

Numerical analysis of a one-dimensional elastodynamic model of dry friction and unilateral contact

Y. Renard *

SIGMAS Project, Laboratoire de Modélisation et Calcul LMC-IMAG, BP 53X, F-38041 Grenoble, France

Abstract

This paper deals with a numerical analysis of a one-dimensional dynamic purely elastic (i.e. hyperbolic) model with dry friction. Since we consider a Coulomb friction law with a slip velocity dependent coefficient, generally, the problem has more than one solution. A mass perturbation approach is developed to regain the uniqueness and to perform the numerical analysis. This study can be viewed as a first step in the numerical analysis of more elaborated dynamic purely elastic problems with dry friction. © 2001 Elsevier Science B.V. All rights reserved.

1. Introduction

Friction laws with a slip velocity dependent coefficient were introduced to modelize the stick-slip phenomenon, which is the appearance of self-sustained vibration in mechanical systems submitted to dry friction. Even though there is not a universally accepted model of this phenomenon, the dynamic aspect seems to be essential in the behavior of such systems. Some mathematical and numerical results tend to prove that multi-dimensional systems submitted to dry friction and unilateral contact can develop instabilities even if the simplest Coulomb law with a constant coefficient is chosen (see [3,15,11,16]). Because of the results on systems with finite number of degree of freedom, the Coulomb law with a slip velocity dependent coefficient is still considered as a good candidate for the modelization of stick-slip motion of elastic structures.

The model described above is a one-dimensional purely elastic model with dry friction and unilateral contact. It is well known that there are major mathematical difficulties to treat such dynamical model, due to the bad regularity of the solutions (discontinuities in velocity). This work is a continuation of a paper [19] which treats the modeling part and also of the work of Ionescu and Paumier [8]. This one-dimensional model, because of its simplicity, allows to enhance some fundamental properties of the dynamic behavior of elastic structure under dry friction and unilateral contact. The purpose here is to perform a numerical analysis of the purely elastic case, with a mass perturbation approach. This mass perturbation approach is also studied in [19] where it is proven, under certain restrictions, that the sequence of solutions of the perturbed problem converges toward a particular solution of the non-perturbed problem. This particular solution can be related to the selection made by the perfect delay criterion (see [8,10]).

* Present address: Département de Génie Mathématique, INSA de Toulouse, 31077 Toulouse Cedex 4, France.

E-mail address: yves.Renard@gmm.insa-tlse.fr (Y. Renard).

2. The one-dimensional elastic model with dry friction

The model we deal with in this paper is a slight extension of the one introduced in [8,19], and was also studied in [17] and fully described in [18]. This model represents the dynamic motion of an elastic slab which slides with friction on a moving rigid foundation (see Fig. 1). The slab is assumed to be linearly elastic with constant Lamé coefficients λ and G , density ρ and height H .

With convenient initial data and lateral boundary conditions, it is possible to consider motions of the slab which only depend on the vertical coordinate x_3 . So, if we denote by $\underline{u}(t, x_3) = (u_1(t, x_3), u_2(t, x_3), u_3(t, x_3))$ the displacement of the slab, the evolution of the vertical displacement $u_3(t, x_3)$ is described by the equation

$$\partial_t^2 u_3(t, x_3) = c_2 \partial_{x_3 x_3}^2 u_3(t, x_3) \quad \text{in }]0, T] \times]0, H[, \quad (1)$$

where $c_2 = \sqrt{(\lambda + 2G)/\rho}$ is the velocity of longitudinal waves. The slab is assumed to be fixed on its top

$$u_3(t, H) = -D \quad \text{with } D > 0 \quad \forall t \in]0, T], \quad (2)$$

and it is submitted to a unilateral contact condition on $x_3 = 0$:

$$\partial_{x_3} u_3(t, 0) \leq 0, \quad -u_3(t, 0) \leq 0, \quad \partial_{x_3} u_3(t, 0) u_3(t, 0) = 0 \quad \forall t \in]0, T]. \quad (3)$$

It is convenient to express this later condition by the inclusion

$$(\lambda + 2G) \partial_{x_3} u_3(t, 0) \in -J_N(-u_3(t, 0)) \quad \forall t \in]0, T], \quad (4)$$

where the multi-valued map

$$J_N(x) = \begin{cases} \{0\} & \text{if } x < 0, \\ [0, +\infty[& \text{if } x = 0, \end{cases}$$

is the sub-differential of the convex indicator function of the interval $]-\infty, 0]$.

The evolution of the horizontal displacement $u_T(t, x_3) = (u_1(t, x_3), u_2(t, x_3))$ is described by the equation

$$\partial_t^2 u_T(t, x_3) = c_1 \partial_{x_3 x_3}^2 u_T(t, x_3), \quad \text{in }]0, T] \times]0, H[, \quad (5)$$

where $c_1 = \sqrt{G/\rho}$ is the velocity of transversal waves. On $x_3 = H$ it is assumed that the slab is fixed

$$u_T(t, H) = 0 \quad \forall t \in]0, T]. \quad (6)$$

The friction is modeled by a Coulomb law of friction with a slip velocity dependent coefficient on the contact boundary $x_3 = 0$:

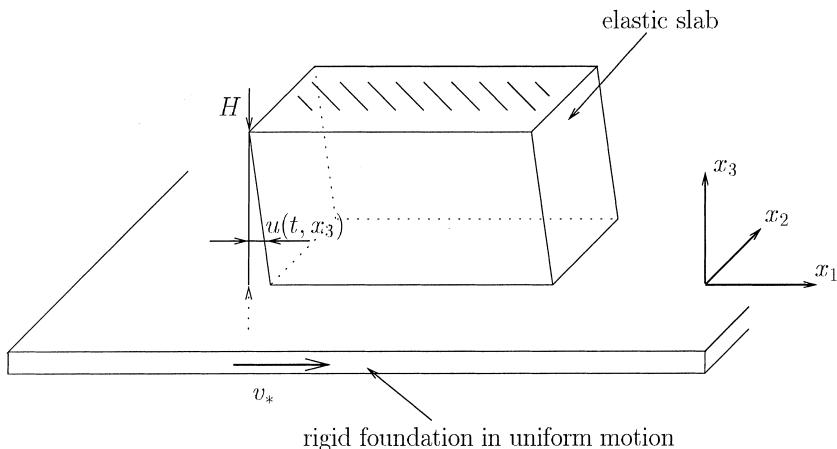


Fig. 1. Elastic slab in sliding with dry friction on a rigid foundation which is in uniform motion v_* along the x_1 axis.

$\forall t \in]0, T]$

$$\text{if } \partial_t u_T(t, 0) = v_* \text{ then } |G\partial_{x_3} u_T(t, 0)| \leq -(\lambda + 2G)\partial_{x_3} u_3(t, 0)\mu(0), \quad (7)$$

$$\text{if } \partial_t u_T(t, 0) \neq v_* \text{ then } \begin{cases} |G\partial_{x_3} u_T(t, 0)| = -(\lambda + 2G)\partial_{x_3} u_3(t, 0)\mu(\|\partial_t u_T(t, 0) - v_*\|), \\ \exists \lambda \geq 0, \partial_t u_T(t, 0) - v_* = \lambda \partial_{x_3} u_T(t, 0), \end{cases} \quad (8)$$

where $\mu : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ is the friction coefficient, and v_* the constant planar velocity of the rigid foundation. It is easy to see that this friction condition is equivalent to the following inclusion:

$\forall t \in]0, T]$,

$$G\partial_{x_3} u_T(t, 0) \in -(\lambda + 2G)\partial_{x_3} u_3(t, 0)\mu(\|\partial_t u_T(t, 0) - v_*\|)\text{Dir}(\partial_t u_T(t, 0) - v_*), \quad (9)$$

where the multi-valued map

$$\text{Dir}(v) = \left\{ \frac{v}{\|v\|} \right\} \text{ if } v \neq 0, \quad \text{Dir}(0) = \overline{B}(0, 1),$$

is defined on \mathbb{R}^2 , and is the sub-differential of the convex map $v \mapsto \|v\|$. The problem is completed by initial conditions

$$\underline{u}(0, x_3) = \underline{u}^0(x_3), \quad \partial_t \underline{u}(0, x_3) = \underline{u}^1(x_3) \quad \forall x_3 \in [0, H]. \quad (10)$$

Finally, introducing the matrix

$$\mathbf{C} = \begin{pmatrix} c_1 & 0 & 0 \\ 0 & c_1 & 0 \\ 0 & 0 & c_2 \end{pmatrix},$$

and with the previous notations, the problem can be written as follows:

$$(P1) \quad \partial_{tt}^2 \underline{u}(t, x_3) = \mathbf{C} \partial_{x_3 x_3}^2 \underline{u}(t, x_3) \quad \text{in }]0, T] \times]0, H[, \quad (11)$$

$$\underline{u}(t, H) = (0, 0, -D) \quad \text{with } D > 0 \quad \forall t \in]0, T], \quad (12)$$

$$\underline{u}(0, x_3) = \underline{u}^0(x_3), \quad \partial_t \underline{u}(0, x_3) = \underline{u}^1(x_3) \quad \forall x_3 \in [0, H], \quad (13)$$

$$(\lambda + 2G)\partial_{x_3} u_3(t, 0) \in -J_N(-u_3(t, 0)) \quad \forall t \in]0, T], \quad (14)$$

$$G\partial_{x_3} u_T(t, 0) \in -(\lambda + 2G)\partial_{x_3} u_3(t, 0)\mu(\|\partial_t u_T(t, 0) - v_*\|)\text{Dir}(\partial_t u_T(t, 0) - v_*) \quad \forall t \in]0, T], \quad (15)$$

Due to the bad behavior of both the unilateral contact and the friction conditions, it cannot be expected to have any solution of the problem in a classical sense (i.e. twice continuously differentiable). It is easy to see that, most of the time, the velocity $\partial_t \underline{u}$ is not continuous (see [8]). One can also notice that it is possible to solve independently the problem in vertical displacement u_3 .

