

# Melonic large $N$ limit of 5-index irreducible random tensors

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*Tensor Journal Club*  
February 16, 2022

Results based on:

- ▶ [arXiv:2104.03665](#) – Commun. Math. Phys. (2022) with S. Harribey.



but also

- ▶ [arXiv:1712.00249](#) – Commun. Math. Phys. (2019) with D. Benedetti, R. Gurau and M. Kolanowski.
- ▶ [arXiv:1803.02496](#) – JHEP (2018)

# OUTLINE

Introduction

$O(N)$  irreducible random tensors with complete graph interaction

Existence of the large  $N$  limit

Melonic dominance at leading order

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## RANDOM TENSORS

Space of tensors  $T = T_{a_1 \dots a_p}$ ,  $a_i \in \{1, \dots, N\}$ , equipped with measure of the form:

$$d\nu(T) = d\mu_{\mathbf{P}}(T)e^{-S_N(T)}$$

- ▶  $d\mu_{\mathbf{P}}$  is Gaussian with covariance  $\mathbf{P}$ :

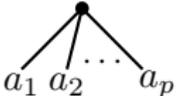
$$\int d\mu_{\mathbf{P}}(T) T_{a_1 \dots a_p} T_{b_1 \dots b_p} = \mathbf{P}_{a_1 \dots a_p, b_1 \dots b_p}$$

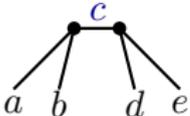
- ▶ both  $\mathbf{P}$  and  $S_N$  are invariant under the action of some unitary group:  $O(N)$ ,  $U(N)$  or  $Sp(N)$ .

What type of universal behaviour can we obtain in the asymptotic limit  
 $N \rightarrow \infty$  ?

## TENSORS AND INVARIANTS

Symmetric tensor:

$$T_{a_1 a_2 \dots a_p} =$$


$$\sum_{c=1}^N T_{abc} T_{cde} =$$


Connected invariants:

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$$T_{a_1 a_2 \dots a_p} = \begin{array}{c} \bullet \\ \diagdown \quad \diagup \\ a_1 \quad a_2 \quad \dots \quad a_p \end{array}$$

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Connected invariants:

$$p = 1$$



$$(\phi_a \phi^a)$$

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Connected invariants:

$$p = 2 \quad \begin{array}{c} \circ \\ \bullet \end{array} \quad \begin{array}{c} \circ \\ \bullet \quad \bullet \end{array} \quad \begin{array}{c} \circ \\ \bullet \quad \bullet \quad \bullet \end{array} \quad \begin{array}{c} \circ \\ \bullet \quad \bullet \quad \bullet \quad \bullet \end{array} \quad \dots \quad (\text{tr}(M^n))$$

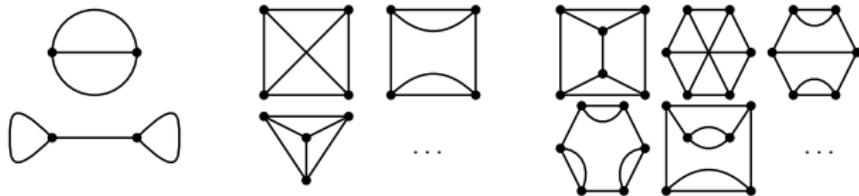
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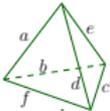
 $p = 3$ 

$$\#\{\text{invariants of order } 2n\} \sim \left(\frac{3}{2}\right)^n n!$$

⇒ Rapid growth of theory space for  $p \geq 3$ . Universal features at large  $N$ ?

QG IN  $D \geq 3$  AS A TENSOR INTEGRAL?

$$\mathcal{F}(\lambda) = \ln \int dT \exp \left( -T_{abc}T_{abc} + \frac{\lambda}{N^\alpha} T_{aeb}T_{bfc}T_{ced}T_{dfa} \right)$$



[Ambjørn, Durhuss, Jónsson '91; Gross '91; Sasakura '91;...]

► Challenges:

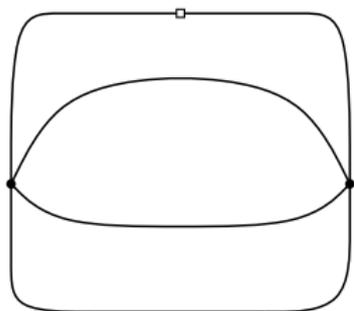
- matrix techniques not available (spectral representation)
- interplay between combinatorics and topology: nice global properties from local Feynman rules?
- large- $N$  expansion ?

► Path to progress: [Gurau '09; Gurau, Rivasseau, Bonzom,... '10s]

- more symmetry:  $U(N)^D \rightarrow$  colored tensor models
- tractable combinatorics, mapping to sufficiently regular topological spaces.

⇒ universal large- $N$  expansion, in any  $D \geq 3$

## MELONS



[BONZOM, GURAU, RIELLO, RIVASSEAU '11...]

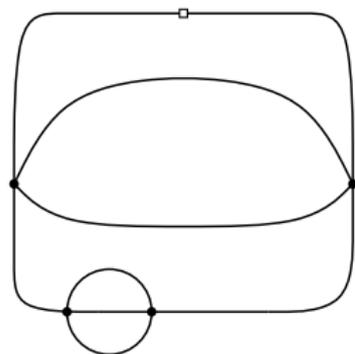
$$\omega(\Delta) = 0 \quad \Leftrightarrow \quad \Delta \text{ is melonic}$$

→ special triangulations of the  $D$ -sphere, with a tree-like combinatorial structure.

