the linearized system in Theorem 1:

$$a_{1} = \left(\mu + \gamma_{u} + \tau \frac{\lambda_{p}}{\tau^{*}u_{\infty} + \gamma_{p}} + \alpha n^{2}u_{\infty}^{n-1} + \rho \frac{\alpha}{\mu}u_{\infty}^{n} + \gamma_{p} + \tau u_{\infty} + \sigma + \delta\right),$$

$$a_{2} = \left(\mu + \gamma_{u} + \alpha n^{2}u_{\infty}^{n-1} + \rho \frac{\alpha}{\mu}u_{\infty}^{n}\right)(\gamma_{p} + \tau u_{\infty} + \sigma + \delta) + \gamma_{p}\sigma + (\gamma_{p} + \tau u_{\infty})\delta$$

$$+ \mu \left(\gamma_{u} + \tau \frac{\lambda_{p}}{\tau^{*}u_{\infty} + \gamma_{p}} + \alpha n^{2}u_{\infty}^{n-1} + \rho \frac{\alpha}{\mu}u_{\infty}^{n}\right) + \rho\alpha nu_{\infty}^{n} + \tau (\gamma_{p} + \delta)\frac{\lambda_{p}}{\tau^{*}u_{\infty} + \gamma_{p}},$$

$$a_{3} = \left(\mu + \gamma_{u} + \alpha n^{2}u_{\infty}^{n-1} + \rho \frac{\alpha}{\mu}u_{\infty}^{n}\right)(\gamma_{p}\sigma + (\gamma_{p} + \tau u_{\infty})\delta) + (\gamma_{p}\delta + (\gamma_{p} + \delta)\mu)\tau \frac{\lambda_{p}}{\tau^{*}u_{\infty} + \gamma_{p}}$$

$$+ \left\{\mu \left(\gamma_{u} + \alpha n^{2}u_{\infty}^{n-1} + \rho \frac{\alpha}{\mu}u_{\infty}^{n}\right) + \rho\alpha nu_{\infty}^{n}\right\}(\gamma_{p} + \tau u_{\infty} + \sigma + \delta),$$

$$a_{4} = \mu\gamma_{p}\delta\tau \frac{\lambda_{p}}{\tau^{*}u_{\infty} + \gamma_{p}} + \left\{\mu \left(\gamma_{u} + \alpha n^{2}u_{\infty}^{n-1} + \rho \frac{\alpha}{\mu}u_{\infty}^{n}\right) + \rho\alpha nu_{\infty}^{n}\right\}(\gamma_{p}\sigma + (\gamma_{p} + \tau u_{\infty})\delta).$$

Appendix B: Lyapunov functional

Here we detail a Lyapunov function Φ which is the key ingredient to prove global stability of system (8–11) in Proposition 2. This function appears to be a bit tricky,

but determining it rest upon the backward method described for instance Chapter 4, p. 120, in the book by Khalil (1996). It consists in investigate an expression of the derivative Φ' and then going back to chose the parameters Φ such as Φ' is negative definite. After tedious calculus, a Liapunov function Φ for system (8–11) is given by

$$\begin{split} \Phi &= \frac{1}{2} \left(\frac{2\gamma_p}{\delta} \right) s_1 \theta_1^2 + \frac{1}{2} \left(1 + 2 \frac{\delta + \gamma_u + \rho (A_\infty + \theta_1)}{\sigma} \right) \theta_2^2 + \frac{1}{2} \left(\frac{2\gamma_p}{\delta} \right) \theta_3^2 \\ &+ \frac{1}{2} \left(\frac{\sigma}{\gamma_p} \right) \theta_4^2 + \left(\frac{\rho p_\infty}{\gamma_u + \rho A_\infty + \mu} \right) \theta_1 \theta_2 + \theta_1 \theta_3 \\ &+ \left(\frac{\rho p_\infty}{\gamma_u + \rho A_\infty + \mu} + 1 + \frac{\rho}{\tau} \right) \theta_1 \theta_4 + \theta_2 \theta_3 + 2\theta_2 \theta_4 + \left(\frac{2\gamma_p}{\delta} \right) \theta_3 \theta_4, \end{split}$$

where $\theta_1 = A - A_\infty$, $\theta_2 = u - u_\infty$, $\theta_3 = p - p_\infty$, $\theta_4 = b - b_\infty$, with $s_1 = \max(T_1, T_2)$ such that