Friction laws such as Coulomb law with a sliding velocity dependent coefficient have a long history. Rabinowicz already studies such dependencies in 1958 in [14]. Important experimental works on the measurement of this dependency has been made (see for instance [2,5,7,10,21]). Most of these experiments have been made on the steady state because transient states are extremely difficult to measure.

Very often, a decreasing coefficient of friction is used in order to modelize stick-slip phenomenon. The value of the coefficient for a sliding velocity equal to zero is called the static coefficient μ_s , and the asymptotic value for a huge sliding velocity is called the dynamic coefficient of friction μ_d .

3. On the existence and uniqueness

For convenience and from now on, we assume $T < (H/c_2)$. This is not strictly necessary, but the analysis will be greatly simplified, and this is not a real restriction, since all the following results can be extended for $T > (H/c_2)$ in studying the problem successively on intervals

$$\left[0, \left(\frac{H}{c_2}\right)\right], \quad \left[\frac{H}{c_2}, 2\frac{H}{c_2}\right], \quad \left[2\frac{H}{c_2}, 3\frac{H}{c_2}\right], \dots$$

Let us introduce the quantities

$$b_+(t, x_3) = \partial_t \underline{u}(t, x_3) + \mathbf{C} \partial_3 \underline{u}(t, x_3), \quad (16)$$

$$b_-(t, x_3) = \partial_t \underline{u}(t, x_3) + \mathbf{C} \partial_3 \underline{u}(t, x_3), \quad (17)$$

which are propagated along the characteristic lines of Eq. (11), i.e. that for a convenient δt and for ε is + or – one has

$$b_{1\varepsilon}(t + \delta t, x_3) = b_{1\varepsilon}(t, x_3 + \varepsilon c_1 \delta t), \quad (18)$$

$$b_{2\varepsilon}(t + \delta t, x_3) = b_{2\varepsilon}(t, x_3 + \varepsilon c_1 \delta t), \quad (19)$$

$$b_{3\varepsilon}(t + \delta t, x_3) = b_{3\varepsilon}(t, x_3 + \varepsilon c_2 \delta t), \quad (20)$$

where $b_{1\varepsilon}$, $b_{2\varepsilon}$ and $b_{3\varepsilon}$ are the components of b_ε . Since we assume $T < H/c_2$, and thanks to these equalities, if we set

$$u_N(t) = -u_3(t, 0), \quad v_s = \partial_t u_T(t, 0) - v_*, \quad b_{T\varepsilon} = \begin{pmatrix} b_{1\varepsilon} \\ b_{2\varepsilon} \end{pmatrix}, \quad (21)$$

and if we denote by

$$S(t) = -\frac{(\lambda + 2G)}{c_2} \left(b_{3+}(t, 0) + \frac{d}{dt} u_N(t) \right), \quad (22)$$

the contact pressure which is always positive, then Problem (P1) becomes

(P2) Find u_N and $v_s(t)$ such that :

$$\frac{d}{dt} u_N(t) \in -b_{3+}(t, 0) - \frac{c_2}{\lambda + 2G} J_N(u_N(t)), \quad (23)$$

$$u_N(0) = u_{N0} = u_3^0(0, 0), \quad (24)$$

$$\frac{G}{c_1} (b_{T+} - v_*) \in \frac{G}{c_1} v_s(t) + S(t) \mu(\|v_s(t)\|) \text{Dir}(v_s(t)). \quad (25)$$

One can notice that those two equations are not of the same kind. Eqs. (23) and (24) form a first-order differential inclusion Cauchy problem in the variable u_N , and Eq. (25) is a scalar inclusion in variable v_s . Problem (23), (24) is still independent of Eq. (25) and the following result holds:

Proposition 1. If $b_{3+}(\cdot, 0) \in L^1(0, T)$ and if the initial condition satisfies $u_{N0} \leq 0$, the following Cauchy problem

$$\begin{cases} \frac{d}{dt} u_N(t) \in -b_{3+}(t, 0) - \frac{c_2}{\lambda + 2G} J_N(u_N(t)), & \text{a.e. in } [0, T], \\ u_N(0) = u_{N0}, \end{cases} \quad (26)$$

has a unique absolutely continuous solution, and this solution satisfies $u_N(t) \leq 0$ on $[0, T]$.

The proof of this result comes directly from the result of Deimling [4] on differential inclusions, and the fact that the multi-valued map

$$(t, x) \mapsto -b_{3+}(t, 0) - \frac{c_2}{\lambda + 2G} J_N(x), \quad (27)$$

has convex values, is measurable, upper semi-continuous and one-sided Lipschitz. \square

Unfortunately, inclusion (25) has generally not a unique solution. We can only prove the following result:

Proposition 2. *If the friction coefficient μ is a Lipschitz continuous and bounded positive function of the sliding velocity v_s , then Inclusion (25) has at least one solution. Moreover, if the friction coefficient μ is an increasing function of v_s this solution is unique.*

The proof of this result is a simple study of the shape of the right-hand side of Inclusion (25). \square

In fact, when the friction coefficient has at least one decreasing branch, it is always possible to exhibit a case where there is more than one solution to Inclusion (25). See [8,19] for such examples. Moreover, this is the interesting case, since it seems to be necessary to use a coefficient with a decreasing branch to modelize well stick-slip phenomena.

4. Mass perturbation and perfect delay criterion

One of the major problems is thus the non-uniqueness of the solution of Problem (P1). In fact, certain considerations of mechanical order (see [19]) allow to say that most of the solutions of Inclusion (25) are not acceptable. Authors like [8,10] proposed a criterion in order to choose the convenient solution of the problem. This criterion is a dynamic extension of the perfect delay criterion which comes from the catastrophe theory and which can be formulated as follows:

the system only jump when it has no other choice. (28)

This criterion is not sufficiently accurate to always select a unique solution (it does in experience of pure loading or pure unloading, see [8]).

Another approach to this problem is to introduce a viscous term in the elastic law and to make this term to vanish. In this framework, it is possible to prove that the visco-elastic version of the problem has a unique solution. But at this moment there is no theoretical results of convergence when the viscous term vanishes.

The approach we introduced in a work [19], is to consider a mass perturbation of the friction condition. The classical friction condition

$$G\partial_{x_3}u_T(t,0) \in S(t)\mu(\|\partial_tu_T(t,0) - v_*\|)\text{Dir}(\partial_tu_T(t,0) - v_*), \quad (29)$$

is replaced by the following condition

$$\varepsilon\partial_u^2u_T + G\partial_{x_3}u_T(t,0) \in S(t)\mu(\|\partial_tu_T(t,0) - v_*\|)\text{Dir}(\partial_tu_T(t,0) - v_*), \quad (30)$$

where ε is the perturbation parameter. The additional term $\varepsilon\partial_u^2u_T$ can be assimilated as a surface mass on the contact boundary. This kind of perturbations, similarly to existing friction laws as Dieterich–Ruina laws (see [20]), make a differential operator to appear on the contact boundary. Further development of such boundary perturbations in upper dimensions are presented in [13].

With this mass perturbation, the Problem (P2) becomes

(P3) Find u_N and $v_s(t)$ such that:

$$\frac{d}{dt}u_N(t) \in -b_{3+}(t,0) - \frac{c_2}{\lambda + 2G}J_N(u_N(t)), \quad (31)$$

$$u_N(0) = u_N^0 = u_3^0(0,0), \quad (32)$$

$$\varepsilon\frac{d}{dt}v_s(t) \in \frac{G}{c_1}(b_{T+} + v_s(t) - v_*) + S(t)\mu(\|v_s(t)\|)\text{Dir}(v_s(t)), \quad (33)$$

$$v_s(0) = v_s^0 = u_T^1(0,0). \quad (34)$$

The first result is that Problem (P3) has a unique solution.

Proposition 3. If $b_{T+} \in (L^1(0, T))^2$ and if the contact pressure S is a positive function and belongs to $L^1(0, T)$ (this is the case if S comes from the solution of (31), (32) and if conditions of Proposition (1) hold), then the following Cauchy problem:

$$\begin{cases} \varepsilon \frac{d}{dt} v_s(t) \in \frac{G}{c_1} (b_{T+} + v_s(t) - v_*) + S(t) \mu(\|v_s(t)\|) \text{Dir}(v_s(t)), \\ v_s(0) = v_s^0, \end{cases} \quad (35)$$

has a unique absolutely continuous solution.

Similarly as for Proposition 1, this result comes from the fact that the multi-valued map

$$(t, v) \mapsto \frac{G}{c_1} (b_{T+} + v - v_*) + S(t) \mu(\|v\|) \text{Dir}(v), \quad (36)$$

has convex values, is upper semi-continuous and one-sided Lipschitz. \square

In a more restrictive framework, in [19], we proved that when ε goes to zero the solution of the perturbed problem converges toward a particular solution of the initial problem. The connection between this and the perfect delay criterion is that if this criterion selects a unique solution, then this is also the limit of the solutions of the perturbed problems.

5. Regularized unilateral contact condition

For the convergence proof of Section 6, we need the solution to have a certain regularity. The regularity of u_N is not sufficient. One of the most usual way to regularize the unilateral contact condition is to consider a continuous approximation of the map J_N . It is the case, for example, for the well-known normal compliance law developed by Oden and Martins in [12]. Here we consider the following simple piecewise linear approximation J_N^η of J_N :

$$J_N^\eta(x) = \begin{cases} 0 & \text{if } x < 0, \\ \frac{x}{\eta} & \text{if } x \geq 0, \end{cases} \quad (37)$$

The solutions of the differential equation

$$\frac{d}{dt} u_N(t) = -b_{3+}(t, 0) - \frac{c_2}{\lambda + 2G} J_N^\eta(u_N(t)), \quad (38)$$

are such that the time derivative of u_N is Lipschitz continuous since $b_{3+}(t, 0)$ is. This regularization corresponds to a certain interface elasticity.