Closed equation for their generating function:

$$G(\lambda) = 1 + \lambda G(\lambda)^{D+1} \quad (\text{Fuss-Catalan})$$

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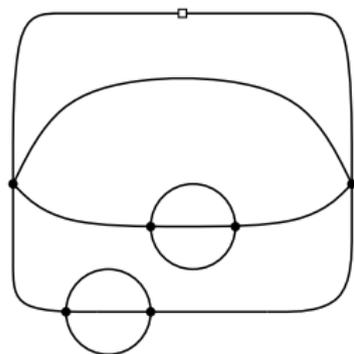
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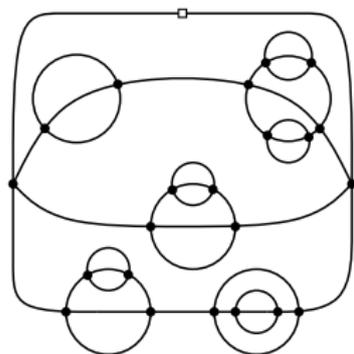
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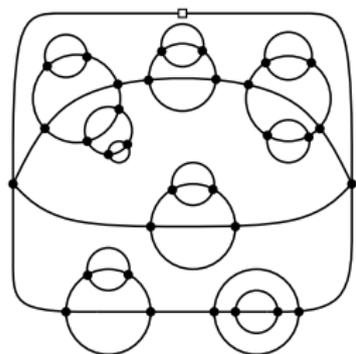
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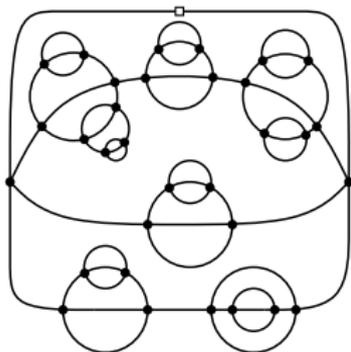
Critical regime / continuum limit:

Melons are **branched polymers** [Gurau, Ryan '13]

i.e. they converge to the **continuous random tree** [Aldous '91].

$$d_{\text{spectral}} = 4/3$$

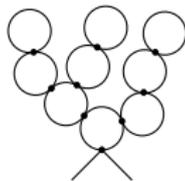
⇒ **strong universality**: limit independent of  $D$ !



# THE MELONIC LIMIT IN LARGE- $N$ QFT

Vector field  $\phi_a(x)$

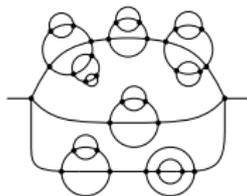
Bubble diagrams



Easy

Tensor field  $T_{abc}(x)$

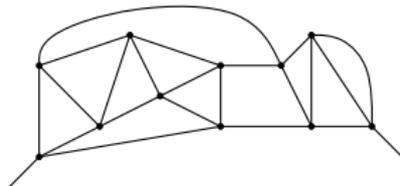
Melon diagrams



Tractable

Matrix field  $M_{ab}(x)$

Planar diagrams

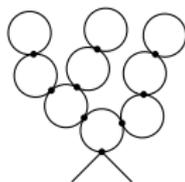


Hard

## THE MELONIC LIMIT IN LARGE-N QFT

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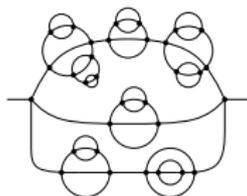
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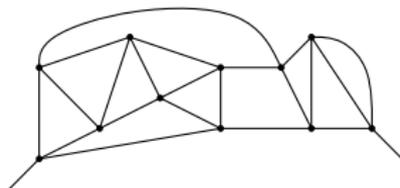
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Tractable

Matrix field  $M_{ab}(x)$ 

Planar diagrams



Hard

- ▶ Strong-coupling regime of the **SYK model**: disordered model of Majorana fermions [Kitaev; Maldacena, Stanford; Gross, Rosenhaus,...]
- ▶ **Tensor model** realization: fermionic tensor field  $\Psi_{abc}(t)$ 
  - ▶ no disorder [Witten '16; Klebanov, Tarnopolsky '16]
  - ▶ natural QFT generalizations [Klebanov et al., Gurau, Benedetti, Haribey, Suzuki, Lettera,...]

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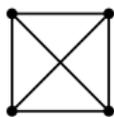
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## MODELS

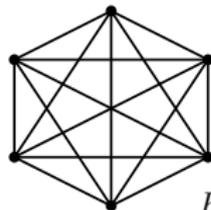
$p$ -index tensor  $T_{a_1 \dots a_p}$ , with  $p$  odd and measure of the form:

$$d\nu(T) = d\mu_P(T) e^{-S_N(T)}$$

- ▶  $P$  = orthogonal projector on an irreducible representation of  $O(N)$ ;
- ▶  $S_N = -\frac{\lambda}{N^\alpha} \text{Inv}(T)$ , where  $\text{Inv}(T)$  is a complete-graph invariant (graph  $K_{p+1}$ ).



$K_4$  ( $p = 3$ )



$K_6$  ( $p = 5$ )

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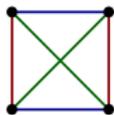
# PRECURSORS: COLORED $O(N)$ MODELS

$T_{a_1 a_2 \dots a_p}$ , in fundamental representation of  $O(N) \times O(N) \times \dots \times O(N)$ :

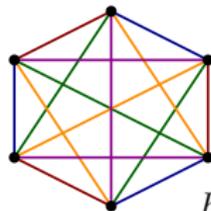
▶  $P_{a_1 a_2 \dots a_p, b_1 b_2 \dots b_p} = \delta_{a_1 b_1} \delta_{a_2 b_2} \dots \delta_{a_p b_p}$



▶ The complete-graph  $K_{p+1}$  is  $p$ -edge-colorable



$K_4$  ( $p = 3$ )



$K_6$  ( $p = 5$ )

Theorem: (Ferrari, Rivasseau, Valette '17)

A melonic large  $N$  limit exists for **prime**  $p$ .