$$T_{1} = \frac{\rho^{2} \delta u_{\infty}^{2} \left(1 + 2\frac{1+\delta}{\sigma}\right)}{8\mu\gamma_{p}} + \frac{(\gamma_{p} + \mu)^{2} \left(\frac{\delta}{2\gamma_{p}}\right)^{2}}{4\gamma_{p}\mu} + \frac{\left[\left(\delta + \mu\right) \left(\frac{\rho p_{\infty}}{\gamma_{u} + \rho A_{\infty} + \mu} + 1\right) + \left(\sigma + \delta + \mu\right)\frac{\rho}{\tau} + 2\rho u_{\infty}\right]^{2}}{8\mu\sigma},$$

and $T_2 = \Gamma \left(\frac{\delta}{2\gamma_p}\right)^2 T'_2$ with

$$T_{2}' = \left(\frac{\rho p_{\infty}}{\gamma_{u} + \rho A_{\infty} + \mu}\right)^{2} \left\{ \frac{2\sigma + \delta}{2\gamma_{p}} + \left(\frac{\delta}{2\gamma_{p}}\Gamma\right)^{-1} \left(\frac{1}{1 + 2\frac{\delta + \gamma_{u}}{\sigma}}\right) \right\} \\ + \frac{\rho p_{\infty}}{\gamma_{u} + \rho A_{\infty} + \mu} \left\{ 2 + 4\frac{\rho}{\tau}\frac{\delta + \gamma_{u}}{\sigma} \right\} + \frac{\delta}{2\gamma_{p}} \left\{ \frac{\rho}{\tau} \left(2 + \frac{\rho}{\tau}\right) + \frac{\sigma + 2(\delta + \gamma_{u})}{\gamma_{p}} \right\} \\ + \left(1 + 2\frac{\delta + \gamma_{u}}{\sigma}\right) \left\{ \frac{\rho}{\tau} \left(1 + \frac{\rho}{\tau}\right) + \frac{\delta}{2\gamma_{p}}\frac{\sigma}{\gamma_{p}} - 1 \right\}$$

and

$$\Gamma = \frac{1}{\left(1 + 2\frac{\delta + \gamma_u}{\sigma} - \frac{\delta}{2\gamma_p}\right)\left(\frac{\delta}{2\gamma_p}\frac{\sigma}{\gamma_p} - 1\right)}.$$

We remark that $T_1 > 0$ so that $s_1 > 0$, and then we deduce that the Lyapunov function Φ is positive when condition $\left(1 + 2\frac{\delta + \gamma_u}{\sigma}\right) > \frac{\delta}{2\gamma_p} > \frac{\gamma_p}{\sigma}$ holds true. In such case, its derivative along the solutions of system (8–11) is given by

$$\begin{split} \Phi' &= -\left(\mu s_1 + \rho u \frac{\delta}{2\gamma_p} \cdot \frac{\rho p_\infty}{\gamma_u + \rho A_\infty + \mu}\right) \theta_1^2 - \rho u_\infty \frac{\delta}{2\gamma_p} \left(1 + 2 \frac{\gamma_u + \rho (A_\infty + \theta_1) + \delta}{\sigma}\right) \theta_1 \theta_2 \\ &- \frac{\delta}{2\gamma_p} \left(\frac{2(\gamma_u + \rho (A_\infty + \theta_1) + \tau p)(\gamma_u + \rho (A_\infty + \theta_1) + \delta)}{\sigma} + \gamma_u + \rho (A_\infty + \theta_1)\right) \theta_2^2 \\ &- \frac{\delta}{2\gamma_p} \left((\delta + \mu) \left(\frac{\rho p_\infty}{\gamma_u + \rho A_\infty + \mu} + 1\right) + (\sigma + \delta + \mu) \frac{\rho}{\tau} + 2\rho u_\infty\right) \theta_1 \theta_4 \\ &- \left(\frac{\delta \tau u}{2\gamma_p} + \gamma_p\right) \theta_3^2 - \delta \left(\frac{\sigma}{\gamma_p} \frac{\delta}{2\gamma_p}\right) \theta_4^2 - \frac{\delta}{2\gamma_p} (\gamma_p + \mu) \theta_1 \theta_3. \end{split}$$

and remains nonpositive. Furthermore, $\Phi' = 0$ if and only if $\theta_1 = \theta_2 = \theta_3 = \theta_4 = 0$. The conclusion holds by the LaSalle Invariance Principle LaSalle (1976).

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