6. Numerical analysis

In this section we present a numerical scheme for the perturbed problem. Once again, the results will be established for $T < H/c_2$ even though they can be easily extended for $T > H/c_2$ and of course even though the numerical scheme has not himself such a restriction. The scheme is a coupling of a Lax–Wendroff scheme for the interior of the slab and of implicit Euler schemes for the equations on the boundary. The implicit Euler scheme is chosen for the simplicity of the proof. More accurate schemes can be used (in [18] we use a theta method, with a parameter 1/2, which is also an two points Adams–Moulton scheme, and gives a better numerical approximation). The convergence is obtained for the regularized unilateral contact condition. This is an open problem to know if it should be possible to work with the normal contact condition. This analysis can be viewed as a first step to improve numerical analysis results for purely elastic friction problems with a complex friction law.

6.1. Description of the scheme

The slab is discretized with a constant step $\Delta z = H/N_z$ and a constant time step $\Delta t = T/N_t$. The principal variables are $\underline{u}_{,t} = \partial_t \underline{u}$ and $\underline{u}_{,3} = \partial_{x_3} \underline{u}$. Let us denote by

$$\underline{u}^{i,n} = \begin{pmatrix} u_1^{i,n} \\ u_2^{i,n} \\ u_3^{i,n} \end{pmatrix}, \quad \underline{u}_{,t}^{i,n} = \begin{pmatrix} u_{1,t}^{i,n} \\ u_{2,t}^{i,n} \\ u_{3,t}^{i,n} \end{pmatrix}, \quad \text{and} \quad \underline{u}_{,3}^{i,n} = \begin{pmatrix} u_{1,3}^{i,n} \\ u_{2,3}^{i,n} \\ u_{3,3}^{i,n} \end{pmatrix},$$

the approximated values of, respectively \underline{u} , $\underline{u}_{,t}$ and $\underline{u}_{,3}$ at the time $t^n = n\Delta t$ and at $x_3^i = i\Delta z$. We also use the following notations for the tangential part:

$$\underline{u}_T^{0,n} = \begin{pmatrix} u_1^{0,n} \\ u_2^{0,n} \end{pmatrix}, \quad \underline{u}_{T,t}^{0,n} = \begin{pmatrix} u_{1,t}^{0,n} \\ u_{2,t}^{0,n} \end{pmatrix} \quad \text{and} \quad \underline{u}_{T,3}^{0,n} = \begin{pmatrix} u_{1,3}^{0,n} \\ u_{2,3}^{0,n} \end{pmatrix}.$$

With the matrix

$$\mathbf{D} = \frac{\Delta t}{\Delta z} C = \begin{pmatrix} c_1 \frac{\Delta t}{\Delta z} & 0 & 0 \\ 0 & c_1 \frac{\Delta t}{\Delta z} & 0 \\ 0 & 0 & c_2 \frac{\Delta t}{\Delta z} \end{pmatrix} = \begin{pmatrix} d_1 & 0 & 0 \\ 0 & d_1 & 0 \\ 0 & 0 & d_2 \end{pmatrix},$$

The Lax–Wendroff scheme on the interior of the slab, for $1 \leq i \leq N_z - 1$, reads

$$\underline{u}_{,t}^{i,n+1} = \underline{u}_{,t}^{i,n} + \frac{\mathbf{D}}{2} \left(\underline{u}_{,t}^{i-1,n} - 2\underline{u}_{,t}^{i,n} + \underline{u}_{,t}^{i+1,n} \right) + \frac{\mathbf{DC}}{2} \left(\underline{u}_{,3}^{i+1,n} - \underline{u}_{,3}^{i-1,n} \right), \quad (39)$$

$$\underline{u}_{,3}^{i,n+1} = \underline{u}_{,3}^{i,n} + \frac{\mathbf{D}}{2} \left(\underline{u}_{,3}^{i-1,n} - 2\underline{u}_{,3}^{i,n} + \underline{u}_{,3}^{i+1,n} \right) + \frac{\mathbf{DC}^{-1}}{2} \left(\underline{u}_{,t}^{i+1,n} - \underline{u}_{,t}^{i-1,n} \right). \quad (40)$$

With

$$\underline{b}_+^{i,n} = \underline{u}_{,t}^{i,n} + \mathbf{C} \underline{u}_{,3}^{i,n} \quad \text{and} \quad \underline{b}_-^{i,n} = \underline{u}_{,t}^{i,n} - \mathbf{C} \underline{u}_{,3}^{i,n}, \quad (41)$$

this scheme reads also

$$\underline{b}_+^{i,n+1} = \underline{b}_+^{i,n} + \mathbf{D}(\underline{b}_+^{i+1,n} - \underline{b}_+^{i,n}), \quad 0 \leq i \leq N_z - 1, \quad (42)$$

$$\underline{b}_-^{i,n+1} = \underline{b}_-^{i,n} + \mathbf{D}(\underline{b}_-^{i-1,n} - \underline{b}_-^{i,n}), \quad 1 \leq i \leq N_z, \quad (43)$$

The approximation of the Dirichlet condition on $z = H$ is

$$\underline{u}_{,3}^{N_z,n+1} = \underline{u}_{,3}^{N_z,n} + \mathbf{D}(\underline{u}_{,3}^{N_z-1,n} - \underline{u}_{,3}^{N_z,n}) - \mathbf{DC}^{-1} \underline{u}_{,t}^{N_z-1,n}, \quad (44)$$

$$\underline{u}_{,t}^{N_z,n+1} = 0. \quad (45)$$

or

$$\underline{b}_+^{N_z,n+1} = -\underline{b}_-^{N_z,n+1}. \quad (46)$$

The unilateral contact condition, which corresponds to the differential inclusion

$$\frac{d}{dt} u_3(t, 0) \in b_{3+}(t, 0) + J_{\mathcal{N}}(-u_3(t, 0)),$$

for the nonregularized case or to the differential equation

$$\frac{d}{dt} u_3(t, 0) = b_{3+}(t, 0) + \frac{c_2}{\lambda + 2G} J_{\mathcal{N}}^{\eta}(-u_3(t, 0)),$$

for the regularized case, is approximated with an implicit Euler scheme. This gives for the non regularized case

$$u_3^{0,n+1} = (u_3^{0,n} + \Delta t b_{3+}^{0,n+1})_+, \quad (47)$$

where $(x)_+ = \max(0, x)$ is the positive part of x . The value of $u_{3,t}^{0,n+1}$ is approximated by

$$u_{3,t}^{0,n+1} = \begin{cases} b_{3+}^{0,n+1} & \text{if } u_3^{0,n+1} \neq 0, \\ 0 & \text{otherwise,} \end{cases} \quad (48)$$

For the regularized case, the implicit Euler scheme gives

$$u_3^{0,n+1} = (u_3^{0,n} + \Delta t b_{3+}^{0,n+1})_+ + \frac{1}{1 + \frac{\Delta t}{\tilde{\eta}}} (u_3^{0,n} + \Delta t b_{3+}^{0,n+1})_-, \quad (49)$$

$$u_{3,t}^{0,n+1} = b_{3+}^{0,n+1} + J_N^{\tilde{\eta}}(-u_3^{0,n+1}), \quad (50)$$

where $\tilde{\eta} = \frac{\lambda+2G}{c_2} \eta$.

The contact pressure is thus given by

$$S^{n+1} = -\frac{\lambda+2G}{c_2} (b_{3+}^{0,n+1} - u_{3,t}^{0,n+1}). \quad (51)$$

Those schemes are such that $S^{n+1} \geq 0$. Moreover, for the nonregularized case, one has

$$u_3^{0,n+1} \geq 0, \quad S^{n+1} u_3^{0,n+1} = 0,$$

similarly as the continuous case.

Now, concerning the perturbed friction condition, we have to approximate the following differential inclusion:

$$\varepsilon \frac{d}{dt} \underline{u}_{T,t} \in \frac{G}{c_1} (b_{T+}(t) - \underline{u}_{T,t}) - S(t) \mu(\|\underline{u}_{T,t} - \underline{V}_e(t)\|) \text{Dir}(\underline{u}_{T,t} - \underline{V}_e(t)),$$

The implicit Euler scheme reads

$$\underline{u}_{T,t}^{0,n+1} \in \underline{u}_{T,t}^{0,n} + \frac{\Delta t}{\varepsilon} \left(\frac{G}{c_1} (b_{T+}^{0,n+1} - \underline{u}_{T,t}^{0,n+1}) - S^{n+1} \mu(\|\underline{u}_{T,t}^{0,n+1} - \underline{V}_e(t^{n+1})\|) \text{Dir}(\underline{u}_{T,t}^{0,n+1} - \underline{V}_e(t^{n+1})) \right), \quad (52)$$

6.2. Properties of the Lax–Wendroff scheme

It is well known that the Lax–Wendroff scheme is not stable if the Courant–Friedrichs–Levi condition $d_2 = \frac{\Delta t}{\Delta z} c_2 \leq 1$ does not hold. Thus we assume that the matrix \mathbf{D} has all its coefficients between 0 and 1.