( $p = 3$ : [Tanasa, SC '15])

# PRECURSORS: COLORED O(N) MODELS

( $p = 3$ )

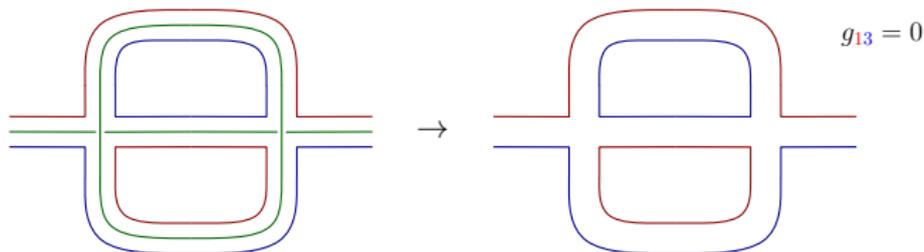
$$\frac{\lambda}{N^{3/2}} T_{aeb} T_{cfb} T_{ced} T_{afd}$$



- ▶  $A(G) \sim N^{-\omega}$  with  $\omega = 3 + \frac{3}{2}V - F \geq 0$
- ▶  $G$  leading order  $\Leftrightarrow \omega = 0 \Leftrightarrow G$  is a melon diagram

Idea of proof:

- ▶ Euler relation:  $\omega := g_{13} + g_{12} + g_{23} \in \frac{\mathbb{N}}{2}$ , where  $g_{ij}$  = genus of the jacket ( $ij$ ).
- ▶ melons are "super-planar" i.e. they have planar jackets



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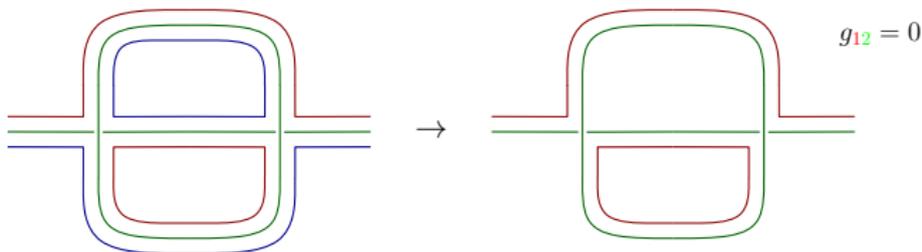
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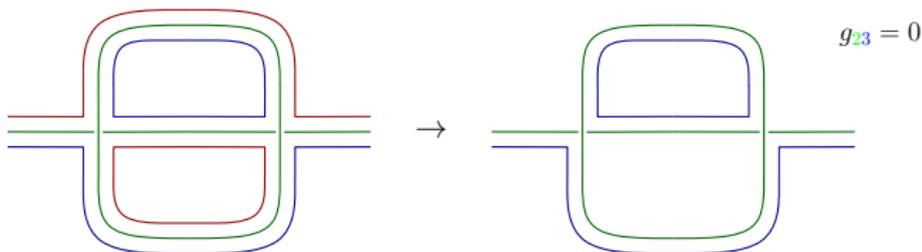
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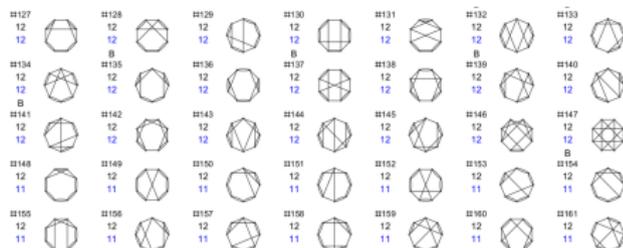


### Conjecture (Klebanov–Tarnopolsky '17)

For  $p = 3$ ,  $\exists$  melonic large  $N$  limit for  $O(N)$  symmetric **traceless** tensors.

Evidence. Explicit numerical check of all diagrams up to order  $\lambda^8$ .

[Klebanov, Tarnopolsky, JHEP '17]



### Proof and further generalizations.

1.  $O(N)$  irreducible,  $p = 3$

[Benedetti, SC, Gurau, Kolanowski, Commun. Math. Phys. '19; SC, JHEP '18]

2.  $Sp(N)$  irreducible,  $p = 3$

[SC, Pozsgay, Nucl. Phys. B '19]

3.  $O(N)$  irreducible,  $p = 5$

[SC, Haribey '21]

Much more involved and subtle constructions than in the colored case.

## IRREDUCIBLE TENSORS – PROPAGATOR

$\mathbf{P}$  = orthogonal projector on one of the irreducible tensor spaces.

example: for traceless tensors with symmetry

1	2
3	

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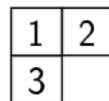
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$$\begin{aligned} \int d\mu_{\mathbf{P}}(T) T_{a_1 a_2 a_3} T_{b_1 b_2 b_3} &= \mathbf{P}_{a_1 a_2 a_3, b_1 b_2 b_3} = \frac{1}{3} (\delta_{a_1 b_1} \delta_{a_2 b_2} \delta_{a_3 b_3} - \delta_{a_1 b_3} \delta_{a_2 b_2} \delta_{a_3 b_1}) \\ &+ \frac{1}{6} (\delta_{a_1 b_2} \delta_{a_2 b_1} \delta_{a_3 b_3} + \delta_{a_1 b_1} \delta_{a_2 b_3} \delta_{a_3 b_2}) \\ &- \frac{1}{6} (\delta_{a_1 b_2} \delta_{a_2 b_3} \delta_{a_3 b_1} + \delta_{a_1 b_3} \delta_{a_2 b_1} \delta_{a_3 b_2}) \\ &+ \frac{1}{2(N-1)} (\delta_{a_1 b_3} \delta_{a_2 a_3} \delta_{b_1 b_2} + \delta_{a_1 a_2} \delta_{a_3 b_1} \delta_{b_2 b_3}) \\ &- \frac{1}{2(N-1)} (\delta_{a_1 b_1} \delta_{a_2 a_3} \delta_{b_2 b_3} + \delta_{a_1 a_2} \delta_{a_3 b_3} \delta_{b_1 b_2}) \end{aligned}$$