The Lax–Wendroff scheme has important properties concerning the L^∞ stability. If we consider the following quantities:

$$B_{\max}^n = \max_{0 \leq i \leq N_z} (\|\underline{b}_+^{i,n}\|, \|\underline{b}_-^{i,n}\|), \quad (53)$$

$$dB_{\max}^n = \max_{0 \leq i \leq N_z-1} (\|\underline{b}_+^{i+1,n} - \underline{b}_+^{i,n}\|, \|\underline{b}_-^{i+1,n} - \underline{b}_-^{i,n}\|), \quad (54)$$

then from Eqs. (42) and (43) it comes

$$\|\underline{b}_+^{i,n+1}\| = \|\underline{b}_+^{i,n} + \mathbf{D}(\underline{b}_+^{i+1,n} - \underline{b}_+^{i,n})\| \leq B_{\max}^n, \quad 0 \leq i \leq N_z - 1,$$

$$\|\underline{b}_-^{i,n+1}\| = \|\underline{b}_-^{i,n} + \mathbf{D}(\underline{b}_-^{i-1,n} - \underline{b}_-^{i,n})\| \leq B_{\max}^n, \quad 1 \leq i \leq N_z,$$

and

$$\begin{aligned} \|\underline{b}_+^{i,n+1} - \underline{b}_+^{i+1,n+1}\| &= \|\underline{b}_+^{i,n} - \underline{b}_+^{i+1,n} + \mathbf{D}(\underline{b}_+^{i+1,n} - \underline{b}_+^{i+2,n} - \underline{b}_+^{i,n} + \underline{b}_+^{i+1,n})\| \\ &= dB_{\max}^n, \quad 0 \leq i \leq N_z - 2, \end{aligned}$$

$$\begin{aligned} \|\underline{b}_-^{i,n+1} - \underline{b}_-^{i-1,n+1}\| &\leq \|\underline{b}_-^{i,n} - \underline{b}_-^{i-1,n} + \mathbf{D}(\underline{b}_-^{i-1,n} - \underline{b}_-^{i-2,n} - \underline{b}_-^{i,n} + \underline{b}_-^{i-1,n})\| \\ &\leq dB_{\max}^n, \quad 2 \leq i \leq N_z. \end{aligned}$$

That is to say that the Lax–Wendroff scheme on the interior of the slab makes B_{\max}^n and dB_{\max}^n only to decrease. It is not difficult to see that the Dirichlet condition do the same

$$\begin{aligned}\|\underline{b}_+^{N_z,n+1}\| &= \|\underline{b}_-^{N_z,n+1}\| \leq B_{\max}^n, \\ \|\underline{b}_+^{N_z-1,n+1} - \underline{b}_+^{N_z,n+1}\| &= \|\underline{b}_+^{N_z-1,n} + \mathbf{D}(\underline{b}_+^{N_z,n} - \underline{b}_+^{N_z-1,n}) - \underline{b}_+^{N_z,n} + \mathbf{D}(\underline{b}_-^{N_z-1,n} - \underline{b}_-^{N_z,n})\|, \\ &\leq dB_{\max}^n.\end{aligned}$$

Thus we can conclude

$$\begin{aligned}B_{\max}^{n+1} &\leq \max(B_{\max}^n, \|\underline{b}_-^{0,n+1}\|), \\ dB_{\max}^{n+1} &\leq \max(dB_{\max}^n, \|\underline{b}_-^{1,n+1} - \underline{b}_-^{0,n+1}\|).\end{aligned}$$

Moreover, for $n \leq N_z$, the following estimation holds:

$$\|\underline{b}_+^{0,n}\| \leq B_{\max}^0, \quad (55)$$

$$\|\underline{b}_+^{1,n} - \underline{b}_+^{0,n}\| \leq dB_{\max}^0. \quad (56)$$

The reason is that if one develops $\underline{b}_+^{0,n}$, one has

$$\begin{aligned}\underline{b}_+^{0,n} &= (1 - \mathbf{D})\underline{b}_+^{0,n-1} + \mathbf{D}\underline{b}_+^{1,n-1}, \\ &= (1 - \mathbf{D})^2\underline{b}_+^{0,n-2} + 2(1 - \mathbf{D})\mathbf{D}\underline{b}_+^{1,n-2} + \mathbf{D}^2\underline{b}_+^{2,n-2}, \\ &= \dots \\ &= \sum_{i=0}^j \binom{j}{i} \mathbf{D}^i (1 - \mathbf{D})^{j-i} \underline{b}_+^{i,n-j}, \quad \text{for } 0 \leq i \leq \min(n, N_z),\end{aligned}$$

where the $\binom{j}{i}$ are the combinations of i elements among j . Since $\sum_{i=0}^j \binom{j}{i} \mathbf{D}^i (1 - \mathbf{D})^{j-i} = (1 - \mathbf{D} + \mathbf{D})^j = 1$ the result hold. The same computation can be made with $\|\underline{b}_+^{1,n} - \underline{b}_+^{0,n}\|$.

6.3. Stability of the schemes

We first deal with the u_3 component of the displacement, i.e. the component submitted to the unilateral contact condition. We consider $B_{3\max}^n$ and $dB_{3\max}^n$:

$$B_{3\max}^n = \max_{0 \leq i \leq N_z} (|b_{3+}^{i,n}|, |b_{3-}^{i,n}|), \quad (57)$$

$$dB_{3\max}^n = \max_{0 \leq i \leq N_z-1} (|b_{3+}^{i+1,n} - b_{3+}^{i,n}|, |b_{3-}^{i+1,n} - b_{3-}^{i,n}|), \quad (58)$$

We just saw that

$$\begin{aligned}B_{3\max}^{n+1} &\leq \max(B_{3\max}^n, |b_{3-}^{0,n+1}|), \\ dB_{3\max}^{n+1} &\leq \max(dB_{3\max}^n, |b_{3-}^{1,n+1} - b_{3-}^{0,n+1}|).\end{aligned}$$

Proposition 4. *If the CFL condition $d_2 \leq 1$ holds, schemes (39)–(46), (47), (48) has the following properties:*

$$B_{3\max}^n \leq B_{3\max}^0,$$

$$|u_3^{0,n}| \leq |u_3^{0,0}| + TB_{3\max}^0.$$

Proof. We have

$$b_{3-}^{0,n+1} = \begin{cases} -b_{3+}^{0,n+1} & \text{if } u_{3,t}^{0,n+1} = 0, \\ b_{3+}^{0,n+1} & \text{if } u_{3,t}^{0,n+1} = b_{3+}^{0,n+1}, \end{cases}$$

Thus $|b_{3-}^{0,n+1}| = |b_{3+}^{0,n+1}|$ and

$$B_{3 \max}^{n+1} \leq B_{3 \max}^n \leq \cdots \leq B_{3 \max}^0.$$

For $u_3^{0,n}$ now we have

$$|u_3^{0,n+1} - u_3^{0,n}| \leq \Delta t |b_{3+}^{0,n+1}|,$$

from the scheme, and thus

$$|u_3^{0,n}| \leq TB_{3 \max}^0. \quad \square$$

Proposition 5. *If exists $d_{\min} > 0$ such that $d_{\min} \leq d_2 \leq 1$ then schemes (39)–(46), (49), (50) have the following properties: Exist C_1, C_2, C_3 and C_4 four positive constants which only depend on initial data such that*

$$B_{3 \max}^n \leq C_1,$$

$$|u_3^{0,n}| \leq C_2,$$

$$dB_{3 \max}^n \leq C_3 dB_{3 \max}^0 + \Delta t C_4.$$

Proof. We saw that

$$|b_{3+}^{0,n}| \leq B_{3 \max}^0, \quad |b_{3+}^{1,n} - b_{3+}^{0,n}| \leq dB_{3 \max}^0, \quad (59)$$

and it follows from scheme (49), (50) that

$$\begin{aligned} |u_3^{0,n+1}| &\leq |\Delta t b_{3+}^{0,n+1} + u_3^{0,n}|, \\ &\leq |u_3^{0,n}| + \Delta t B_{3 \max}^0, \end{aligned}$$

thus, since we assume $n < N_z$:

$$\begin{aligned} |u_3^{0,n}| &\leq |u_3^{0,0}| + N_z \Delta t B_{3 \max}^0, \\ &\leq |u_3^{0,0}| + \frac{H}{c_2} B_{3 \max}^0. \end{aligned}$$

Moreover, we can also deduce from schemes (49), (50) that

$$\begin{aligned} |u_3^{0,n+1} - u_3^{0,n}| &\leq \Delta t \left(|b_{3+}^{0,n+1}| + \frac{|u_3^{0,n+1}|}{\tilde{\eta}} \right), \\ &\leq \Delta t B_{3 \max}^0 (1 + 2N_z \Delta t + |u_3^{0,0}|). \end{aligned}$$

and because

$$b_{3-}^{0,n+1} = b_{3+}^{0,n+1} + 2c_2 u_{3,3}^{0,n+1},$$

and

$$u_{3,3}^{0,n+1} = \begin{cases} 0 & \text{if } u_3^{0,n+1} \geq 0, \\ \frac{u_3^{0,n+1}}{c_2 \tilde{\eta}} & \text{if } u_3^{0,n+1} < 0, \end{cases}$$

we have

$$\begin{aligned} |b_{3-}^{0,n+1}| &\leq |b_{3+}^{0,n+1}| + \frac{2|u_3^{0,n+1}|}{\tilde{\eta}}, \\ &\leq B_{3 \max}^0 \left(1 + 2 \frac{H}{c_2} \right) + 2|u_3^{0,0}|. \end{aligned}$$

That is to say

$$B_{3\max}^n \leq B_{3\max}^0 \left(1 + 2 \frac{H}{c_2} \right) + 2|u_3^{0,0}|.$$

Now if we compute

$$\begin{aligned} |b_{3-}^{1,n+1} - b_{3-}^{0,n+1}| &= |b_{3-}^{1,n} + D(b_{3-}^{0,n} - b_{3-}^{1,n}) - b_{3-}^{0,n+1}|, \\ &\leq (1 - d_2)|b_{3-}^{1,n} - b_{3-}^{0,n}| + 2|c_2 u_{3,3}^{0,n} - c_2 u_{3,3}^{0,n+1}| + |b_{3+}^{0,n+1} - b_{3+}^{0,n}|, \\ &\leq (1 - d_2)dB_{3\max}^n + d_2 dB_{3\max}^0 + 2 \frac{\Delta t}{\tilde{\eta}} B_{3\max}^0 (1 + N_z \Delta t + |u_3^{0,0}|), \end{aligned}$$

it comes

$$dB_{3\max}^{n+1} \leq (1 - d_2)dB_{3\max}^n + d_2 dB_{3\max}^0 + 2 \frac{\Delta t}{\tilde{\eta}} B_{3\max}^0 (1 + 2N_z \Delta t + |u_3^{0,0}|),$$

and by induction

$$\begin{aligned} dB_{3\max}^n &\leq (1 - d_2)^n dB_{3\max}^0 + \left(d_2 dB_{3\max}^0 + 2 \frac{\Delta t}{\tilde{\eta}} B_{3\max}^0 (1 + N_z \Delta t) \right) \sum_{i=0}^{n-1} (1 - d_2)^i, \\ &\leq (1 - d_2)^n dB_{3\max}^0 + \frac{1}{d_2} \left(d_2 dB_{3\max}^0 + 2 \frac{\Delta t}{\tilde{\eta}} B_{3\max}^0 (1 + N_z \Delta t + |u_3^{0,0}|) \right), \\ &\leq (1 + (1 - d_2)^n) dB_{3\max}^0 + \Delta t \frac{2}{\tilde{\eta} d_2} B_{3\max}^0 \left(1 + \frac{H}{c_2} \right), \\ &\leq 2dB_{3\max}^0 + \Delta t \frac{2}{\tilde{\eta} d_{\min}} B_{3\max}^0 \left(1 + \frac{H}{c_2} + |u_3^{0,0}| \right). \quad \square \end{aligned}$$

Now, for the complete scheme, including the friction condition, we consider $B_{T\max}^n$ and $dB_{T\max}^n$:

$$B_{T\max}^n = \max_{0 \leq i \leq N_z} (\|b_{T+}^{i,n}\|, \|b_{T-}^{i,n}\|), \quad (60)$$

$$dB_{T\max}^n = \max_{0 \leq i \leq N_z - 1} (\|b_{T+}^{i+1,n} - b_{T+}^{i,n}\|, \|b_{T-}^{i+1,n} - b_{T-}^{i,n}\|), \quad (61)$$

where $b_{T+}^{i,n}$ and $b_{T-}^{i,n}$ are the horizontal components of $b_T^{i,n}$ and $b_T^{i,n}$.

Proposition 6. *If exists $d_{\min} > 0$ such that $d_{\min} \leq d_2 \leq 1$ then schemes (39)–(46), (51), (52) combined with either schemes (47), (48) or (49), (50) has the following properties: there exist four positive constant C_1, C_2, C_3 and C_4 which only depend on the initial condition such that*

$$\begin{aligned} B_{T\max}^n &\leq C_1, \\ \|u_3^{0,n}\| &\leq C_2, \\ dB_{T\max}^n &\leq C_3 dB_{T\max}^0 + \Delta t C_4. \end{aligned}$$

Proof. From (52), it comes

$$\underline{u}_{T,t}^{0,n+1} \in \underline{u}_{T,t}^{0,n} + \frac{\Delta t}{\varepsilon} \left(\frac{G}{c_1} (b_{T+}^{0,n+1} - \underline{u}_{T,t}^{0,n+1}) - S^{n+1} \mu (\|u_{T,t}^{0,n+1} - \underline{V}_e(t^{n+1})\|) \mathbf{Dir}(u_{T,t}^{0,n+1} - \underline{V}_e(t^{n+1})) \right),$$

i.e.

$$\underline{u}_{T,t}^{0,n+1} \in \underline{u}_{T,t}^{0,n} + \Delta t F(t^{n+1}, \underline{u}_{T,t}^{0,n+1}), \quad (62)$$

with a multi-valued map F which has a bounded growth

$$\|F(t^n, x)\| \leq c(t^n)(1 + \|x\|),$$

with $c(t^n)$ defined by

$$c(t^n) = \frac{1}{\varepsilon} \max \left(\frac{G}{c_1}, 1 \right) (\|b_{T+}^{0,n}\| + \|S^n\| \|\mu\|_\infty). \quad (63)$$

From the previous propositions, $\|S^n\|$ is bounded independently of n . Moreover, since we assume $n < N_z$, $\|\underline{b}_{T+}^{0,n}\| \leq B_{T \max}^0$ the quantity $c(t^n)$ is bounded independently of n and Δt for $t < \frac{H}{c_2} d_{\min}$. Setting

$$c_{\max} = \max_{0 \leq i < 2N_z} c(t^i),$$

and from (62) we deduce that

$$\|\underline{u}_{T,t}^{0,n+1}\| \leq \|\underline{u}_{T,t}^{0,n}\| + \Delta t c_{\max} (1 + \|\underline{u}_{T,t}^{0,n+1}\|).$$

We assume $\Delta t < \frac{1}{2c_{\max}}$, thus

$$\begin{aligned} \|\underline{u}_{T,t}^{0,n+1}\| &\leq \frac{1}{1 - \Delta t c_{\max}} (\|\underline{u}_{T,t}^{0,n}\| + c_{\max} \Delta t) \\ &\leq (1 + 2c_{\max} \Delta t) (\|\underline{u}_{T,t}^{0,n}\| + c_{\max} \Delta t). \end{aligned}$$

And by induction

$$\|\underline{u}_{T,t}^{0,n}\| \leq (1 + 2c_{\max} \Delta t)^n \|\underline{u}_{T,t}^0\| + \Delta t c_{\max} \sum_{i=1}^n (1 + 2c_{\max} \Delta t)^i.$$

Since $(1 + 2c_{\max} \frac{T}{N})^N \leq e^{2Tc_{\max}}$:

$$\|\underline{u}_{T,t}^{0,n}\| \leq e^{4\frac{H}{c_2} d_{\min} c_{\max}} \left(\|\underline{u}_{T,t}^0\| + 2c_{\max} \frac{H}{c_2} d_{\min} \right).$$

Setting

$$C_1^0 = e^{4\frac{H}{c_2} d_{\min} c_{\max}} \left(\|\underline{u}_{T,t}^0\| + 2 \frac{H}{c_2} c_{\max} d_{\min} \right),$$

we deduce immediately

$$\begin{aligned} \|\underline{b}_{T-}^{0,n+1}\| &\leq \|\underline{b}_{T+}^{0,n+1}\| + 2\|\underline{u}_{T,t}^{0,n}\|, \\ &\leq B_{T \max}^0 + 2C_1^0. \end{aligned}$$

Thus

$$\begin{aligned} \|\underline{u}_{T,t}^{0,n+1} - \underline{u}_{T,t}^{0,n}\| &\leq \Delta t \|F(t^n + 1, \underline{u}_{T,t}^{0,n+1})\|, \\ &\leq \Delta t c_{\max} (1 + \|\underline{u}_{T,t}^{0,n+1}\|), \\ &\leq \Delta t c_{\max} (1 + C_1^0), \end{aligned}$$

and

$$\begin{aligned} \|\underline{b}_{T-}^{1,n+1} - \underline{b}_{T-}^{0,n+1}\| &\leq \|\underline{b}_{T-}^{1,n} + d_1 (\underline{b}_{T-}^{0,n} - \underline{b}_{T-}^{1,n}) - \underline{b}_{T-}^{0,n+1}\|, \\ &\leq (1 - d_1) \|\underline{b}_{T-}^{0,n} - \underline{b}_{T-}^{1,n}\| + 2\|\underline{u}_{T,t}^{0,n+1} - \underline{u}_{T,t}^{0,n}\| + \|\underline{b}_{T+}^{0,n+1} - \underline{b}_{T+}^{0,n}\|, \\ &\leq (1 - d_1) dB_{T \max}^n + d_1 dB_{T \max}^0 + 2\Delta t c_{\max} (1 + C_1^0). \end{aligned}$$

And now just as for Proposition 5, we conclude by induction that

$$dB_{T_{\max}}^n \leq 2dB_{T_{\max}}^0 + 2 \frac{\Delta t}{d_1} c_{\max} (1 + C_1^0). \quad \square$$

6.4. Convergence

To simplify the proof, we will assume that the ratio $\frac{\Delta t}{\Delta z}$ is constant, i.e. N_z is proportional to N_t :

$$N_z = \frac{1}{\eta} N_t,$$

where η is independent of N_t and satisfies the CFL condition

$$\eta \geq \frac{c_2 T}{H},$$

Let us denote by $\underline{u}_{,t}^{i,n,N_t}, \underline{u}_{,3}^{i,n,N_t}$ the discrete solution given by one of the schemes and $\underline{u}_{,t}^{N_t}(t, x_3), \underline{u}_{,3}^{N_t}(t, x_3)$ the linear interpolations of the discrete solutions

$$\underline{u}_{,t}^{N_t}(n\Delta t + \delta t, i\Delta z + \delta z) = (1 - \delta t)((1 - \delta z)\underline{u}_{,t}^{i,n,N_t} + \delta z \underline{u}_{,t}^{i+1,n,N_t}) + \delta t((1 - \delta z)\underline{u}_{,t}^{i,n+1,N_t} + \delta z \underline{u}_{,t}^{i+1,n+1,N_t}),$$

with $0 \leq \delta t \leq \Delta t$ and $0 \leq \delta z \leq \Delta z$. An analogous definition is taken for $\underline{u}_{,3}^{N_t}(t, x_3)$.