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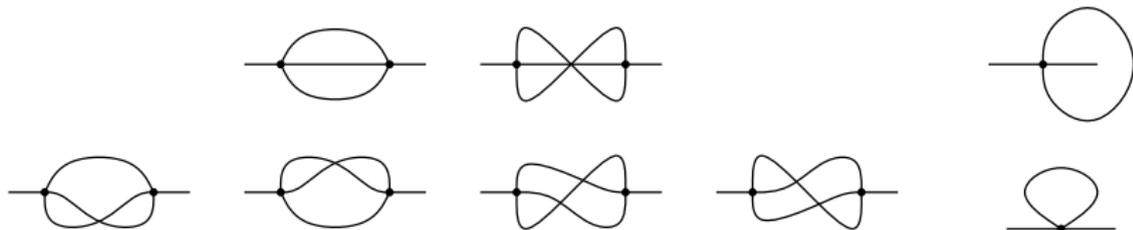
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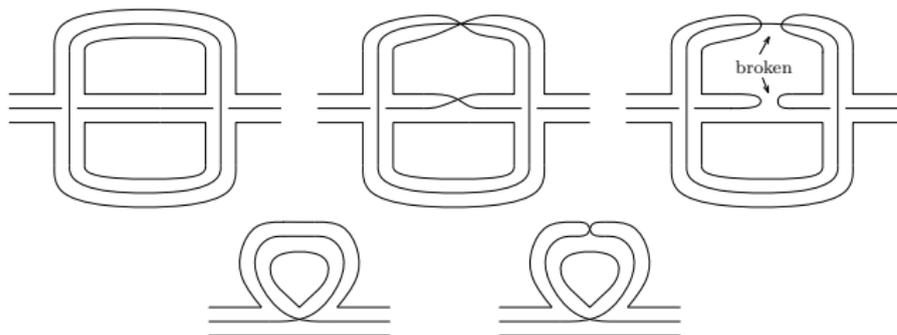
$$\begin{aligned}
 T_{a_1 a_2 a_3} \text{ ————— } T_{b_1 b_2 b_3} = & \frac{1}{3} \left( \begin{array}{c} a_1 \text{-----} b_1 \\ a_2 \text{-----} b_2 \\ a_3 \text{-----} b_3 \end{array} - \begin{array}{c} \diagdown \quad \diagup \\ \diagup \quad \diagdown \end{array} \right) \\
 & + \frac{1}{6} \left( \begin{array}{c} \text{||||} \\ \diagdown \quad \diagup \\ \diagup \quad \diagdown \end{array} + \begin{array}{c} \diagdown \quad \diagup \\ \text{||||} \\ \diagup \quad \diagdown \end{array} \right) \\
 & - \frac{1}{6} \left( \begin{array}{c} \diagdown \quad \diagup \\ \diagdown \quad \diagup \\ \diagup \quad \diagdown \end{array} + \begin{array}{c} \diagdown \quad \diagup \\ \diagdown \quad \diagup \\ \text{||||} \end{array} \right) \\
 & + \frac{1}{2(N-1)} \left( \begin{array}{c} \cup \quad \cup \\ \diagdown \quad \diagup \\ \diagup \quad \diagdown \end{array} + \begin{array}{c} \diagdown \quad \diagup \\ \cup \quad \cup \\ \diagup \quad \diagdown \end{array} \right) \\
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 \end{aligned}$$

## IRREDUCIBLE TENSORS – MAPS AND STRANDED GRAPHS

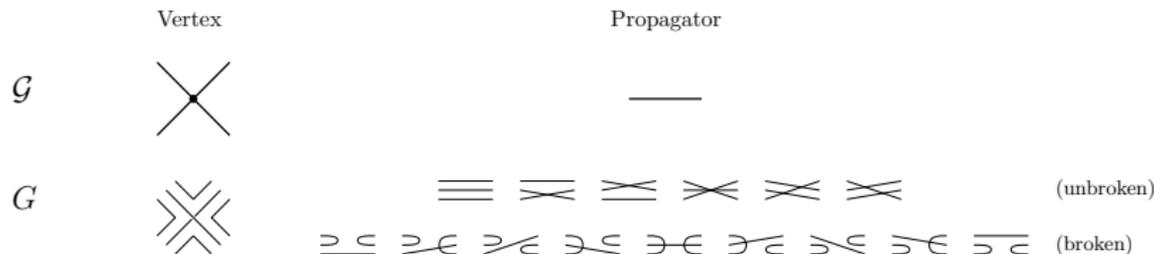
- Feynman expansion  $\rightarrow$  combinatorial maps  $\mathcal{G}$ :



- Decomposition of propagators  $\rightarrow$  stranded graphs  $\mathcal{G}$ :



# IRREDUCIBLE TENSORS – FEYNMAN AMPLITUDES



$$\mathcal{F}_N(\lambda) = \sum_{\text{connected maps } \mathcal{G}} \frac{\lambda^{V(\mathcal{G})}}{s(\mathcal{G})} A(\mathcal{G})$$