Lemma 1. *If the sequence $(\underline{u}_{,t}^{N_t}, \underline{u}_{,3}^{N_t})_{N_t}$ uniformly converges toward a pair of functions $(\underline{u}_{,t}, \underline{u}_{,3})$ and if the functions $\underline{u}_{,t}$ and $\underline{u}_{,3}$ are continuous on $[0, T] \times [0, H]$ then these functions satisfy the relations (18)–(20) on the characteristic lines.*

Proof. We establish the result on only one relation and one component, the vertical component, because the proof is identical for the other cases. As before, we define

$$\begin{aligned} b_{3+}(t, x_3) &= u_{3,t}(t, x_3) + c_2 u_{3,3}(t, x_3), \\ b_{3+}^{N_t}(t, x_3) &= u_{3,t}^{N_t}(t, x_3) + c_2 u_{3,3}^{N_t}(t, x_3), \\ b_{3+}^{i,n,N_t} &= u_{3,t}^{i,n,N_t} + c_2 u_{3,3}^{i,n,N_t}. \end{aligned}$$

For any z_0, t_0, δ_0 such that

$$\delta_0 \in \left[0, \max \left(T, \frac{H}{c_2}\right)\right], \quad z_0 \in \left[0, H - \eta \delta_0 \frac{H}{T}\right], \quad t_0 \in [0, T - \delta_0].$$

Let us define

$$n_1^{N_t} = \left\lfloor \frac{N_t t_0}{T} \right\rfloor, \quad i_2^{N_t} = \left\lfloor \frac{N_t z_0}{\eta H} \right\rfloor, \quad n_2^{N_t} = \left\lfloor \frac{N_t(t_0 + \delta_0)}{T} \right\rfloor,$$

$$\Delta_n^{N_t} = n_2^{N_t} - n_1^{N_t},$$

where $\lfloor x \rfloor$ is the integer part of x . Of course, $n_1^{N_t} \Delta t$ converges toward t_0 when N_t goes to infinity; and $(n_2^{N_t} \Delta t, i_2^{N_t} \Delta z)$ converges toward $(t_0 + \delta_0, z_0)$. Thus

$$b_{3+}^{i_2^{N_t}, n_2^{N_t}, N_t} \xrightarrow{N_t \rightarrow +\infty} b_{3+}(t_0 + \delta_0, z_0),$$

From the Lax–Wendroff scheme, it comes

$$\begin{aligned}
 b_{3+}^{i_2^{N_t}, n_2^{N_t}, N_t} &= (1 - d_2) b_{3+}^{i_2^{N_t}, n_2^{N_t} - 1, N_t} + d_2 b_{3+}^{i_2^{N_t} + 1, n_2^{N_t} - 1, N_t}, \\
 &= (1 - d_2)^2 b_{3+}^{i_2^{N_t}, n_2^{N_t} - 2, N_t} + 2(1 - d_2)d_2 b_{3+}^{i_2^{N_t} + 1, n_2^{N_t} - 2, N_t} + d_2^2 b_{3+}^{i_2^{N_t} + 2, n_2^{N_t} - 2, N_t}, \\
 &= \dots \\
 &= \sum_{i=0}^{\Delta_n^{N_t}} \binom{i}{\Delta_n^{N_t}} d^i (1 - d)^{\Delta_n^{N_t} - i} b_{3+}^{i_2^{N_t} + i, n_1^{N_t}, N_t}, \\
 &= \sum_{i=0}^{\Delta_n^{N_t}} B_i^{\Delta_n^{N_t}} (d_2) b_{3+}^{i_2^{N_t} + i, n_1^{N_t}, N_t},
 \end{aligned}$$

where $B_i^{\Delta_n^{N_t}}$ is a Bernstein polynomial. With the decomposition

$$b_{3+}^{i_2^{N_t} + i, n_1^{N_t}, N_t} = b_{3+} \left(t_0, z_0 + \frac{i}{\Delta_n^{N_t}} \left(\eta \delta_0 \frac{H}{T} \right) \right) + \left(b_{3+}^{i_2^{N_t} + i, n_1^{N_t}, N_t} - b_{3+} \left(t_0, z_0 + \frac{i}{\Delta_n^{N_t}} \left(\eta \delta_0 \frac{H}{T} \right) \right) \right),$$

it comes

$$\begin{aligned}
 b_{3+}^{i_2^{N_t}, n_2^{N_t}, N_t} &= \sum_{i=0}^{\Delta_n^{N_t}} B_i^{\Delta_n^{N_t}} (d_2) b_{3+} \left(t_0, z_0 + \frac{i}{\Delta_n^{N_t}} \left(\eta \delta_0 \frac{H}{T} \right) \right) \\
 &\quad + \sum_{i=0}^{\Delta_n^{N_t}} B_i^{\Delta_n^{N_t}} (d_2) \left(b_{3+}^{i_2^{N_t}, n_1^{N_t}, N_t} - b_{3+} \left(t_0, z_0 + \frac{i}{\Delta_n^{N_t}} \left(\eta \delta_0 \frac{H}{T} \right) \right) \right).
 \end{aligned}$$

From the approximation with Bernstein polynomials, we know that for each continuous function $f : [0, 1] \rightarrow \mathbb{R}$, the function $f^N = \sum_{i=0}^N B_i^N f(\frac{i}{N})$ uniformly converges toward f when N goes to $+\infty$. Since the function $b_{3+}^{N_t}$ is continuous, the expression

$$\sum_{i=0}^{\Delta_n^{N_t}} B_i^{\Delta_n^{N_t}} (d_2) b_{3+} \left(t_0, z_0 + \frac{i}{\Delta_n^{N_t}} \left(\eta \delta_0 \frac{H}{T} \right) \right),$$

converges towards the value

$$b_{3+} \left(t_0, z_0 + d_2 \left(\eta \delta_0 \frac{H}{T} \right) \right) = b_{3+} (t_0, z_0 + \delta_0 c_2),$$

when $\Delta_n^{N_t}$ goes to $+\infty$.

Now, due to uniform convergence, it comes

$$\forall \varepsilon, \exists N_\varepsilon > 0; \quad \forall N_t > N_\varepsilon, \quad \forall (t, z) \in [0, T] \times [0, H], |b_{3+}^{N_t}(t, z) - b_{3+}(t, z)| < \varepsilon.$$

Considering the uniform continuity of b_{3+} on $[0, T] \times [0, H]$ we get

$$\forall \varepsilon, \exists \delta_\varepsilon; \quad \forall (t_1, z_1), \forall (t_2, z_2), \quad d((t_1, z_1), (t_2, z_2)) < \delta_\varepsilon \Rightarrow |b_{3+}(t_1, z_1) - b_{3+}(t_2, z_2)| < \varepsilon.$$

And because

$$|t_0 - n_1^{N_t} \Delta t| \leq \Delta t,$$

and

$$\begin{aligned}
 \left| (i_2^{N_t} + i) \Delta z - z_0 + \frac{i}{\Delta_n^{N_t}} \left(\eta \delta_0 \frac{H}{T} \right) \right| &\leq \Delta z + \Delta z \frac{i}{\Delta_n^{N_t}} \left(\left\lfloor \frac{t_0}{\Delta t} \right\rfloor - \left\lfloor \frac{t_0}{\Delta t} \right\rfloor + \frac{\delta_0}{\Delta t} \right) - \frac{\delta_0}{\Delta t}, \\
 &\leq 3 \Delta z,
 \end{aligned}$$

then

$$\forall \varepsilon, \exists N_\varepsilon > 0; \quad \forall N_t > N_\varepsilon, \quad |b_{3+}^{N_t}(n_1^{N_t} \Delta t, (i_2^{N_t} + i) \Delta z) - b_{3+}(n_1^{N_t} \Delta t, (i_2^{N_t} + i) \Delta z)| < \varepsilon, \\ |b_{3+}(n_1^{N_t} \Delta t, (i_2^{N_t} + i) \Delta z) - b_{3+}(t_0, z_0 + \delta_0 c_2)| < \varepsilon.$$

and

$$\forall \varepsilon, \exists N_\varepsilon > 0; \quad \forall N_t > N_\varepsilon, \quad |b_{3+}^{N_t}(n_1^{N_t} \Delta t, (i_2^{N_t} + i) \Delta z) - b_{3+}(t_0, z_0 + \delta_0 c_2)| < 2\varepsilon.$$

then the expression

$$\sum_{i=0}^{N_t} B_i^{\Delta_n^{N_t}}(d_2) \left(b_{3+}^{N_t}(n_1^{N_t} \Delta t, (i_2^{N_t} + i) \Delta z) - b_{3+} \left(t_0, z_0 + \frac{i}{\Delta_n^{N_t}} \left(\eta \delta_0 \frac{H}{T} \right) \right) \right),$$

converges to zero when N_t goes to $+\infty$.

We just established that the relation

$$b_{3+}(t_0 + \delta_0, z_0) = b_{3+}(t_0, z_0 + c_2 \delta_0), \quad (64)$$

is satisfied for all $\delta_0 \in [0, \max(T, \frac{H}{c_2})]$, and for all pairs (t_0, z_0) satisfying

$$z_0 \in \left[0, H - \eta \delta_0 \frac{H}{T} \right], \quad t_0 \in [0, T - \delta_0].$$

It is not sufficient to conclude, because we need to establish the result for all pairs (t_0, z_0) satisfying

$$z_0 \in [0, H - c_2 \delta_0], \quad t_0 \in [0, T - \delta_0].$$

But, equality (64) can be split in two parts:

$$b_{3+} \left(t_0 + \frac{\delta_0}{2}, z_0 + c_2 \frac{\delta_0}{2} \right) = b_{3+} \left(t_0, z_0 + c_2 \frac{\delta_0}{2} \right), \\ b_{3+} \left(t_0 + \delta_0, z_0 \right) = b_{3+} \left(t_0 + \frac{\delta_0}{2}, z_0 + c_2 \frac{\delta_0}{2} \right).$$

which are established for

$$z_0 \in \left[0, H - c_2 \frac{\delta_0}{2} - \eta \frac{\delta_0}{2} \frac{H}{T} \right].$$

Using this reasoning, one sees that relation (64) is satisfied for

$$z_0 \in \left[0, H - c_2 \frac{\delta_0}{2} \right].$$

And since the function b_{3+} is continuous, this relation is also satisfied for

$$z_0 \in \left[0, H - c_2 \frac{\delta_0}{2} \right]. \quad \square$$

Proposition 7. *If the following hypotheses are satisfied:*

- the initial data \underline{u}_0 is differentiable and its derivative is Lipschitz continuous,
- the initial data \underline{u}_1 is Lipschitz continuous,
- the ratio $\eta = \frac{N_t}{N_z}$ is constant and satisfies $\eta \geq \frac{c_2 T}{H}$,

then schemes (39)–(46), (49)–(52) is convergent.

Proof. Let us still consider $\underline{u}_t^{N_t}(t, x_3)$, $\underline{u}_3^{N_t}(t, x_3)$ the linear interpolations of the discrete solutions of schemes (39)–(46), (49)–(52). From Propositions 5 and 6 these functions are bounded in $(L^\infty([0, T] \times [0, H]))^3$

independently of N_t , and moreover exist C_3 and C_4 , two positive constants which only depend on the initial data such that

$$\begin{aligned} & \max_{0 \leq i \leq N_t-1} (\|\underline{u}_{,t}^{i+1,n,N_t} - \underline{u}_{,t}^{i,n,N_t}\|, \|\mathbf{C}\underline{u}_{,3}^{i+1,n,N_t} - \mathbf{C}\underline{u}_{,3}^{i,n,N_t}\|) \\ & < C_3 \max_{0 \leq i \leq N_t-1} (\|\underline{u}_1((i+1)\Delta t) - \underline{u}_1(i\Delta t)\|, \|\mathbf{C}\partial_{x_3}\underline{u}_0((i+1)\Delta t) - \mathbf{C}\partial_{x_3}\underline{u}_0(i\Delta t)\|) + C_4\Delta t. \end{aligned}$$

Since we assume \underline{u}_0 is differentiable and is derivative is Lipschitz and \underline{u}_1 is Lipschitz, we conclude that $\underline{u}_{,t}^{N_t}$ and $\underline{u}_{,3}^{N_t}$ are bounded in $(W^{1,\infty}(]0, T[\times]0, H[))^3$ independently of N_t .

From the Rellich–Kondrachov Theorem (see [1] for instance) the space $W^{1,\infty}(]0, T[\times]0, H[)$ is included in the space $\mathcal{C}(]0, T[\times]0, H[)$ of the continuous functions on $[0, T] \times [0, H]$, and the inclusion is compact. Thus, from the sequence $(\underline{u}_{,t}^{N_t}, \underline{u}_{,3}^{N_t})_{N_t}$ it is possible to extract a sub-sequence $(\underline{u}_{,t}^{N'_t}, \underline{u}_{,3}^{N'_t})_{N'_t}$ which uniformly converges toward a pair of continuous functions $(\underline{u}_{,t}, \underline{u}_{,3})$.

From Lemma 1, we know that $\underline{u}_{,t}(t, x_3)$ and $\underline{u}_{,3}(t, x_3)$ satisfy to the relations (18)–(20) on the characteristic lines. In particular, for $t < \frac{H}{c_2}$, we get

$$\begin{aligned} b_{1+}(t, 0) &= u_{11}(c_1 t) + c_1 \partial_{x_3} u_{10}(c_1 t), \\ b_{2+}(t, 0) &= u_{21}(c_1 t) + c_1 \partial_{x_3} u_{20}(c_1 t), \\ b_{3+}(t, 0) &= u_{31}(c_2 t) + c_2 \partial_{x_3} u_{30}(c_2 t). \end{aligned}$$

And, because of the uniform convergence, the functions $b_{1+}^{N'_t}(., 0)$, $b_{2+}^{N'_t}(., 0)$ and $b_{3+}^{N'_t}(., 0)$ uniformly converge toward the respective values $u_{11}(c_1 t) + c_1 \partial_{x_3} u_{10}(c_1 t)$, $u_{21}(c_1 t) + c_1 \partial_{x_3} u_{20}(c_1 t)$ and $u_{31}(c_2 t) + c_2 \partial_{x_3} u_{30}(c_2 t)$.

The expression of the implicit Euler scheme is

$$\begin{aligned} u_3^{0,n+1,N'_t} &= \left(u_3^{0,n,N'_t} + \Delta t b_{3+}^{0,n+1,N'_t} \right)_+ + \frac{1}{1 + \frac{\Delta t}{\tilde{\eta}}} \left(u_3^{0,n,N'_t} + \Delta t b_{3+}^{0,n+1,N'_t} \right)_-, \\ S^{n+1,N'_t} &= -\frac{\lambda + 2G}{c_2} \left(b_{3+}^{0,n+1,N'_t} - u_{3,t}^{0,n+1,N'_t} \right), \\ \underline{u}_{T,t}^{0,n+1,N'_t} &\in \underline{u}_{T,t}^{0,n,N'_t} + \frac{\Delta t}{\varepsilon} \left(\frac{G}{c_1} \left(b_{T+}^{0,n+1,N'_t} - \underline{u}_{T,t}^{0,n+1,N'_t} \right) \right. \\ &\quad \left. - S^{n+1,N'_t} \mu \left(\|\underline{u}_{T,t}^{0,n+1,N'_t} - \underline{V}_e(t^{n+1})\| \right) \text{Dir} \left(\underline{u}_{T,t}^{0,n+1,N'_t} - \underline{V}_e(t^{n+1}) \right) \right), \end{aligned}$$

which can be rewritten as

$$\begin{aligned} u_3^{0,n+1,N'_t} &= \left(u_3^{0,n,N'_t} + \Delta t b_{3+}(t^{n+1}, 0) \right)_+ + \frac{1}{1 + \frac{\Delta t}{\tilde{\eta}}} \left(u_3^{0,n,N'_t} + \Delta t b_{3+}(t^{n+1}, 0) \right)_- + \epsilon_3^{n+1,N'_t}, \\ S^{n+1,N'_t} &= -\frac{\lambda + 2G}{c_2} \left(b_{3+}(t^{n+1}, 0) - u_{3,t}^{0,n+1,N'_t} \right), \\ \underline{u}_{T,t}^{0,n+1,N'_t} &\in \underline{u}_{T,t}^{0,n,N'_t} + \frac{\Delta t}{\varepsilon} \left(\frac{G}{c_1} \left(b_{T+}(t^{n+1}, 0) - \underline{u}_{T,t}^{0,n+1,N'_t} \right) \right. \\ &\quad \left. - S^{n+1,N'_t} \mu \left(\|\underline{u}_{T,t}^{0,n+1,N'_t} - \underline{V}_e(t^{n+1})\| \right) \text{Dir} \left(\underline{u}_{T,t}^{0,n+1,N'_t} - \underline{V}_e(t^{n+1}) \right) \right) + \epsilon_T^{n+1,N'_t}. \end{aligned}$$

This scheme is an implicit Euler scheme for the Cauchy problem corresponding to Eq. (38) and the Cauchy problem (35). Now, under the conditions that the right-hand side of the differential inclusion of (35) is one-sided Lipschitz and satisfies the growth condition (63) and also that the quantities ϵ_3^{n+1,N'_t} , and ϵ_T^{n+1,N'_t} satisfy

$$\lim_{N_t \rightarrow \infty} \left[\max_{1 \leq j \leq N} (\|\epsilon_T^{n,N'_t}\|, |\epsilon_3^{n,N'_t}|) \right] = 0,$$

it is a result of Lempio (see [9,6]) that this scheme is convergent. The quantities ϵ_3^{n+1,N'_t} , and ϵ_T^{n+1,N'_t} satisfy these conditions because the following estimations hold:

$$\begin{aligned}
|\epsilon_3^{n,N'_t}| &\leq \left(\frac{1}{1 + \frac{\Delta t}{\eta}} \right) |b_{3+}(t^{n+1}, 0) - b_{3+}^{0,n+1,N'_t}|, \\
&\leq 2\Delta t |b_{3+}(t^{n+1}, 0) - b_{3+}^{0,n+1,N'_t}|, \\
\|\epsilon_T^{n,N'_t}\| &\leq \frac{\Delta t}{\varepsilon} \left(\frac{G}{c_1} + \frac{\lambda + 2G}{c_2} \|\mu\|_\infty \right) \|b_{T+}(t^{n+1}, 0) - b_{T+}^{0,n+1,N'_t}\|,
\end{aligned}$$

and because the functions $b_{T+}^{N'_t}(., 0)$ and $b_{T+}^{N'_t}(., 0)$ uniformly converge, respectively toward $b_{T+}(., 0)$ and $b_{3+}(., 0)$.

In conclusion, for $t \in [0, \frac{H}{c_2}]$ the function $u(t, 0)_{3,t}$ is the derivative with respect to time of the unique solution of the Cauchy problem associated to the differential equation (38) and the function $u(t, 0)_{T,t}$ is the unique solution of (35). Additionally, From Lemma (1), the functions $\underline{u}_{,t}(t, x_3)$ and $\underline{u}_3(t, x_3)$ satisfy to the relations 18,18,20 on the characteristic lines. From the uniqueness of this solution, we can conclude that all the uniformly convergent sub-sequences of the sequence $(\underline{u}_{,t}^{N'_t}, \underline{u}_3^{N'_t})_{N_t}$ converge toward $(\underline{u}_{,t}, \underline{u}_3)$. This is sufficient to say that the whole sequence $(\underline{u}_{,t}^{N_t}, \underline{u}_3^{N_t})_{N_t}$ uniformly converges toward $(\underline{u}_{,t}, \underline{u}_3)$. \square

7. Numerical experiments

In this section, we present a few numerical experiments made with scheme (39)–(48), 51,52, i.e. with the non-regularized unilateral contact condition. Mechanical parameters for the material corresponds to a little metallic piece ($\lambda = 107 \times 10^9$ Pa, $G = 77 \times 10^9$ Pa, $\rho = 7800$ Kg/m³) of height 5 cm. The coefficient of friction has the shape of Fig. 2, with $\mu_s = 0.8$ and $\mu_d = 0.5$.

7.1. Numerical test of convergence

We here test the stability and the convergence of schemes (39)–(48), 51 on an example. This convergence has not been established theoretically. In order to have a relevant case, we choose an example where the contact pressure varies, and where sometimes the contact does not hold. The perturbation parameter for the friction condition is chosen constant ($\varepsilon = 10.0$). The computations are made with different values of Δt and Δz with a constant ratio $\frac{\Delta t}{\Delta z}$.

The results of computations are presented on Figs. 3 and 4.

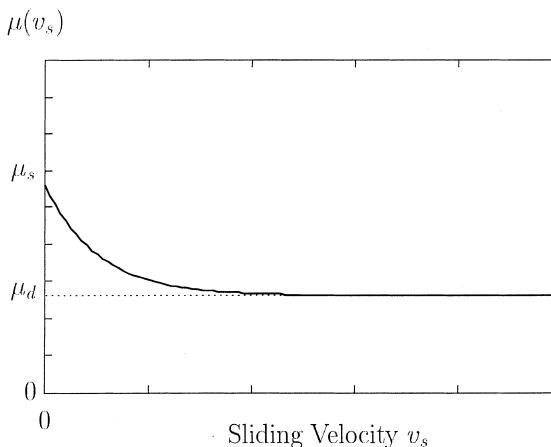


Fig. 2. Example of friction coefficient.