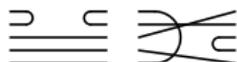
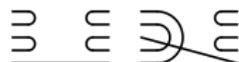
$\mathcal{G}$  decomposes into up to  $15^{E(\mathcal{G})}$  stranded graphs  $G$ :

$$A(\mathcal{G}) = \sum_G A(G), \quad A(G) \sim N^{-\omega(G)}$$

$$\omega(G) = 3 + \frac{3}{2}V(G) + B(G) - F(G)$$

$$V = \#\{\text{vertices}\}, B = \#\{\text{broken edges}\}, F = \#\{\text{faces}\}$$

## IRREDUCIBLE TENSORS – 5-INDEX TENSORS

*Unbroken**Broken**Doubly-broken*

Map  $\mathcal{G}$  decomposes into up to  $945^{E(\mathcal{G})}$  stranded graphs  $G$ :

$$A(\mathcal{G}) = \sum_G A(G), \quad A(G) \sim N^{-\omega(G)}$$

$$\omega(G) = 5 + 5V(G) + B_1(G) + 2B_2(G) - F(G)$$

$$B_1 = \#\{\text{broken edges}\}, \quad B_2 = \#\{\text{doubly - broken edges}\}$$

## MAIN THEOREMS

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# OUTLINE

Introduction

$O(N)$  irreducible random tensors with complete graph interaction

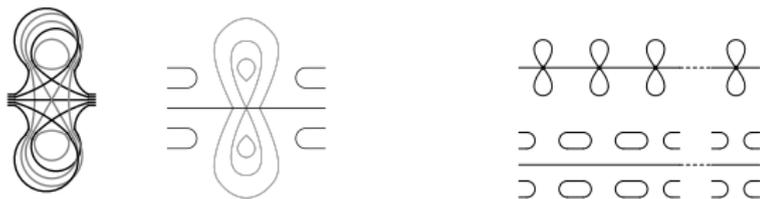
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Melonic dominance at leading order

# MAIN DIFFICULTIES

Natural conjecture: For any stranded graph  $G$ ,  $\omega(G) \geq 0$ .

× **Incorrect!** × Counter-example: chain of "bad double-tadpoles"



- ▶ Subclass of graphs with good scaling properties: **stranded graphs** containing **no melon** or **double-tadpole** subgraphs.



- ▶ No global constraint such as Euler's relation available  $\Rightarrow$  analysis of **local combinatorial structure** of  $G$ .

## PROOF STRATEGY

1. Eliminate **melon** and **double-tadpole** 2-point functions at the Feynman map level:

$$\begin{array}{c} \text{melon} \end{array} = \mathcal{O}\left(\frac{1}{N}\right) \text{---} \quad \begin{array}{c} \text{double-tadpole} \end{array} = \mathcal{O}(1) \text{---}$$

*This is where the **irreducibility assumption** plays a crucial role.*

2. Obtain  $\mathcal{G}$  with no melon and no double-tadpole.

**Proposition:** For any stranded configuration  $G$  of  $\mathcal{G}$ ,  $\omega(G) \geq 0$ .

*Proof.* Induction on  $V = \#\{\text{vertices}\}$ . Conceptually straightforward but challenging by its complexity.

# IDEA OF PROOF – COMBINATORIAL MOVES

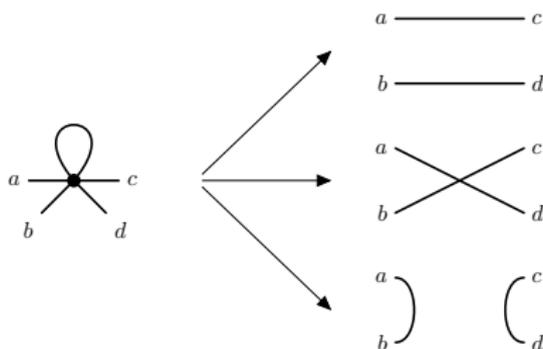
Find **local combinatorial moves** that:

- ▶ decrease  $V$ ;
- ▶ decrease  $\omega$ ;
- ▶ preserve constraints: connectedness,  $\emptyset$  melon,  $\emptyset$  double-tadpole.

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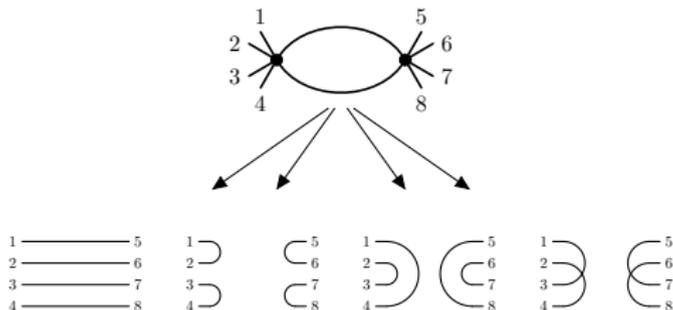
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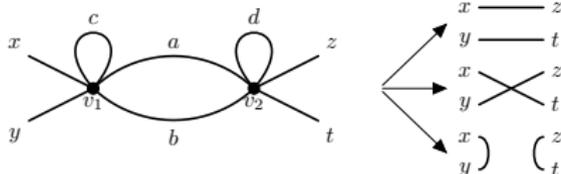
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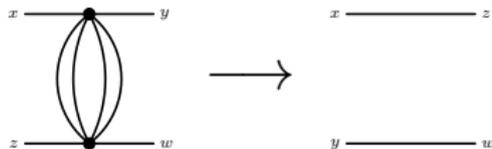
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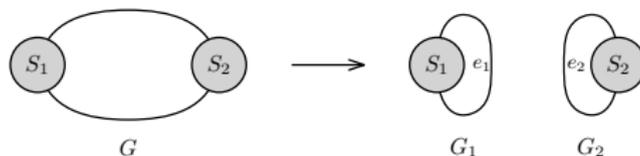
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Find *local combinatorial moves* that:

- ▶ decrease  $V$ ;
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*End graphs*

- ▶ Ring graphs ( $V = 0$ ):

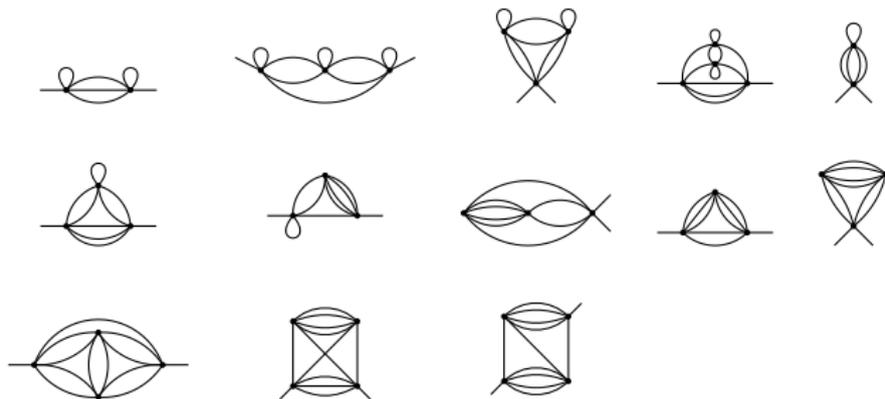


- ▶  $G$  with no face of length 1 or 2  $\Rightarrow \omega(G) > 0$ .
- ▶ Special cases that need to be treated separately.

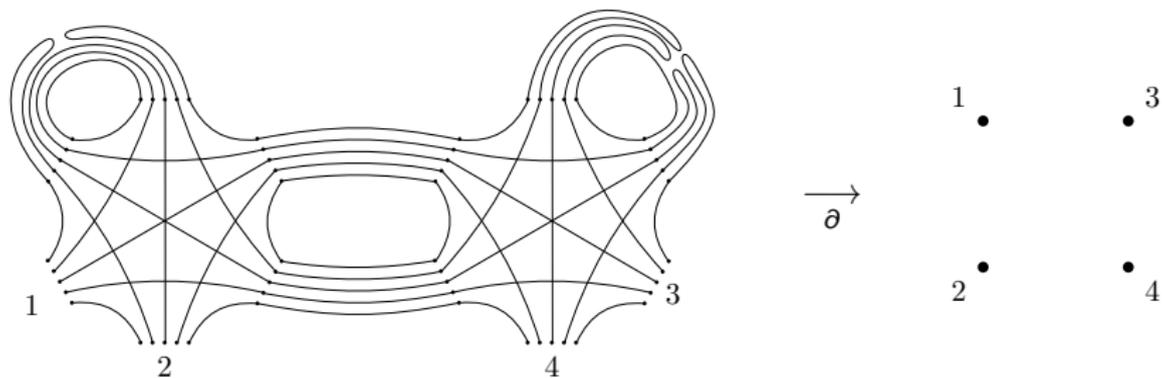
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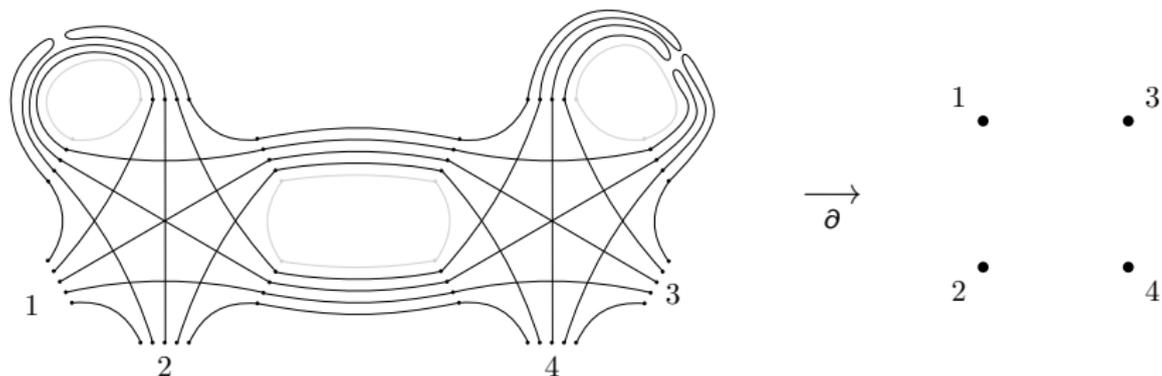
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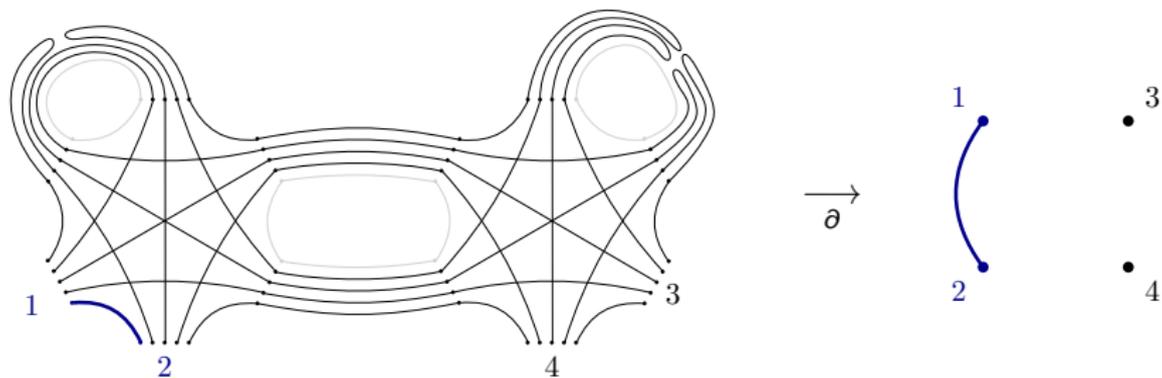
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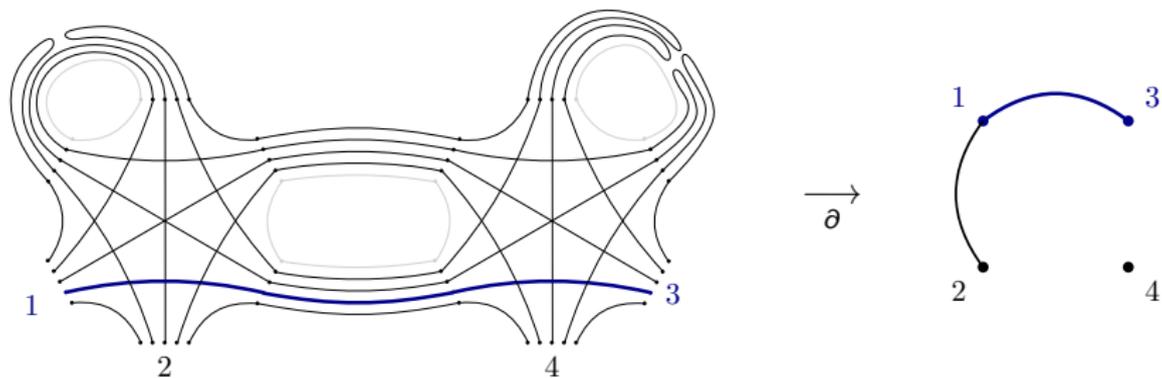
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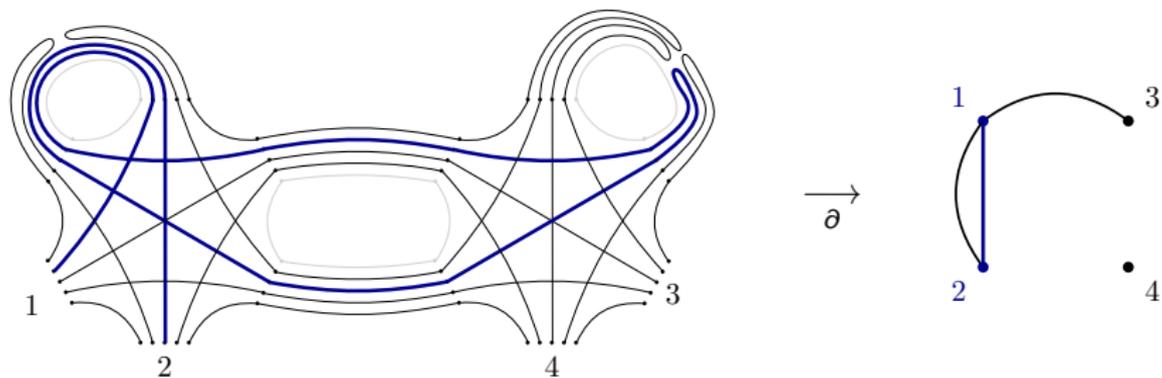
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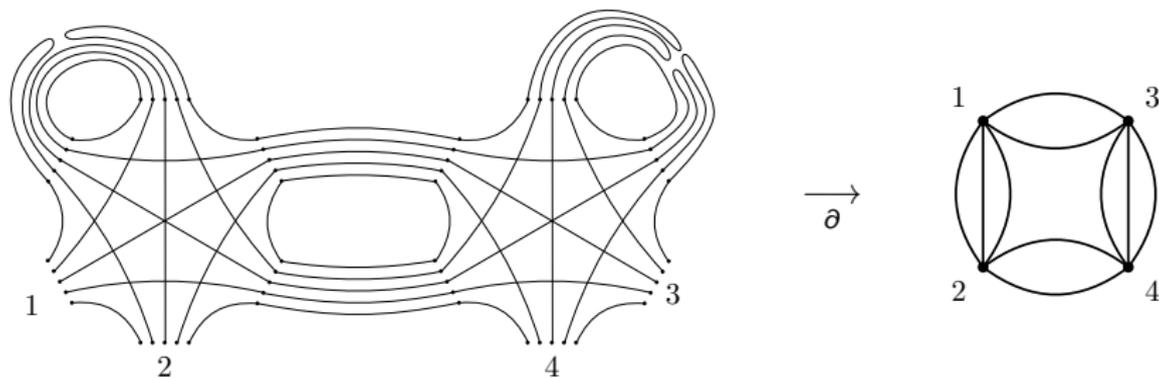
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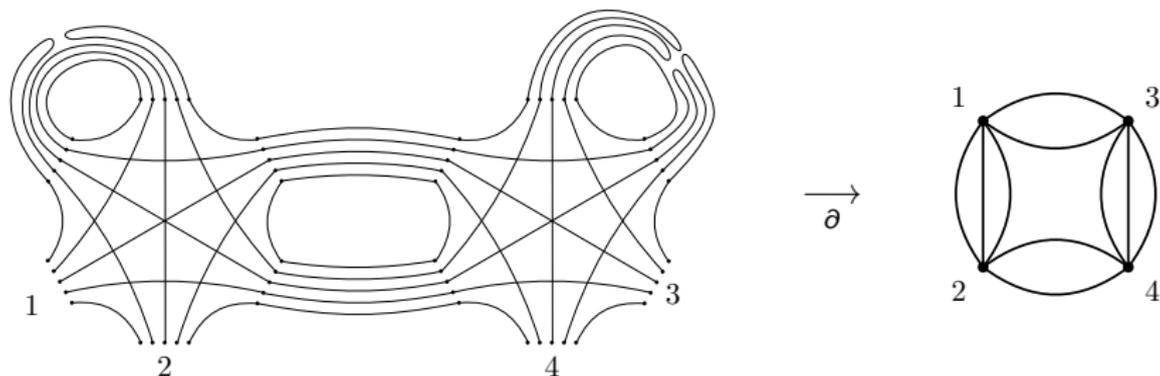
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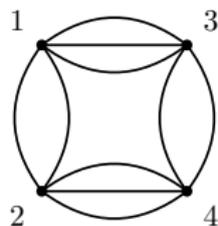
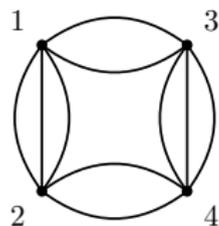
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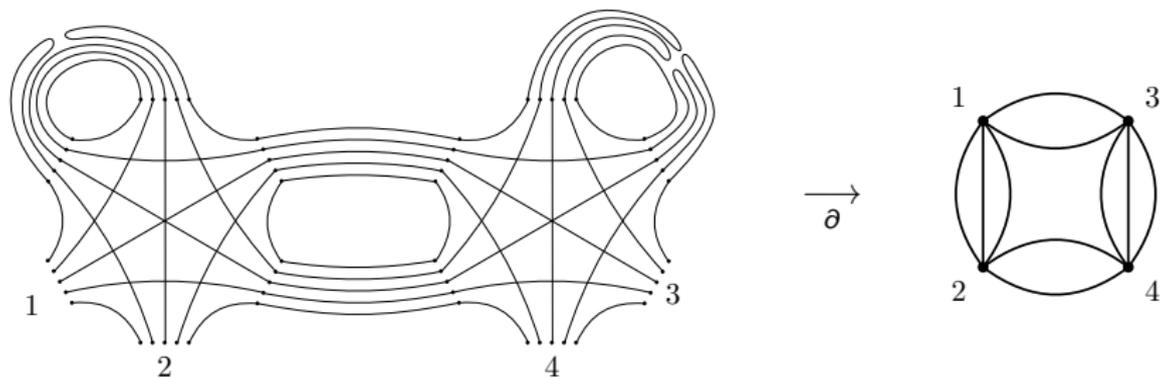
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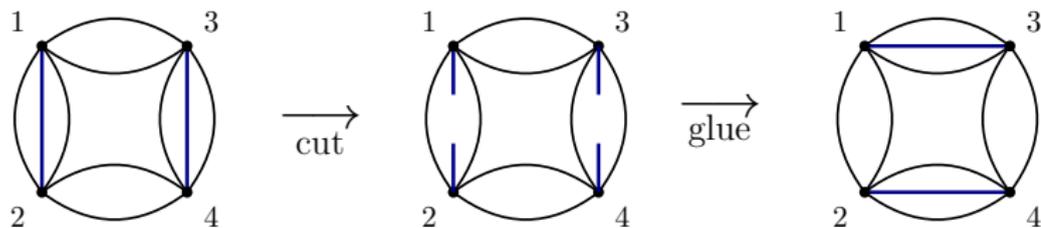
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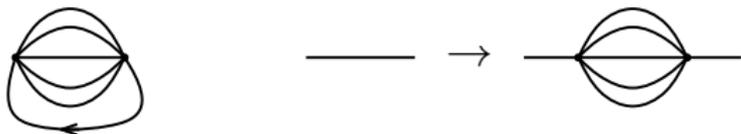
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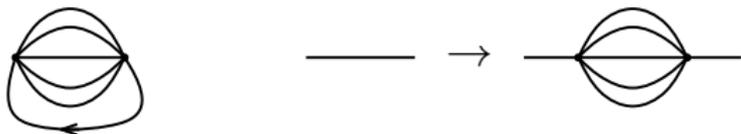