Experiment	Δt	Δz	$\ v_x^n - v_x^{n-1}\ _\infty$	$\ u_z^n - u_z^{n-1}\ _\infty$
1	10^{-7}	10^{-4}		
2	4×10^{-8}	4×10^{-5}	0.3678	2.5342×10^{-6}
3	10^{-8}	10^{-5}	0.1277	9.2298×10^{-7}
4	4×10^{-9}	4×10^{-6}	0.0905	7.1579×10^{-7}
5	10^{-9}	10^{-6}	0.0237	2.8334×10^{-7}

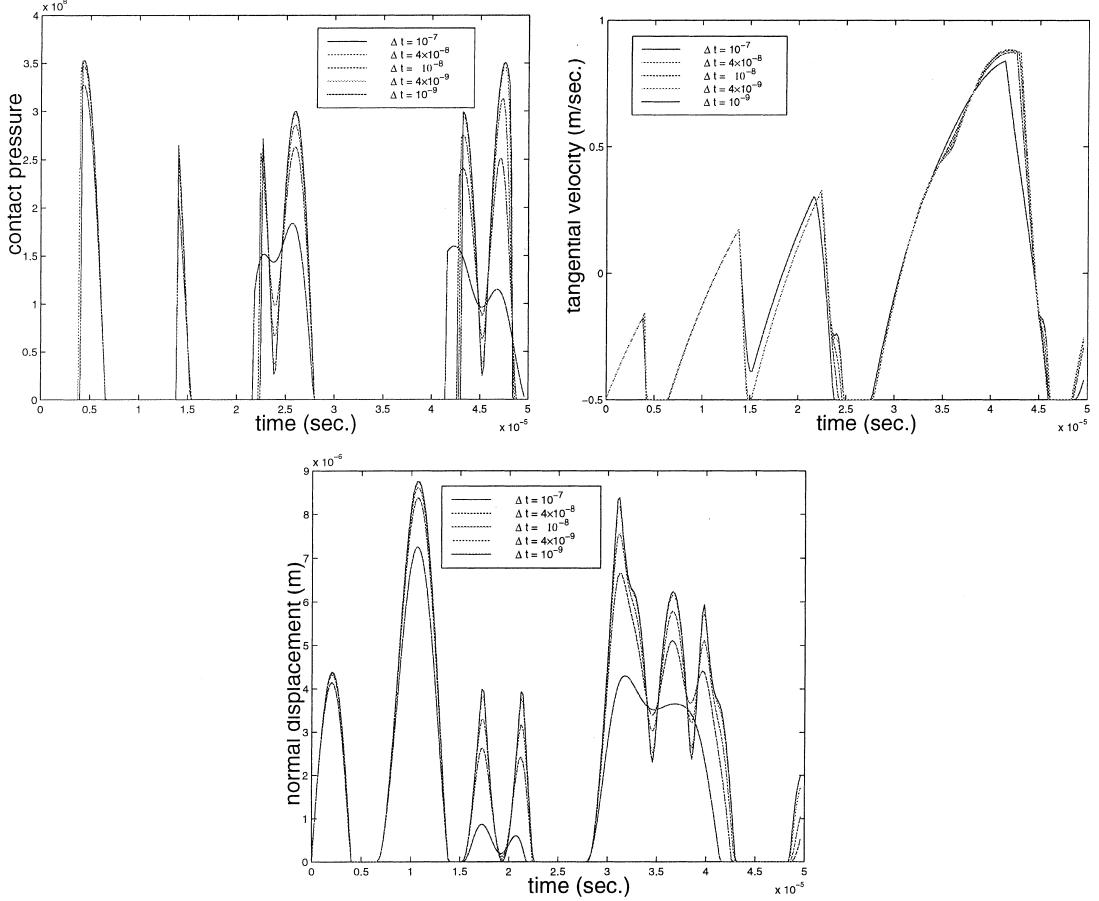
Fig. 3. Different values of Δt , Δz and maximal difference between two successive computations.

Fig. 4. Evolution of contact pressure, tangential velocity and normal displacement.

There is a numerical convergence, but, due to the bad regularity of the continuous problem solution, the rate of convergence seems to be lower than linear on Δt .

7.2. Example of global behavior

Now, we choose a small perturbation parameter ($\varepsilon = 0.1$) and a constant contact pressure.

Fig. 5 presents the evolution of tangential displacement on the contact boundary. One can see that there is a chaotic behavior at the beginning of the simulation, but, after a certain time of simulation the motion becomes more regular. Fig. 6 presents a detail of the simulation where the motion is regular. This dynamical system seems to be attracted by a solution which is nearly periodical. It is a classical result that a one degree of freedom elastic system with such a law of friction often has a periodic attractor. Here, for the infinite slab, the difference is the infinite number of degree of freedom and the presence of jump in velocity.

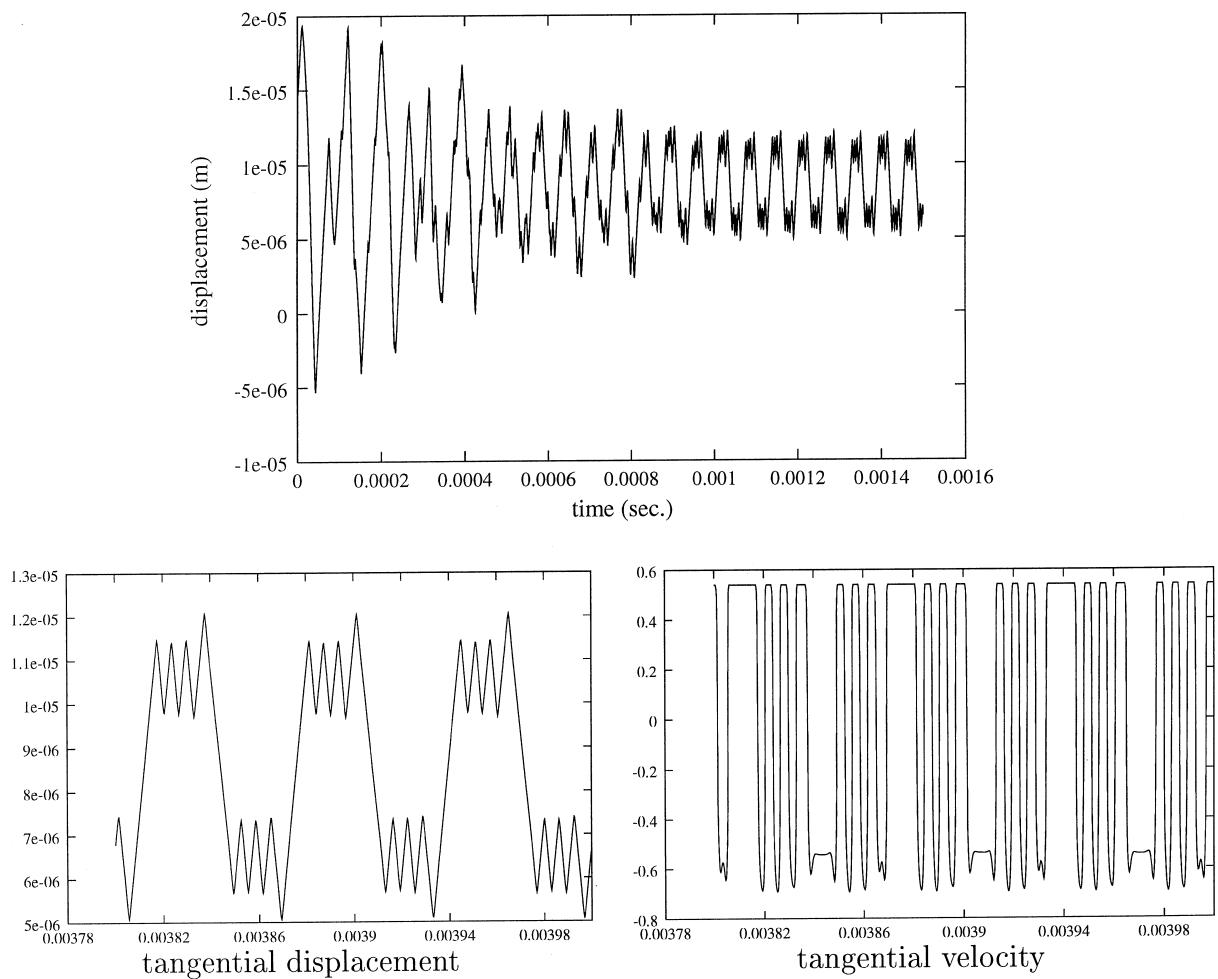


Fig. 5. Evolution of the tangential displacement on the contact boundary.

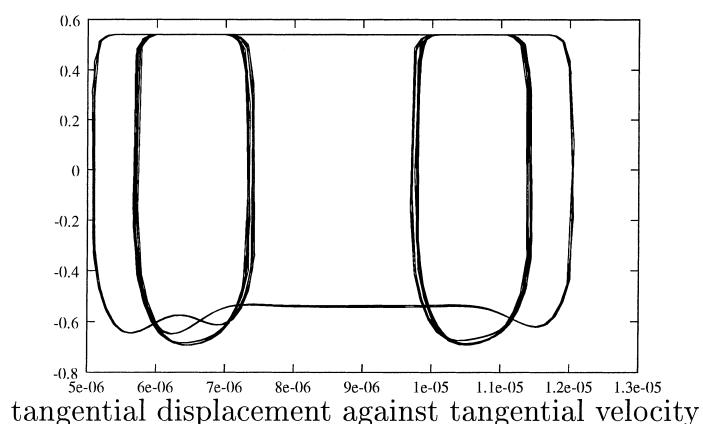


Fig. 6. Detail of the evolution of tangential displacement and velocity.

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