*Proof.* Cauchy-Schwarz inequalities on maps for which we do not already have strict bounds e.g.

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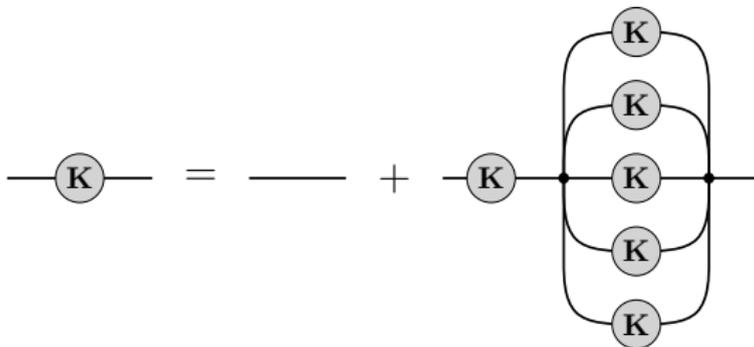


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# SCHWINGER–DYSON EQUATION

Hallmark of melonic limit: the 2-point function verifies a closed SDE



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## CONCLUSION AND OUTLOOK

Large  $N$  limit of  $p$ -index **irreducible** random tensors with  $p = 5$ :

Complete graph interaction +  $O(N)$  symmetry  $\Rightarrow$  melonic limit.

- ▶ In contrast to the matrix case, the **irreducibility condition is essential**.
- ▶ Other interactions, other symmetry groups (e.g.  $Sp(N)$ ), as well as fermionic tensors can be analyzed with the same method.
- ▶ Subleading orders can in principle be characterized as well.

Open questions:

- ▶ Generalization to **arbitrary (prime)  $p$** ?
- ▶ Useful applications to strongly-coupled QFT?
- ▶ Towards a general theory of random tensors?